



International Agreement Report

A Study of the Dispersed Flow Interfacial Heat Transfer Model of RELAP5/MOD2.5 and RELAP5/MOD3

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Abstract

In this work, the model of interfacial heat transfer for the dispersed flow regime used in the RELAP5 computer code is investigated. The validity of the 1-D approach used for calculating the heat exchange between the droplets and the vapour in the dispersed flow region above a quench front was shown to be questionable for conditions of low quality at the quench front and low mass flux. Under such conditions, the interfacial heat transfer calculated assuming a uniform distribution of droplets over the cross-sectional area of the channel is necessarily overpredicted, and the vapour superheat is strongly underpredicted. The purpose of the present paper is to show that the limitations of the 1-D approach, obtained from the steady-state analyses of slow reflooding experiments, has some impact on the performance of the 1-D transient computer codes like RELAP5/MOD2.5 and RELAP5/MOD3. As an example, the transient analysis of a low flooding rate experiment in a tube was performed. An early completion of the quench process and a fast desuperheating of the vapour at the tube exit was obtained by both codes. The too high quench front velocity (four times higher than in the experiment) could not, however, be put univocally in relation to the underprediction of the vapour temperature, and the consequent increase of the precursory cooling, as many coupled thermal and hydraulic transient effects prevailed. Quasi steady-state analyses of two runs, where the boundary conditions for the post-dryout region could be better controlled for a predetermined position of the quench front, were thus performed. These analyses show that the vapour superheat at the tube exit is strongly underpredicted, confirming the limitations of the 1-D model for interfacial heat transfer in the dispersed flow region.

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1 Introduction

An extensive work of assessment, development and applications of the RELAP5 and TRAC/B reactor safety computer programs has been underway for some time at the Thermal-Hydraulics Laboratory of the Paul Scherrer Institute (PSI). One of the reactor safety issues to which the research has been specially addressed is the reflooding phase of the LOCA: indeed, the accurate calculation of the rod temperature during this period is of importance for the estimation of the safety margins required for the reactor operation.

The main results of this research [1],[2],[3],[4] and the contributions to the code development in the areas of wall heat transfer [5],[6], interfacial drag [2],[6], and numerical techniques [6],[7] have been fully documented and examples of references have been provided here. For several reflooding experiments, it has been observed that both the trend and the values of the calculated results (wall and vapour temperatures, entrainment, collapsed liquid level, etc.) obtained by the 'frozen' version of the codes were very different to the measured experimental data. The discrepancies between the calculations and the data were noticeably large, especially for low flooding rate conditions. Low flooding rate experiments are characterized by very high void fractions all over the region above the QF for long time periods; in this zone droplets are dispersed in a stream of superheated vapour. The incorrect calculation of either the vapour superheat or the wall heat transfer coefficient resulted in the poor predictions of the 'frozen' versions of the codes.

The implementation of more consistent packages of correlations for wall heat transfer and interfacial shear (for both the wet and dry wall regions), the accurate selection of several empirical constants entering in the various models on the base of sensitivity studies, and the modification of interpolation schemes and numerical techniques lead to improved versions of the programs, which successfully simulate most experiments with *bundles* [2],[3],[5],[6],[7].

Modifications to the closure laws (largest droplet size, interfacial heat transfer) related to the dispersed flow regime above the Quench Front (QF) have also played a role in the achievement of good results, and more research is underway in this area. These studies are important contributions to the development of reliable calculation tools for reactor safety assessment, as the test facilities considered for the assessment are closely reproducing the conditions expected in the nuclear cores.

However, since a few quantities (average temperatures, entrainment rate, quench front velocity, etc.) are measured and can be used for the assessment, the adequacy of the closure laws used for the individual submodels can seldom be demonstrated. In principle, a fortuitous compensation of errors and the lumping of different effects in one single 'best fit' parameter, could mask the deficiencies in the models implemented in the codes. For instance, the lack of the representation of the spacer grids and an incorrect model for interfacial heat transfer and interfacial shear could be 'compensated' by a proper, although unphysical, choice of the empirical criterion for defining the droplet diameter.

In the present report, a different approach to the analysis of the reflooding process is presented, which is more suitable for evaluating the capabilities of one of the submodels used in the codes, namely that for interfacial heat transfer in the dispersed flow region. The attention to this particular model is motivated by the results of a recent investigation [8],[9] on post-CHF heat transfer with dispersed flow during reflooding. It has been shown that, under conditions of low mass flux and low quality at the QF, the interfacial heat transfer must be overpredicted by the 1-D models and some 3-D (such as those implemented in RELAP and TRAC^a series of codes) and this results in large underpredictions of the vapour superheat. It was argued that the neglect of the cross-sectional distribution of the droplets can be the reason for these discrepancies. The clustering of the droplets around the centre with the consequent reduction in the interfacial heat transfer has been proposed [8],[9] as a possible mechanism for explaining the high superheating rates (close to those admissible for 'frozen flow') at some distance from the QF observed in several tube experiments [10]: this conclusion was obtained using the commonly accepted closure laws for interfacial drag and droplet diameter.

The present analysis using RELAP5 is, therefore, aiming at a verification of such a general statement by a specific application of the 1-D approach. In this respect, the most interesting variable for judging the

^aThe TRAC code is often quoted as 3-D code, as it can provide information on the 3-D distributions *in the core*; the large sizes of the radial mesh which can practically be used, however, do not allow to calculate the radial profiles *inside the individual sub-channels*.

capability of the model is the vapour temperature at the exit of the test section, and this will be the main object of the discussions of the results.

For this purpose, it is convenient to assess the codes against data obtained in single-tube experiments; the simpler tube geometry allows to better test the interfacial heat transfer model, as the 'spurious' effects of the grids (which cannot be quantified and cannot be taken into consideration by the code, but are present in the experiments with bundles) on the droplet size and other variables, will not bias the comparison between the calculated vapour temperature and the experimental data. The exercise presented in this work is the transient analysis of the low flooding rate experiment 3051 at 2 bar in the single tube test section of the University of California-Berkeley (UCB) by means of frozen versions of RELAP5/MOD2.5 (Mod2, cycle 36.06)[11] and RELAP5/MOD3 (7vj)[12]; this experiment was chosen because the separate steady-state analyses [9] showed severe difficulties of the 1-D approach.

However, while the transient analysis of the core is the only important application of the reactor safety codes, much deeper insight in the influence of the interfacial heat transfer on the mixture evolution is made possible by considering 'snapshots' at different times during the transient. The values of several variables are considered constant over a short time period, and the analysis is referred to such a 'frozen' picture. The rationale for the consequent 'quasi steady-state analysis' is shortly outlined here, and discussed more in depth in §3. For the 'steady-state' calculation, the conditions above the QF must be imposed (approximately) constant over a short time, and the wall heat flux over the tube length above the QF must be given as input: if good predictions for both liquid and vapour mass flow rates are obtained, the resulting vapour temperature at the tube exit (which depends on the total heat exchanged with the droplets) will give a direct estimate of the adequacy of the interfacial heat transfer closure law used in the code. Indeed, the uniform heat flux which characterizes the single-tube experiments leads to slowly changing axial wall temperature distributions and allows a sufficiently good estimate of the local heat flux at all elevations. This, in turn, allows a rather accurate evaluation of the total power input to the mixture above the QF (on which the vapour superheat strongly depends) at any selected time, so that one of the unknowns which complicates the interpretation of the results of the transient analyses, namely the instantaneous 'actual' power input, can be eliminated.

Therefore, in the present work the quasi-steady state analyses of two UCB experiments (Run 3051 at 2 bar and Run 3053 at 3 bar) for the time of the mid-height quench are also carried-out.

2 Test cases used for the assessment

Well documented experimental results for the reflooding of tubes were obtained at the University of California at Berkeley [13],[14]. The diagram of the experimental test facility is shown in Fig. 1. The test section used for the second series of experiments at 2 and 3 bar was a 14.25 mm inside diameter (0.83 mm thickness) tube, 3.67 m in length. It was quenched by bottom injection of water at flooding rates from 25 to 180 mm/s: the water subcooling, about 80 K, was the same for all the experiments. Constant electrical heat input was used, and the cooling of the tube was initiated after an equilibrium condition had been reached: initial wall temperatures ranged between 316 °C and 760 °C. The tube was *not* insulated, and large heat losses have been allowed. The transient was initiated from a condition of constant wall temperatures, the power input being balanced by the heat losses. The initial axial wall temperature profile was practically flat.

The temporal variation of the tube wall temperature at different elevations and of the fluid temperature at the tube exit were obtained: the vapour temperature was measured by a sheathed microthermocouple (0.254 mm) located at about 5 cm above the tube exit. The flows of liquid and vapour at the tube exit were separated by means of a carefully designed separator, and the entrained liquid was collected in a tank. The vapour mass flow was also measured. Moreover, the water collected inside the test section (from which the collapsed liquid level can be calculated) was obtained by measuring the pressure difference between inlet and outlet of the test section, and neglecting friction and acceleration pressure drops.

During many tests, the wall temperatures at some distance from the QF set to nearly constant values for long times, being the heat input balanced by the heat losses and the heat transfer to the fluid: typical wall temperature time-histories are shown for Run 3051 (experimental conditions specified below) in Fig. 2.

In the present investigation, the conditions at the earliest time such a steady condition (the quality at the QF being positive) for the wall temperatures at the higher elevations is attained are chosen for the 'quasi steady-state' analyses.

In Table 1 the experimental conditions of the two runs used for assessing the codes are summarized.

For the quasi steady-state analyses, the additional information on the equilibrium quality at the QF x_{QF} and the axial distribution of the net heat input to the fluid q''_{in} at the selected time are also required.

The equilibrium quality immediately above the quench front x_{QF} is calculated [15] by adding to the quality immediately below the Quench Front x^- (calculated by an energy balance over the wet region) the quality change across the Quench Front due to the intense vaporization caused by the release of the heat stored in the tube wall:

$$x^+ = x^- + \frac{C_w A_w (T^+ - T^-) U_{QF}}{H_{fg} GA} \quad (1)$$

with:

$$x^- = \frac{P_{tot} \left(\frac{z_{QF}}{L} \right)}{H_{fg} GA} + \frac{\Delta H_{sub}}{H_{fg}}$$

where C_w is the volumetric heat capacity of the wall, A_w the cross sectional area of the wall, U_{QF} the quench front velocity, H_{fg} the vaporization enthalpy, G the total mass flux, A the cross sectional area of the channel, T^+ and T^- the wall temperatures above and below the quench front, ΔH_{sub} the liquid subcooling at the inlet of the test section, L the total length of the tube test section, z_{QF} the distance of the QF from the tube inlet, and P_{tot} the input power.

The net heat flux q''_{in} at each elevation is obtained by subtracting from the inside-wall 'imposed' heat flux q''_o (P_{tot} divided by the internal surface area), the heat losses $q''_{L,iw}$ (referred to the inside-wall surface area) and the sensible heat stored in the wall:

$$q''_{in} = q''_o - q''_{L,iw} - C'_w \frac{dT'_w}{dt} \quad (2)$$

where C'_w is wall heat capacity per unit inside wall area and T'_w the value of the wall temperature (in K) obtained by polynomial interpolation of the experimental temperatures. The outside heat flux (losses) is:

$$q''_L = h_L (T'_w - T_{amb}) \quad (3)$$

where h_L is the global heat transfer coefficient (HTC) from the wall to the surroundings, assumed to be at temperature $T_{amb}=294$ K. The experimental values of h_L could be correlated by the formula:

$$h_L = 11.949 e^{[2.592 \cdot 10^{-3} (T'_w - T_{amb})]} \quad (\text{W/m}^2 \text{ K}) \quad (4)$$

The heat losses per unit area (heat flux to the ambient) referred to the inside-wall heat transfer area is obtained by:

$$q''_{L,iw} = q''_L \frac{D_e}{D}$$

where D_e and D are the external and internal diameters, respectively.

The total heat input to the fluid in the dry region above the Quench Front P_{dry} is calculated by integrating q''_{in} over the wall heat exchange area.

The values of the several variables required for the quasi steady state analyses are given in Table 2 together with the transient times to which they refer.

Run No.	p bar	G kg/m ² s	P_{tot} kW	$T_{w,in}$ K
3051	2.	25.1	4.326	812.
3053	3.	25.	4.326	812.

Table 1: Experimental conditions for test cases UCB-3051 and UCB-3053

Run No.	t s	z_{QF} m	U_{QF} mm/s	T^- K	T^+ K	x^-	x_{QF}	P_{dry} W	q''_{in}	
									$z(m)$	kW/m ²
3051	150.	1.90	7.3	401.	593.	0.10	.12	923.	2.00	16.91
									2.20	15.45
									2.40	13.92
									2.60	12.38
									2.80	10.97
									3.00	9.69
									3.20	8.83
									3.40	8.19
									3.58	8.19
									3053	150.
2.40	14.20									
2.60	13.60									
2.80	12.98									
3.00	12.10									
3.20	10.97									
3.40	9.90									
3.58	9.89									

Table 2: Calculated variables used for the steady-state analyses.

3 Transient analysis of test case 3051

The basic nodalization scheme for the analysis of Run 3051 is shown in Fig. 3. It consists of one pipe (component 120) with 19 cells for the test section, one time-dependent junction (110) for the bottom injection, one junction (125) representing the tube exit, and two time-dependent volumes (105 and 130). The heated wall of the test section is represented by the same number of heat slabs (19), and uniform linear power generation is imposed. The heat losses (calculated from a temperature-dependent heat transfer coefficient h_L) to the ambient depend on the external wall temperature, according to correlations (3) and (4). The values of the heat transfer coefficient h_L has been calculated for several temperatures, and a table of h_L vs temperature has been used for the input deck. The ambient temperature has been imposed constant and equal to 294 K. At the beginning of the transient ($t=0$ s) the test section is filled with superheated steam in thermal equilibrium with the wall of the tube ($T_g=T_{w,in}=812$ K). The transient is initiated by starting ($t=0$ s) the bottom injection of water.

The transient was calculated with the 'frozen' versions of both RELAP5/MOD2.5 and RELAP5/MOD3. For both calculations the correlation for interfacial drag in tubes was chosen. The reflooding option was chosen for the calculation with both versions of RELAP5, so that moving fine heat structure meshes were used in the wall heat transfer calculation. The reflooding heat transfer package was then utilized in the calculation with RELAP5/MOD2.5, whereas this option does not affect the correlations used by RELAP5/MOD3, as a unique heat transfer package is used for all conditions. The input decks for the two calculations are given in App. I.

The calculation by RELAP5/MOD2.5 could not be completed using the nodalization with 19 vols. shown in Fig. 3 (left), as a condition of error (which could not be explained) occurred, at about 160 s, in the subroutine QFMOVE, which calculates the advancement of the quench front (at this time, the quench front elevation was already above 2.44 m). Reduction of maximum and minimum time steps, and changes in the number of radial meshes and maximum number of fine meshes only provoked a negligible time shift of the error condition. Eventually, the solution was to halve the size of the hydrodynamic nodes (from 0.2 to 0.1 m) and, correspondingly, of the heat slabs. The resulting nodalization with 36 vols. is also shown in Fig. 3.

The comparisons of the wall temperatures calculated by both codes with the measurements are shown in Figs. 4 and 5. Even though some minor differences can be observed in the results of the two calculations, the most striking outcome of the analysis is the very large discrepancy between calculated and measured wall temperatures, as in [4]. The test section quenches completely in about 550 s, while both codes predict complete quench within 160 s. While RELAP5/MOD3 predicts correctly a gradual bottom quench progression, RELAP5/MOD2.5 calculates a second quench front, propagating from the top of the test section, as it can be recognized (Fig. 5) from the early quench at 3.35 m (highest measurement station). The results are somewhat surprising, if one recalls that most of the analyses with bundles (e.g., NEPTUN with cosine-shaped heat flux profiles, but also with several spacer grids) indicated the tendency of RELAP5/MOD2.5 to *overpredict* the quench times and the wall peak temperatures, especially at the highest elevations [4]. These results are also in conflict with recent assessment work using data from the Lehigh bundle reflooding experiments, where a uniformly heated bundle was used: for these cases RELAP5/MOD2.5 was generally predicting peak temperatures and quench times with acceptable agreement in comparison to experimental data, and the only underpredictions were found for low elevations [4]. This last assessment work can, however, be barely associated to the previous ones, because of the short length of the test section; in fact, the analysis was focussed on lower elevations than in the NEPTUN case.

RELAP5/MOD3, on the contrary, behaves in accordance with the previous experience [4]: it overpredicts the quench front velocity (or rather the 'cooling front' velocity, as no evident temperature 'knee' can be observed). For the present tube experiment the effect is amplified, as the QF velocity is four times higher than in the experiment.

Moreover, it can be observed (Fig. 6) that early in the transient (30 s for the MOD2.5 and 100 s for the MOD3), the calculated vapour temperatures in the uppermost node drop to low values (about 420–450 K) close to the saturation temperature (about 400 K), whereas the experimental value sets for a long time to a value of around 600 K.

The reason for such a large disagreement between the predicted and the measured quench front velocities is not easy to single out, as several effects can contribute to the inaccuracy of the predictions. On the one hand, too large wall heat fluxes might have driven the wall to rewetting conditions, with consequent drop in the vapour temperature to the saturation value. Moreover, large mass flow oscillations affect both code versions (in Figs. 7 and 8 the vapour and liquid carry-out, as well as the total mass flow rate at the tube exit are shown), and most probably contribute to the excessively fast propagation of the quench front. On the other hand, the large temperature difference between the wall and the vapour might have resulted in an additional increase of the already over-predicted wall heat flux [6] and, consequently in an overprediction of the precursory cooling.

From the comparison of the calculated mass accumulated inside the tube, the collapsed liquid level, the vapour flowing out of the test section and the total liquid carry-out (Fig. 9), with their respective experimental counterparts, one can conclude that both code versions predict the 'average' hydraulic behaviour of the system fairly well.

It should be noted here that the calculation of the total liquid carry-out by RELAP5/MOD2.5 is actually better than what appears in Fig. 9, as the control variable (cntrlvar) which integrates the liquid expelled over the time is affected by an error, and underestimates the total liquid carry-out. A short explanation can be given as follows. The mass flow rate of the phase k at a junction j (in the specific case the junction 125) is obtained from:

$$\dot{M}_{k,j} = \rho_{k,j} \alpha_{k,j} v_{k,j} A$$

and the total mass flow can be calculated by summing up the gas and the liquid mass flows. The total mass expelled out of the test section \dot{M}_{out} is then:

$$\dot{M}_{out} = \text{int}(\text{cntrlvar}) = \int_0^t (\dot{M}_{f,j} + \dot{M}_{g,j}) dt \quad (5)$$

On the other hand, the total carry-out can be calculated by integrating directly the total mass flow at the tube exit \dot{M}_j , which is also calculated by the code:

$$\dot{M}_{out} = \text{int}(\text{mflowj1250000}) = \int_0^t \dot{M}_j dt \quad (6)$$

The comparison of the results of the two expressions is given in Fig. 10: it is observed that a 10% difference exists in the RELAP5/MOD2.5 calculation, while a much smaller discrepancy appears in the RELAP5/MOD3 calculation. As the value of \dot{M}_{out} calculated by Eq. (6) satisfies the mass balance with a high accuracy, it can be concluded that the mass flows calculated by Eq. (5) underestimate the 'actual' mass flows. Indeed, under rather unstable liquid mass flow conditions, the spikes in the mass flow rate (due to the spikes in the liquid mass flow rate) calculated by Eq. (5) are much lower than those present in the mass flow time-history of the variable mflowj1250000 (see Figs. 7 and 8). This can be explained by the different numerical techniques used for obtaining the time and space averages of the various quantities at the junctions in these two different versions of the code.

Even though the comparison between calculations and experiment regards only the exit conditions, it can be assumed that both wall and interfacial drag are correctly predicted over the transient^b and that the axial liquid mass distribution has also been captured with a good approximation.

Using the hypothesis that the low vapour temperatures drive the wall heat flux increase and the wall temperatures drop (and not vice versa), one could then argue that the failure of the RELAP5/MOD3 in predicting vapour and wall temperature, is an overprediction of the interfacial heat transfer, which results in an anticipated desuperheating of the vapour up to large distances from the quench front. An indication in favour of this explanation can be obtained from observing, for example, that at the time of the quench at 2.58 m ($t=93$ s, interpolated value) the vapour temperature at the tube exit (1.2 m downstream) is about 550 K, whereas

^bThis does not necessarily imply that interfacial and wall drag models are correct.

in the experiment at the time this happens ($t=298$ s) the vapour temperature was 50 K higher. However, the comparison is not fully meaningful as one cannot compare the actual quality at the quench front (highly oscillating in the calculation because of the high heat transfer rates across it) and the total heat input to the vapour in the dry region. In fact, one also observes that, because of the high cooling rates (due also to the large wall-to-liquid heat transfer) at large distances from the QF, the calculated and experimental wall temperatures and temperature decrease rates in the dry region are quite different at the time of the 2.58 m quench. This implies different values of the local heat fluxes and, then, of the total heat input to the vapour. Therefore, because of the complex interactions between the different heat transfer processes, the transient effects and the history-dependent thermal-hydraulic conditions, the boundary conditions which influence the vapour superheating rate are different in the calculation and in the experiment, so that the model for interfacial heat transfer in the dispersed flow regime cannot be judged conclusively.

In the case of the MOD2.5 calculation, the occurrence of top quench confuses the picture even more, as the sudden decrease of the vapour temperature close to the exit is not due to the excessive interfacial heat transfer, but to the fact that the vapour cannot receive heat from the wet wall any longer and a large quantity of saturated vapour is generated from the direct wall-to-liquid heat exchange.

From the analysis of this tube experiment, it can be deduced that:

1. The analysis of experiments with tubes can yield rather different results than those obtained for bundles, even for very similar conditions. The results for tubes, even though they are not truly representative of the conditions in a nuclear reactor, provide a useful test for the adequacy of the correlations used in the computer codes, as most of the models implemented are primarily derived for circular channels (the application to bundles implies the use of some 'equivalent' diameter). This can be recognized from the excellent prediction of the hydraulics.
2. The calculated quench front progression is much faster than in the experiment for both versions of the codes. The discrepancy is due to a too large wall heat flux [6] in the post-CHF regime, to which the underprediction of the vapour temperature also contributes.
3. Run 3051 provides a very severe test for the codes. The difficulty of the 1-D approach in predicting interfacial heat transfer demonstrated for more specific conditions [9], cannot be definitely confirmed here, but the large discrepancy in the calculated vapour temperature seems to be consistent with those phenomenological considerations. In particular the MOD3 calculations seem to be affected by the overprediction of interfacial heat transfer, but both calculations are probably dominated by deficiencies in the wall heat transfer models.
4. The transient analysis above does not allow to draw any conclusion about the correctness of the individual models, as they (and their numerical implementation) interact with each other.

With respect to the main goal of the present work, namely the assessment of the interfacial heat transfer model for Dispersed Flow, not much can be concluded from the transient analysis; hence, the analysis of experiment 3053 (for which results similar to those obtained for test 3051 are expected) was not carried-out.

4 Quasi-steady state analysis of Run 3051

The basic idea behind the quasi steady-state analysis has already been introduced in §1; here the mathematical formulation of this approach, as well as the conditions which must be fulfilled in order to arrive to the evaluation of the capabilities of the interfacial heat transfer model using only the experimental vapour temperature at the tube exit, are discussed.

When the quenching process progresses slowly and at constant rate, as usual when the quality at the QF is positive, the thermal-hydraulic conditions at any point in the channel downstream from the QF (wall and fluid temperatures, mass flows, qualities, etc.) vary slowly as well. The quasi steady-state condition is thus realized if the time necessary to the QF to move a significant distance (e.g., the distance between the centres

of the volumes above and below the 'nominal' position of the QF) is larger than the residence time of the droplet in the post-CHF region. If this prerequisite is satisfied, any slow variation of the QF position (and quality) will lead to slow variations of the hydrodynamic variables downstream. Under this condition, the evolution of the mixture can be described by the usual conservation equations, dropping the time-derivatives. The steady-state form of the vapour continuity and energy equation can be written for a straight tube (like in the RELAP5 code [11], but neglecting dissipation and using the enthalpy instead of the internal energy) of cross-sectional area A as:

$$\frac{d}{dz}(\langle \alpha_g \rangle \langle \rho_g \rangle U_g) = \frac{dG_g}{dz} = \langle \Gamma \rangle \quad (7)$$

$$\langle \alpha_g \rangle \langle \rho_g \rangle U_g \frac{dH_b}{dz} = G_g \frac{dH_b}{dz} = Q_{wg} + \langle Q_{ig} \rangle - \langle \Gamma \rangle (H_b - H_{g,s}) \quad (8)$$

where U_g , H_b , G_g , Q_{wg} and $\langle Q_{ig} \rangle$ are the cross-sectionally averaged axial vapour velocity, the bulk vapour enthalpy, the vapour mass flux, the phasic wall heat transfer rate per unit volume and the cross-sectionally averaged interfacial heat transfer rate per unit volume, respectively.

By integrating between the QF elevation z_{QF} and the tube exit elevation z^T Eq. (8) multiplied by the tube flow area A , one obtains:

$$\int_{z_{QF}}^{z^T} AG_g \frac{dH_b}{dz} dz = \int_{z_{QF}}^{z^T} Q_{wg} Adz + \int_{z_{QF}}^{z^T} \langle Q_{ig} \rangle Adz - \int_{z_{QF}}^{z^T} \Gamma (H_b - H_{g,s}) Adz$$

and, integrating by parts the first term in the Left Hand Side of the equation:

$$\begin{aligned} AG_g^T H_b^T - AG_{g,QF} H_{b,QF} - \int_{z_{QF}}^{z^T} AH_b \frac{dG_g}{dz} dz &= \int_{z_{QF}}^{z^T} Q_{wg} Adz + \\ &\int_{z_{QF}}^{z^T} \langle Q_{ig} \rangle Adz + \\ &- \int_{z_{QF}}^{z^T} \Gamma (H_b - H_{g,s}) Adz \end{aligned}$$

Using the continuity equation (7) and replacing $G_g A$ by the vapour mass flow rate \dot{M}_g , it follows:

$$\begin{aligned} \dot{M}_g^T H_b^T - \dot{M}_{g,QF} H_{b,QF} &= \int_{z_{QF}}^{z^T} Q_{wg} Adz + \int_{z_{QF}}^{z^T} \langle Q_{ig} \rangle Adz + \int_{z_{QF}}^{z^T} \Gamma H_b Adz \\ &- \int_{z_{QF}}^{z^T} \Gamma (H_b - H_{g,s}) Adz \end{aligned}$$

Observing that:

$$\int_{z_{QF}}^{z^T} \Gamma Adz = \dot{M}_g^T - \dot{M}_{g,QF}$$

one can write (assuming saturated vapour at the QF):

$$\dot{M}_g^T H_b^T = \dot{M}_{g,QF} H_{b,QF} + P_{g,dry} + \int_{z_{QF}}^{z^T} \langle Q_{ig} \rangle Adz + (\dot{M}_g^T - \dot{M}_{g,QF}) H_{g,s} \quad (9)$$

where $P_{g,dry}$ is the total heat transfer rate to the vapour alone in the dry-wall region. If all the droplet population is represented by an average diameter d , the average interfacial heat transfer rate per unit volume is:

$$\langle Q_{ig} \rangle = \frac{6h_i}{d} \langle (1 - \alpha_g)(T - T_s) \rangle \quad (10)$$

Since in a 1-D model the droplet concentration is assumed uniform over the cross section and the distribution coefficient is taken equal to unity, the relation:

$$\langle (1 - \alpha_g)(T - T_s) \rangle \equiv \langle (1 - \alpha_g) \rangle \langle (T - T_s) \rangle \equiv \langle (1 - \alpha_g) \rangle (T_b - T_s) \quad (11)$$

leads to:

$$\langle Q_{ig} \rangle = \frac{6h_i}{d} \langle (1 - \alpha_g) \rangle \langle (T_b - T_s) \rangle \quad (12)$$

This is the relation that under low mass flux and low quality conditions, will lead to the overprediction of the interfacial heat transfer, and the underprediction of the vapour temperature [9]. To show that RELAP5 is affected by large inaccuracies in the prediction of the vapour temperatures because of the large interfacial heat transfers, and then that Eq. (12) does not hold, is the goal of the following analysis.

From Eq. (9) one easily realizes that if the gas mass flow rate at the QF and the tube exit calculated by the computer code are (approximately) equal to the experimental ones, and the total wall-to-vapour heat transfer is the same as in the experiment, any deviation of the exit enthalpy (temperature) can be attributed to the incorrect prediction of the interfacial heat transfer. This estimation can be done *irrespective of the wall-to-liquid heat transfer* and the consequent vapour production at the wall.

Therefore, for the quasi steady-state analysis four conditions (Fig. 11) must be fulfilled in the calculation:

1. the time needed by the the QF to advance a substantial length about the 'reference' QF position is much longer than the droplet residence time t_r ($\equiv (z^T - z_{QF})/\bar{U}_l$). The average values of the variables during this time (which we will call 'averaging period' t_{av}) can be compared with the values obtained by a steady-state analysis.
2. same (approximately) x_{QF} ($\dot{M}_{g,QF}$) as in the experiment during the averaging period t_{av} .
3. total heat input to the vapour equal to the experimental one.
4. same \dot{M}_g^T as in the experiment.

The first two conditions can be imposed by restricting the analysis to a section of the tube including the region downstream of the 'nominal' position of the QF and a short length upstream of it. The fluid injected at the bottom of this 'reduced' length section is a mixture with a quality close to that inferred from the experiment. The QF progresses from the inlet of the section up to the nominal location of the QF (and the quality at the QF changes accordingly), for which the tube exit vapour temperature must be compared with the experimental value. From this time on, and until the time the first node above z_{QF} quenches, the values of the various variables of interest are considered.

With respect to the third condition, certain assumptions and approximations are needed, since only the total heat flux to the mixture is known at the few measurement stations. As the experimental wall temperatures reach nearly steady values at some time during the run, and drop suddenly only when the QF is very close to the measurement location, the interpolation of the heat fluxes for any elevation of the tube can be calculated with good accuracy, and the total heat input to the mixture in the dry region of the tube can then be estimated with reasonable confidence: with this procedure the values in Table 2 were obtained. Still the partition between the two phases is unknown, so that two approaches can be used in the calculation:

- Impose the total heat flux equal to the experimental one, and let the code partition the heat between the phases. If this option is chosen, the wall-to-liquid heat transfer must be small at least at some distance from the QF, otherwise an underprediction of the vapour temperature can be due to the low heat input to the vapour rather than to the overprediction of the interfacial heat transfer.

- If the code calculates a too large wall-to-liquid heat transfer, increase the total heat flux as much as necessary to transfer to the vapour an amount of heat equal (approximately) to the total experimental one. This second option assumes no wall-to-liquid heat transfer in the experiment. Therefore, this procedure uses the upper limit for the heat input to the vapour, and, if the vapour temperature is still underpredicted, the deficiency in the approach for calculating the interfacial heat transfer will become more evident.

For the fourth condition, namely calculated vapour mass flow close to the experimental value, one can only hope that the computer program predicts liquid and vapour carry-out flow rates close to the experimental ones at the time of the mid-height quench. This expectation is conformed by the results of the transient analysis above, where both calculated liquid and vapour carry-out were in excellent agreement with the measured ones. For the purpose of testing the interfacial heat transfer model, lower vapour mass flows than in the experiment can also be accepted, as the calculated vapour superheat is enhanced by a reduced vapour mass flow. Similarly, the liquid carry-out must be approximately equal to or lower than the experimental value. A calculated liquid mass flow much higher than in the experiment would imply a large concentration of liquid at high elevation: the consequent vapour desuperheating effect could not be related to the 1-D approach for the interfacial heat transfer.

In order to avoid misunderstandings, it is appropriate to repeat here that, as RELAP5 is a transient code, our 'steady-state' analysis consists in running the program over the short time period needed by the QF to advance from the inlet of the tube to the node immediately above the elevation z_{QF} , chosen for the analysis. In summary: The analysis will be restricted to a shorter tube length. The quench front moves from the inlet to a certain predetermined elevation above z_{QF} . From the time the QF is close to z_{QF} , the quality immediately above must be close to x_{QF} , and the power input to the vapour in the dry region must be equal to $P_{g,dry}$; this last condition can be approximated by imposing that the total heat transfer rate to the mixture is equal to the (experimental) P_{dry} . The mass flows at the tube exit will be observed; if they are in close agreement with the experimental data (obtained by differentiating the experimental water and vapour carry-out), the vapour temperature at the tube exit will be compared with the experimental data.

4.1 Calculation with RELAP5/MOD2.5

In Fig. 12 the nodalization used for the 'steady-state' analysis with the MOD2.5 is shown. It consists of ten short (0.03 m) cells below z_{QF} , and nine cells representing the above-QF region having the same size as in the 19 vols. nodalization (Fig. 3). The inlet quality x_{in} was set by a trial-and-error procedure, by checking the value of x_{QF} at the end of the transient, as calculated from an energy balance performed by means of a control variable; at the end x_{in} has been set equal to about 6%. The need for this procedure is due to the oscillations in the calculation of the mass flows immediately above the QF, which render the direct estimation of the quality impossible.

The heat flux at the wall was initially imposed according to Table 2, the initial vapour temperature in the dry region was given a linear profile (the experimental temperature at $t=150$ s being used for the highest node), and the experimental wall temperature was imposed to the heat structures. Unfortunately, the calculation stopped after a few seconds, and the reason was easily found: the vapour temperature upstream of the quench front was exceeding 3000 K. This happened because a very large amount of heat was transferred directly to the vapour, in spite of the fact that a bubbly flow regime was predicted. It was realized that RELAP5, for imposed heat flux conditions, calculates the heat fluxes to the two phases by splitting the total heat input into two *equal* parts, irrespective of the flow regime.

It was then necessary to change our approach of imposing the heat flux, and other boundary conditions were used:

- Impose the power generation.
- Impose the experimental heat losses at $t=150$ s.

- Impose the axial temperature profile. As the heat flux is then calculated by the reflooding heat transfer package, the total heat input to vapour is not controlled, but varies during the short period of time of the 'quasi steady-state' analysis, and depends mainly on the initial wall temperature. Therefore, three different axial wall temperature profiles are imposed, and the resulting heat input to the vapour is calculated. The three axial wall temperature profiles, shown in Fig. 13, are:

1. exp: experimental wall temperature distribution.
2. low: temperatures decreasing with the distance from the QF. These values have been found (by trial and error) to yield (approximately) the desired axial heat flux distribution (Fig. 14).
3. high: temperatures higher than those experimentally observed.

The resulting wall heat fluxes averaged over the last 5 s of the calculations are shown in Fig. 14. The input deck for the calculation using the experimental temperature distribution is given in App. II.

The results of the three computations are shown in Figs 15-21. In Fig. 15 the wall temperatures at several elevations are shown for the low initial wall temperature distribution. The temperature time-histories in the other two cases were similar. Two remarks can be done: 1) the wall temperatures at some distance from the QF ($z \geq 2.2$ m) are practically constant during the calculation time, and 2) in the time period between 10 and 15 s the QF lies somewhere between 1.855 m and 2.0 m. The time needed for the QF to advance from 1.855 m to 2.0 m was found to be larger than the droplet residence time t_r ($\approx 1-2$ s), so that the first requirement for the 'steady-state' analysis is satisfied.

It is first observed that the equilibrium quality at the QF (Fig. 16) varies during the transient time, but is close to the experimental value at least for times larger than 10 s.

The vapour mass flow (Fig. 17) at the tube exit is also in close agreement with the experimental value. On the contrary, the liquid mass flow rate (Fig. 17) exhibits large oscillations: their amplitude, however, do not exceed 30% of the measured average flow rate. The fact that the average values (between 10 and 15 s) of the total mass flow rate and exit quality (Fig. 18) are quite close to the experimental values, suggest that the hydraulic behaviour is in general well predicted.

The total power input to the vapour (Fig. 19) is, in all cases, at least as high as the total heat input to the mixture, so that also the thermal boundary condition on the vapour superheat evolution is satisfied in both cases.

Under the conditions mentioned above, the average values of the several quantities of interest for our analysis during the time period from 10 to 15 s are thus comparable with those obtained by a steady-state analysis for QF 'frozen' at 1.89 m; then, it does make sense to compare the calculated vapour temperature with the measured one. In Fig. 20, it can be observed that after 10 s the vapour temperature (while being unstable) is always underpredicted by at least 50 K, even when the power input to the vapour is twice as much as the experimental total heat transfer rate to the mixture (Fig. 19). It is easily realized that the spikes are due to instantaneous increases of the void fraction due to sudden drops in the liquid mass flow rate (liquid carry-out). In Fig. 21 the average value of the vapour temperature is shown, and the large underprediction (150 K) observed for the initial wall temperature distribution which yields value of P_{dry} close to the experimental ones is in fairly good agreement with that expected on the base of the quasi steady-state separate model [9]. From the RELAP5/MOD2.5 calculations, we can thus conclude that the interfacial heat transfer must be overpredicted. Let us discuss the possible sources contributing to the overprediction:

1. The liquid fraction might have been overestimated. This, however, can be almost excluded by the fairly good prediction of the liquid carry-out, which suggests a correct estimate of the axial liquid distribution.
2. The interfacial heat transfer coefficient (h_i) is possibly overpredicted because of a large average value of the phase velocity difference. Indeed, one can suspect that large oscillations in the liquid velocity (leading even to negative velocities) can still be compatible with a correct average value of the liquid

mass flow, but yield (because of the non-linear dependence of the heat transfer coefficient on the phase velocity difference ($U_g - U_l$)) large values of h_i . This can be also excluded by observing that the velocity difference at two elevations (Fig. 22) are quite stable.

3. The droplet diameter might be underpredicted. The separate steady-state analyses [8] show, however, that for a wide range of droplet diameters the results are very similar to those obtained using the RELAP5/MOD2.5 criterion.
4. The only possibility that is left is, thus, that the interfacial heat transfer rate is based on the incorrect assumption of a uniform droplet concentration over the cross-sectional area.

Other calculations (not shown) have been carried-out using the normal (blowdown) heat transfer package, and slowing down the propagation of the quench front by increasing the thickness of the heat slabs in the region below the QF. The vapour temperature at the tube exit was still underpredicted by at least 100 K.

Furthermore, a version of the code was generated, in which the wall-to-liquid heat transfer was forced to go to zero within 30 cm from QF. The axial wall temperature profile was imposed in such a way to preserve the total experimental power input to the dry region (P_{dry}), so that the heat transfer rate to the vapour ($P_{g,dry}$) was practically equal to the total experimental one. Even under these conditions, the calculated vapour temperature was still much lower than in the experiment.

4.2 Calculation with RELAP5/MOD3

For the analysis with RELAP5/MOD3, the nodalization had to be slightly modified. This was due to the reason that the first calculation (not shown by a figure) showed that the quench front was progressing much faster than in the corresponding RELAP5/MOD2.5 calculation. Under these circumstances, the time necessary for the QF to propagate from the node below z_{QF} to the node above it (time t_{av} over which the thermal-hydraulic conditions of the post-CHF region can be considered approximately insensitive to the actual location of the quench front around the imposed z_{QF}) became comparable with the droplet transit time: the averaging of the variables over the time period t_{av} was consequently becoming of dubious meaning. Therefore, larger nodes (50 mm instead of 30 mm) were used for the section below z_{QF} and the node including z_{QF} (Fig. 23). Moreover, the time, over which the average values of the variables discussed above are calculated, is made dependent on the transient, i.e., the averaging period is between the time the wall temperature in the cell below z_{QF} (cell 9) drops below the quench temperature (580 K) and the time this happens to the first node in the 'dry' zone (cell 11). When the wall temperature in node 11 drops below 580 K the transient is thus terminated.

The first calculation, performed by using initial wall temperatures equal to the experimental ones, also showed a very high heat transfer rate in the dry zone, nearly double as much than in the experiment, and a much larger vapour mass flow, resulting from the large evaporation rates caused by the high wall-to-liquid heat transfer.

Therefore, the low axial wall temperature distribution (Fig. 13) was used for the next calculations. The results are shown in Figs. 24-29. The wall temperatures are shown in Fig. 24: at large distance from the QF they are slowly growing, and the QF lies between 1.825 and 2.0 m for the time between 8 and 12 s. The equilibrium quality at the QF is shown in Fig. 25. It can be remarked that it is highly oscillatory (often, the case found in the RELAP5/MOD3 calculations [4]), due to the large spikes in the heat transfer from the wall to the fluid in the wet region. The lower value of the equilibrium quality, which is representative of the 'steady state' heat transfer rate, is, however, close to the expected value.

The vapour temperature at the tube exit (Fig. 26) is predicted quite well, but one can easily realize from the good prediction of the vapour carry-out (Fig. 27) and the negligible liquid carry-out (Fig. 28) that this is due to the deficiency of liquid in the region far away from the QF. Fig. 29 shows the void fractions at various elevations. The water is accumulated in the region immediately downstream from the QF, so that no heat sink is available to the steam further up. It was originally believed that one of the reasons for the insufficient liquid carry-over could be the larger value of the critical Weber number ($We_{cr}=12$) used in the MOD3 with

respect to that used in the MOD2.5 ($We_{cr}=3$), which implies the calculation of quite large droplets in the mist-flow region and these drops cannot be lifted by the drag force (see below).

It was then decided to perform the calculation with a modified version of the code developed at PSI (version bb) which, among other features, uses the same We_{cr} as in MOD2.5. This modified version has been discussed by Analytis [6], and the description of its features will not be repeated here: it suffices to say that better numerical stability is achieved and several unphysical effects (like the void fraction passing through a minimum immediately below the quench front) are eliminated by the utilization of more physical correlations. The results are shown in Figs. 24-28 and 30. It can be observed that the liquid carry-out (Fig. 28) is a little higher than before, but still much too low: however, the resulting slight modification in the axial void fraction profile (Fig. 30) reduces the average outlet vapour temperature (Fig. 26) by about 30 K.

In order to understand the reasons why the utilization of the same criterion for the droplet size yields different results in RELAP5/MOD2.5 and in RELAP5/MOD3, the overall calculation scheme of the interfacial drag in RELAP5/MOD3 was reviewed. It was realized that, for reflooding calculations, the droplet diameter was bound by a lower limit of 2.5 mm. This value appears arbitrary, and for sensitivity analysis purposes it was changed to 1.25 mm in a new version of the code, which was called UPX. Moreover, it has been found out that the drag coefficient for droplets is defined as:

$$C_D = \min\{C_{D,v}, 0.45\} \quad (13)$$

where the drag coefficient for the viscous regime $C_{D,v}$ is given by:

$$C_{D,v} = \frac{24}{Re}(1. + 0.1Re^{0.687}) \quad (14)$$

It is clear from Eqs. (13)-(14) that for large Re (large droplets) the drag coefficient attains very low values. However, the drag coefficient for a solid sphere is asymptotically approaching the value 0.45 for $Re > 1000$ (Newton regime), being *higher* for lower values of Re [16]. The value of C_D for fluid particles is even higher, as they get distorted [16]. Eq. (13) is then basically *wrong* and even if one wants to consider the droplets as solid spheres, it must be replaced by the usual criterion:

$$C_D = \max\{C_{D,v}, 0.45\} \quad (15)$$

Eq. (15) was then implemented in the version UPX.

The results of the calculation by version UPX are shown in Figs. 25-28 and 31. It can be remarked that the liquid carry-out (Fig. 28) is now larger than in the previous calculations, but it is still too low. The axial void fraction distribution (Fig. 31) shows a larger presence of liquid at the higher elevations, and this determines lower vapour temperature (Fig. 26).

It can be observed that the calculated heat transfer rate to the mixture (Fig. 32) was always higher than in the experiment, but the heat transfer rate to the vapour was about 30% lower (Fig. 33). However, radiation from wall-to-droplets can account for a significant portion of the wall heat transfer [9], so that the actual heat transfer rate to vapour $P_{g,dr}$ is likely to be predicted realistically.

Further attempts to modify the code in order to get the correct liquid mass flow were not performed. As a trend, it has been already established that, for decreasing values of the droplet diameter, more and more liquid can be carried over, and increasing underprediction of the vapour temperature results.

These results are in agreement with those obtained by RELAP5/MOD2.5 and by the steady-state analyses.

5 Quasi-steady state analysis of Run 3053 by RELAP5/MOD2.5

Experiment 3053 was practically the same as Run 3051, the only difference being the higher pressure (3 bar instead of 2). The analysis has been performed only by the RELAP5/MOD2.5, because of the difficulties

in the prediction of the liquid carryout by the MOD3, which render the interpretation of the results more difficult.

The nodalization used for the analysis of run 3053 is shown in Fig. 34. Only, two axial wall temperature distributions (Fig. 35) have been tested. The resulting average axial wall heat flux profiles during the time period between 10 and 14 s are given in Fig. 36.

The transient time for the calculations was 14 s, since the node above the 'nominal' value of z_{QF} quenches at that time. The time-histories of the wall temperature at various elevations for the 'experimental' initial axial wall temperature distribution are shown in Fig. 37. The time span over which the results are representative of a steady quench front at 2.05 m (QF between 2.055 m, vol. 10, and 2.2 m, vol. 12) is between 8 and 14 s.

The results of the calculations are shown in Figs. 38-41. It is noted that the quality at the QF is correctly obtained (Fig. 38) and the heat transfer to the vapour alone is much larger than the experimental *total* power input (Fig. 39). The calculated liquid carry-out (Fig. 40) and the calculated total mass flow at the exit (Fig. 41) decrease from larger values than in the experiment to lower values during the averaging period. The calculated vapour carry-out (Fig. 42) is, however, larger than in the experiment, due to an excessive evaporation in the tube. Therefore, one condition necessary for the fully meaningful comparison (see §3) between calculated and experimental vapour temperature is not satisfied. We note, however, two circumstances: 1) the experimental total mass flow was probably larger than that measured in the test, as a certain error in the mass balance was found [14], and 2) the heat input per unit volume of the flowing vapour is at least as large as in the experiment, since $P_{g,v}$ is substantially larger than the total experimental heat flux (Fig. 39).

It is thus possible to state that the large underprediction of the vapour temperature at the tube exit (Fig. 43) is again due to the overprediction of the interfacial heat transfer, and not to an overprediction of the vapour mass flow.

6 Conclusions

The main goal of the present work was the assessment of the interfacial heat transfer model for the Dispersed Flow Film Boiling region above a quench front implemented in RELAP5/MOD2.5 and RELAP5/MOD3. The interest for this specific aspect of the DFFB model was raised by a separate work [9] which suggested severe limitations of the 1-D models in calculating the interfacial heat transfer rate under low mass flux conditions for low qualities at the quench front.

To this aim, a low-flooding rate reflooding experiment at 2 bar has been analyzed by RELAP5/MOD2.5 and by RELAP5/MOD3. The transient analysis has shown for both codes a large disagreement between the calculated and experimental quench front propagation rates, which can be attributed to an inadequate modelling of the heat transfer phenomena in the vicinity of the QF and confirms other similar analyses, e.g. [4]. In fact the good prediction of the carry-out suggests that the hydrodynamic behaviour is fairly well simulated. The vapour temperature at the tube exit dropped to the saturation value much before than in the experiment, but it could not be conclusively assessed whether a too high interfacial heat transfer was responsible for such bad predictions. This was due to the fact that the total instantaneous power input to the mixture in the post-CHF region at the time the quality at the QF becomes positive, was much larger than in the experiment, as it depends on the axial transient cooling rate distribution and then on the actual wall temperature distribution above the QF, which results from the preceding transient.

In order to control the power input to the mixture, a quasi-steady state analysis has been performed. A short transient was run, during which the quench front is allowed to advance a short distance around the position of interest, imposing (approximately) the quality at the quench front. The power input to the mixture can be controlled, to a certain extent, by imposing the initial wall temperature distribution above the quench front. The RELAP5/MOD2.5 calculated correctly the total carry-out and overpredicted the total heat transfer to the vapour. In spite of this, the vapour temperature was substantially underpredicted. It could be so concluded that interfacial heat transfer was overpredicted.

RELAP5/MOD3 showed a large underprediction of the liquid carry-out, and this led to a good prediction of the vapour temperature. However, when the calculation was repeated using two other versions of the code where the closure laws affecting the droplet dynamics were modified in a way to allow a larger carry-out (though still lower than the experimental one), the vapour temperature was again underpredicted.

The steady-state analysis of a second experiment at 3 bar was also performed by RELAP5/MOD2.5, and results similar to those obtained for the test at 2 bar were obtained.

It can be concluded that, as expected from steady-state analysis, computer codes like RELAP5, which calculate the interfacial heat transfer rate between vapour and droplets assuming a homogeneous distribution of the droplets over the cross section, overpredict the desuperheating effect of the liquid, and hence underpredict the vapour temperature at large distances from the quench front. This underprediction of the vapour temperature, does not always lead to the underprediction of wall temperatures, as other deficiencies in the wall heat transfer modelling and/or hydrodynamics often mask the difficulty of the codes to estimate the correct interfacial heat transfer. Nevertheless, we hope that this work will motivate some efforts to develop some empirical approach to calculate the reduction in the interfacial heat transfer rate under low mass flux/low quality conditions, since a detailed 3-D model which can calculate these effects (like that developed by one of the authors [9]) cannot be easily incorporated in a system code.

Acknowledgements

The interest of Prof. Yadigaroglu for the practical implications of the main results of Andreani's dissertation and his encouragement to carry-out the present investigation were appreciated. The help of M. Sencar in setting up the nodalization and running the codes is gratefully acknowledged.

References

- [1] S.N. Aksan, G.Th. Analytis and D. Lübbesmeyer, 'Switzerland's Code Assessment in support of the International Code Assessment and Application Program (ICAP)', *Proceedings of the 16th Water Reactor Safety Information Meeting* (Oct. 1988), NUREG/CP-0097, Vol. 4, pp. 245-273, March 1989.
- makebox[20pt][l][2]M. Richner, G.Th. Analytis and S.N. Aksan, 'Assessment of RELAP5/MOD2, cycle 36.02, using NEPTUN Reflooding Experimental Data', NUREG/IA-0054, August 1992 (PSI report No. 104, Oct. 1991)
- [3] S.N. Aksan, F. Stierli and G.Th. Analytis, 'Boil-off Experiments with the EIR-NEPTUN Facility: Analysis and Code Assessment Overview Report', NUREG/IA-0040, March 1992 (PSI-Report, No. EIR-629, Sept. 1987).
- [4] M. Sencar and S.N. Aksan, 'Evaluation of Reflooding Models in RELAP5/MOD2.5, RELAP5/MOD3/v7j Codes by using Lehigh University and PSI-NEPTUN Bundle Reflooding Experimental Data', Paper presented at the 2nd *Code Assessment and Maintenance Program (CAMP) Meeting*, Brussels, Belgium, May 1993.
- [5] G.Th. Analytis, 'A Comparative Study of the Post-CHF Wall Heat Transfer Packages of the RELAP5 codes and Preliminary Assessment of Model Changes in RELAP5/MOD3/v7j' Paper presented at the 1st *Code Assessment and Maintenance Program (CAMP) Meeting*, Villigen, PSI, Switzerland, June 1992.
- [6] G.Th. Analytis, 'A Summary of Model Changes and Options in RELAP5/MOD3' Paper presented at the 2nd *Code Assessment and Maintenance Program (CAMP) Meeting*, Brussels, Belgium, May 1993.
- [7] G.Th. Analytis and P. Coddington, 'The Effect of Donor Cell Differencing of some Terms of the Momentum Equations of TRAC-BF1 on the Code Predictions', *Proceedings of the 5th International Topical Meeting on Reactor Thermal-Hydraulics (NURETH-5)*, Vol. III, pp. 745-750, Sept. 21-24, 1992, Salt Lake City, Utah, USA.
- [8] M. Andreani and G. Yadigaroglu, 'Effect of the Cross-Sectional Droplet Distribution in Dispersed Flow Film Boiling at Low Mass Flux', *Proceedings of the 5th International Topical Meeting on Reactor Thermal-Hydraulics (NURETH-5)*, Vol. III, pp. 823-831, Sept. 21-24, 1992, Salt Lake City, Utah, USA.
- [9] M. Andreani, 'Studies of Dispersed Flow Film Boiling with 3-D Lagrangian Hydrodynamics and a 2-D Eulerian Vapour Field', Doctoral Dissertation Nr. 9794, Swiss Federal Institute of Technology, ETHZ, Zurich (1992).
- [10] D.G. Evans, S.W. Webb and J.C. Chen, 'Measurement of Axially Varying Non-Equilibrium in Post-Critical-Heat-Flux Boiling in a Vertical Tube', NUREG/CR-3363, Lehigh-Univ., June 1983.
- [11] Ransom, V.H. et al. 'RELAP5/MOD2 Code Manual', NUREG/CR-4312, Rev. 1, March 1987.

[12] Carlson, K.E. et al. 'RELAP5/MOD3 Code Manual', Vol. 1, 2, 3, and 4, Draft, NUREG/CR-5535, June 1990.

[13] Seban, R., Grief, R., Yadigaroglu, G., Elias, E., Yu, K., Abdollahian, D. and Peake, W. 'UC-B Reflood Program: Experimental Data Report', EPRI NP-743, 1978.

[14] Seban, R.A. 'Reflooding of a Vertical Tube at 1,2 and 3 Atmospheres', Report EPRI NP-3191, 1983.

[15] Yu, K.P. 'An experimental investigation of the Reflooding of a bare tubular test section', Dissertation, Univ. of California at Berkeley, October 1978.

[16] Ishii, M. and Mishima, K. 'Two-Fluid Models and Hydrodynamic Constitutive Relations', Nucl. Eng. and Design, 82 (1984), 107-126.

FIGURES

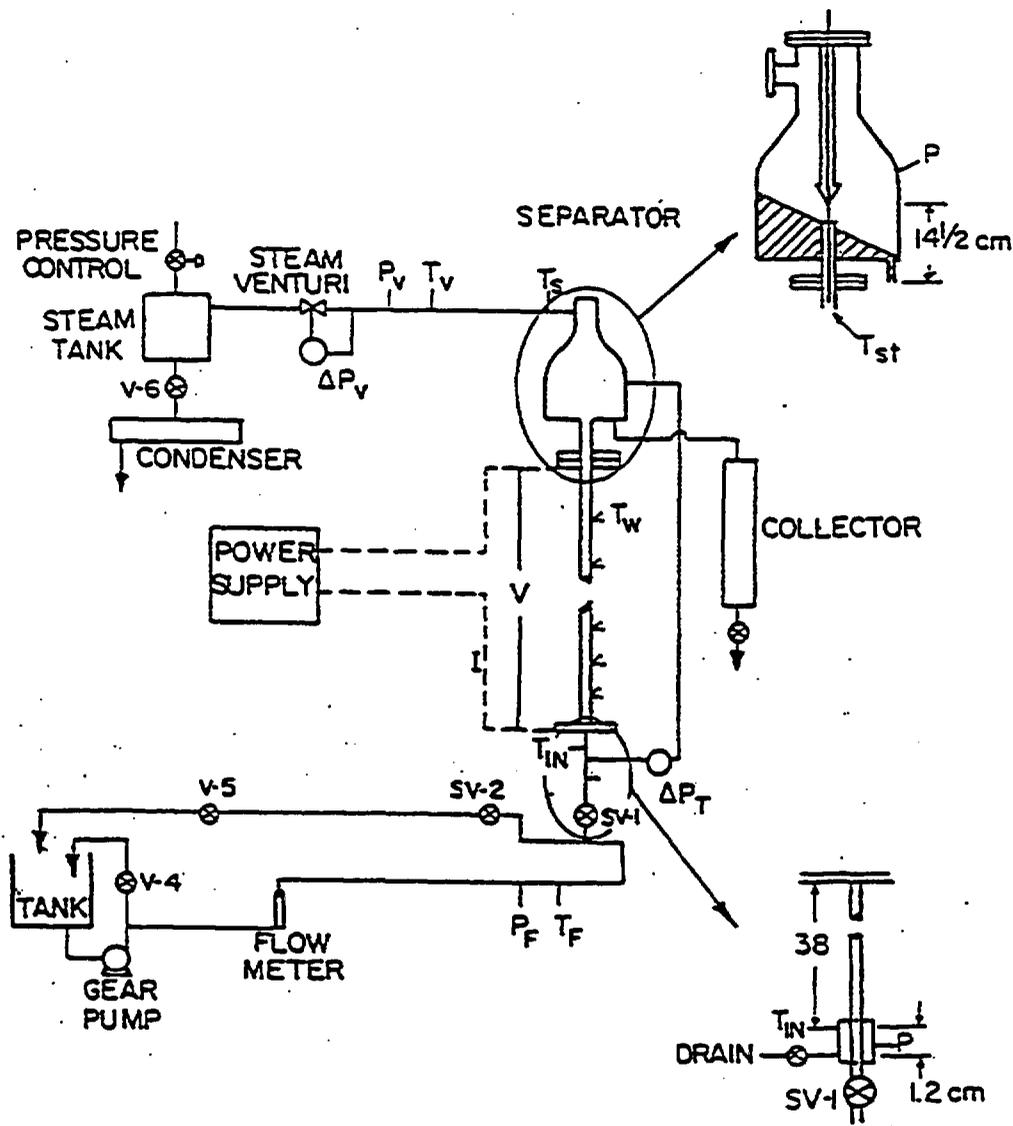


Figure 1: System flow diagram of the UC-B Test facility (Seban et al., 1978).

RUN #3051 (10-1-2)

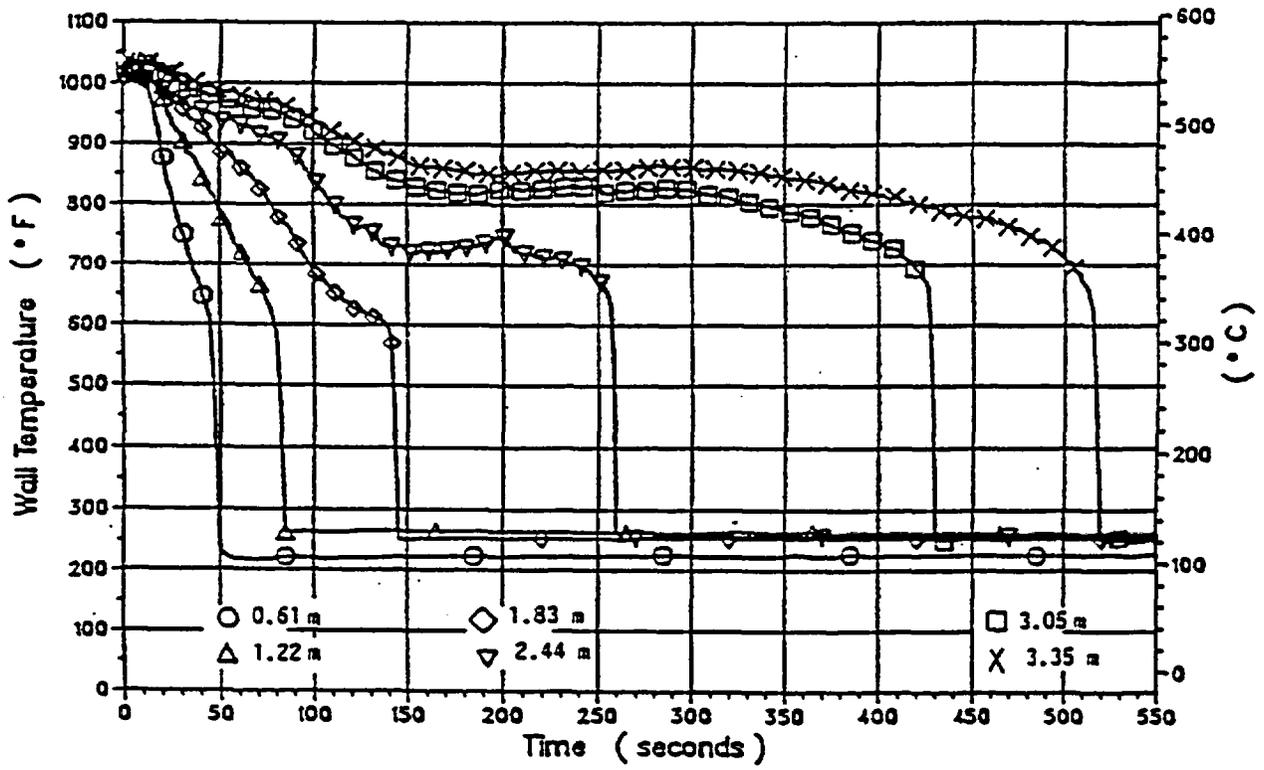


Figure 2: Experimental wall temperatures at various elevations for Run 3051.

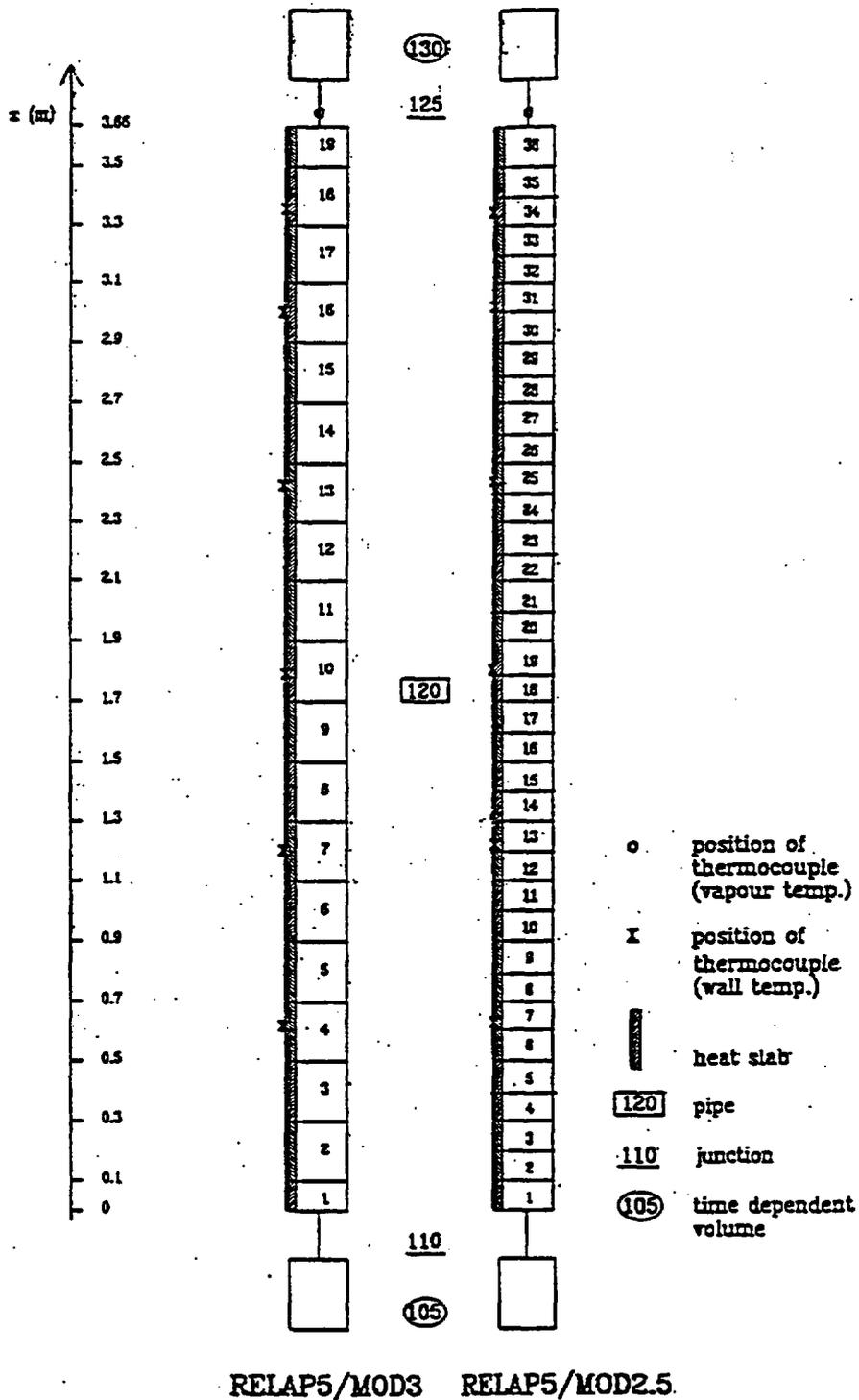


Figure 3: Nodalizations used for the calculations of Run 3051 with RELAP5/MOD2.5 (right) and RELAP5/MOD3 (left).

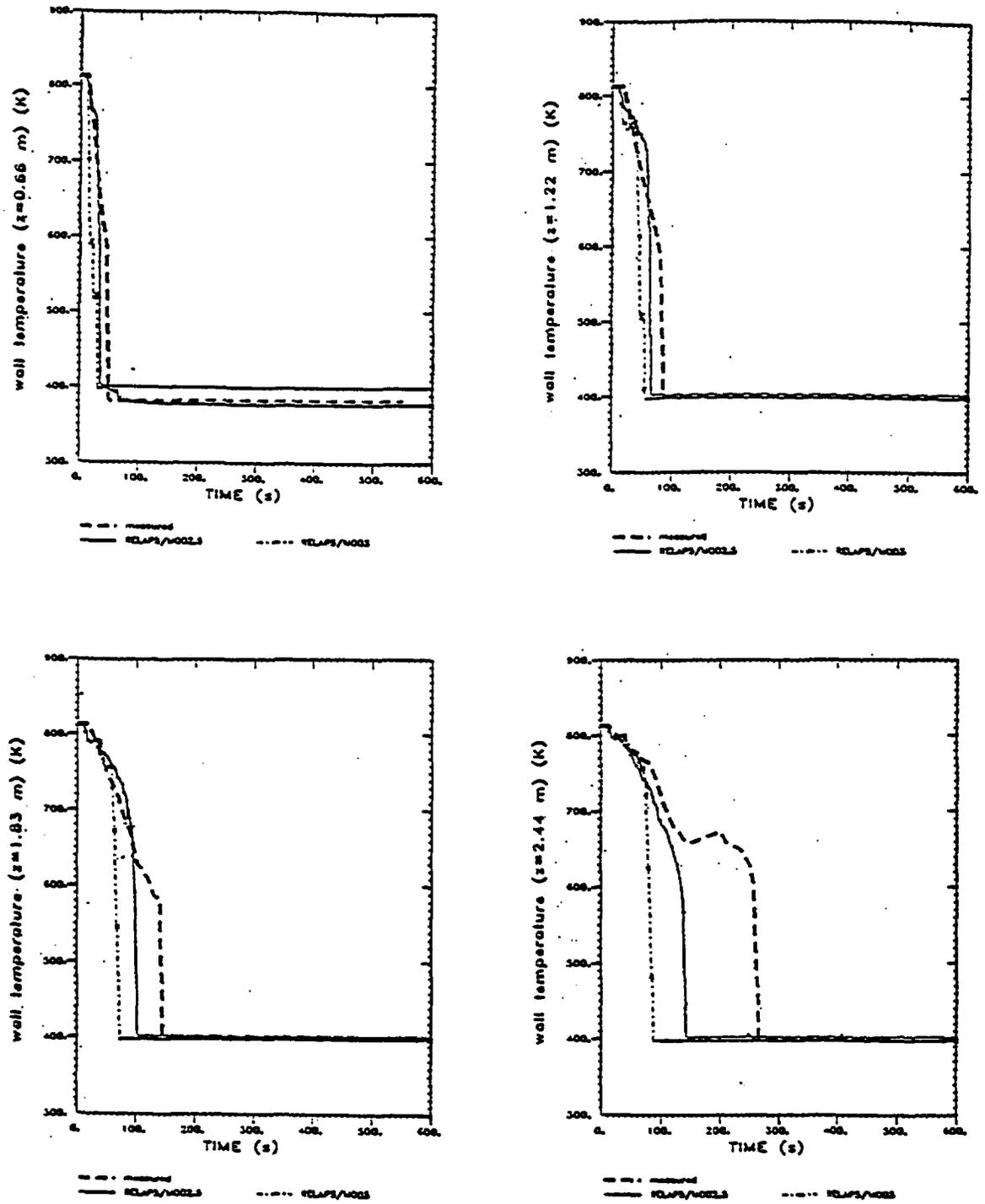


Figure 4: Comparison of the wall temperatures at 0.66, 1.22, 1.83 and 2.44 m calculated by RELAP5/MOD2.5 and RELAP5/MOD3 with the experimental ones for Test 3051.

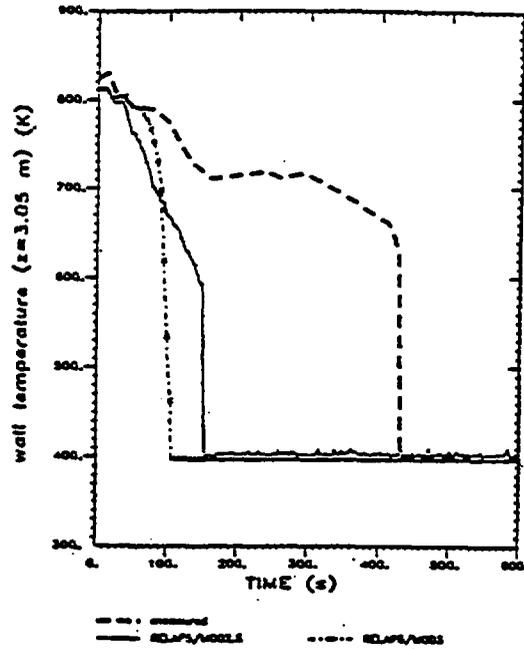
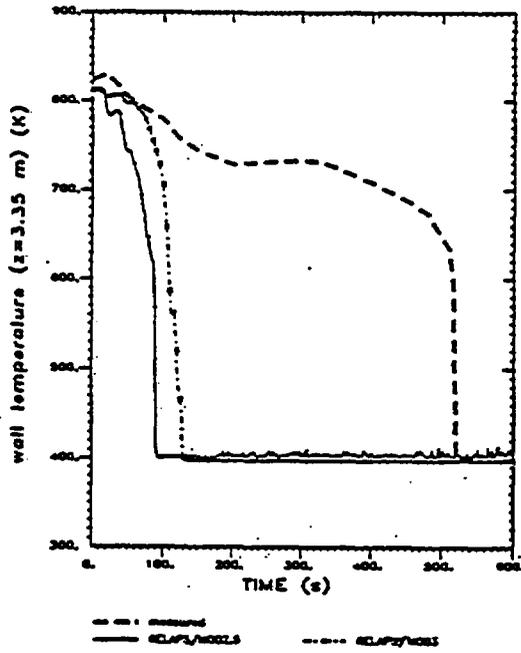


Figure 5: Comparison of the wall temperatures at 3.05 and 3.35 m calculated by RELAP5/MOD2.5 and RELAP5/MOD3 with the experimental ones for Test 3051.

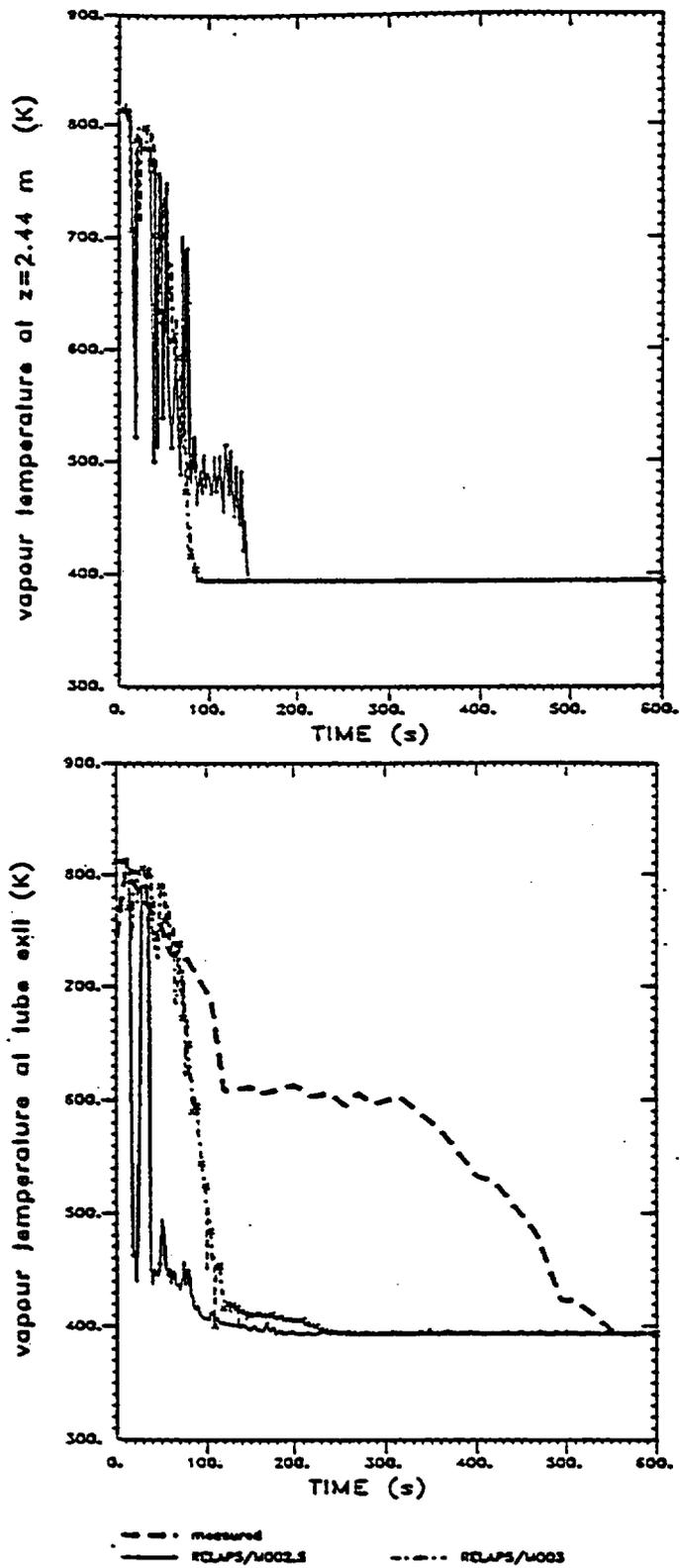


Figure 6: Calculated vapour temperature at 2.44 m and comparison of the calculated vapour temperature at the tube exit with the measured value.

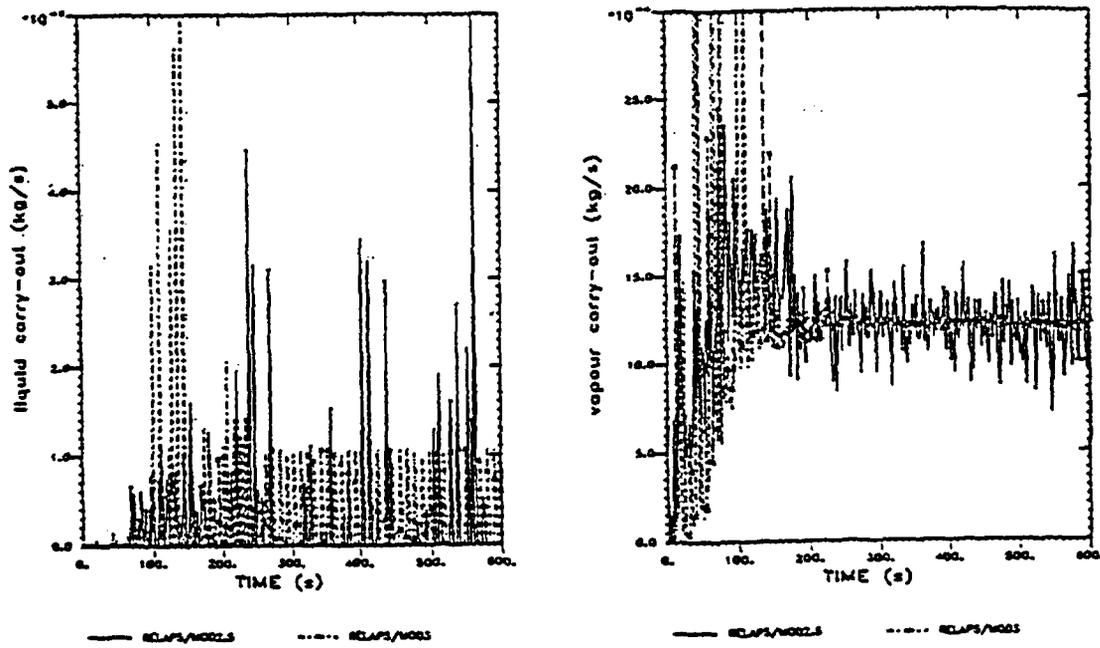


Figure 7: Calculated vapour and liquid carry-out

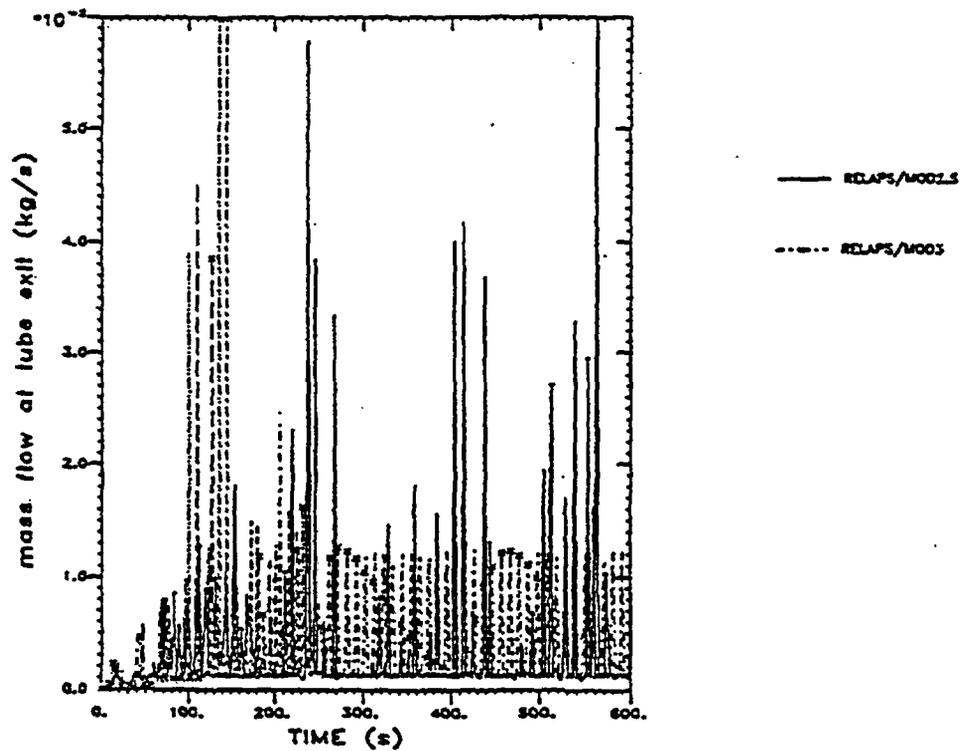


Figure 8: Calculated mass flow rate at the tube exit.

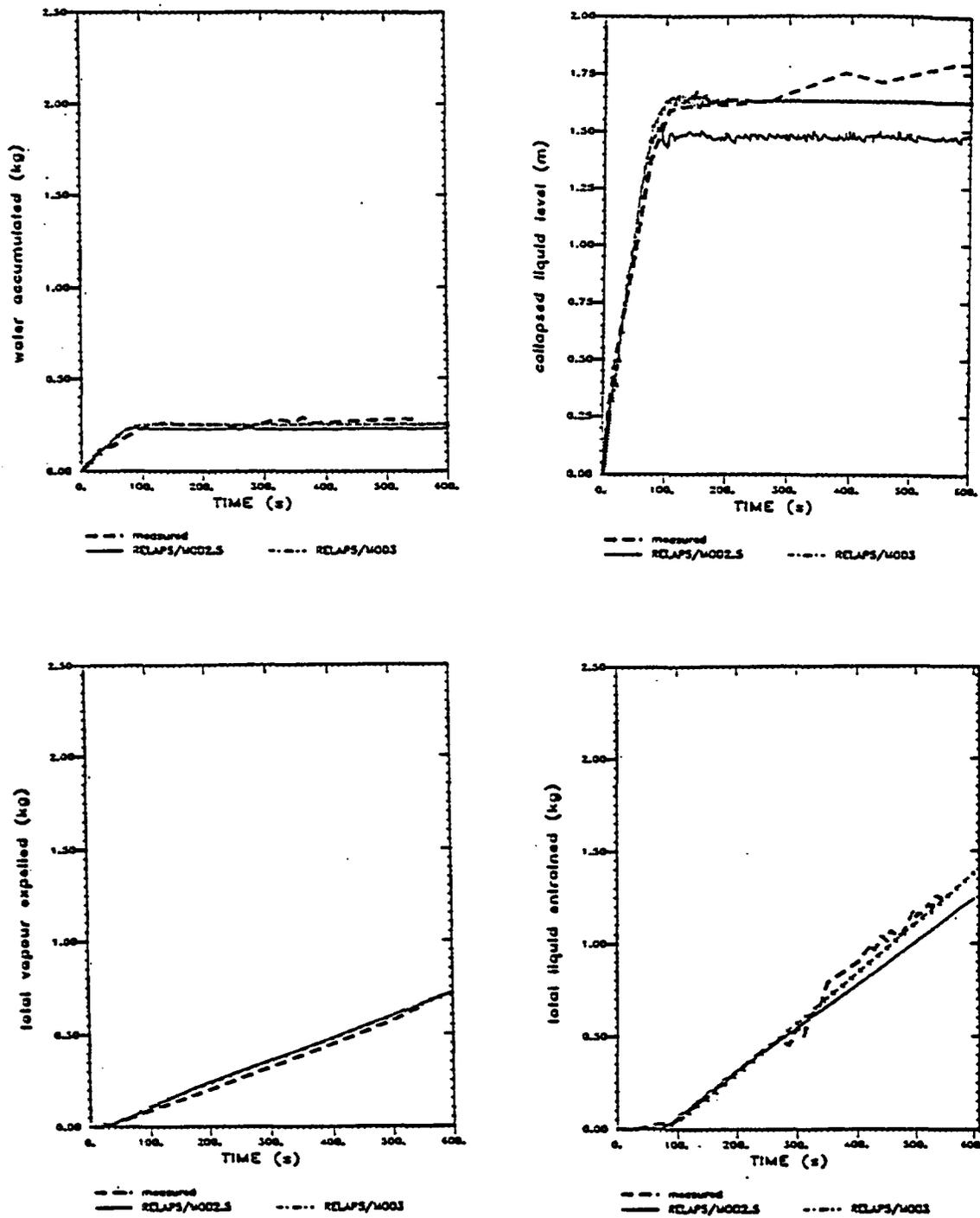


Figure 9: Comparison of the calculated accumulated mass, collapsed liquid level, vapour and liquid entrainment with the experimental data.

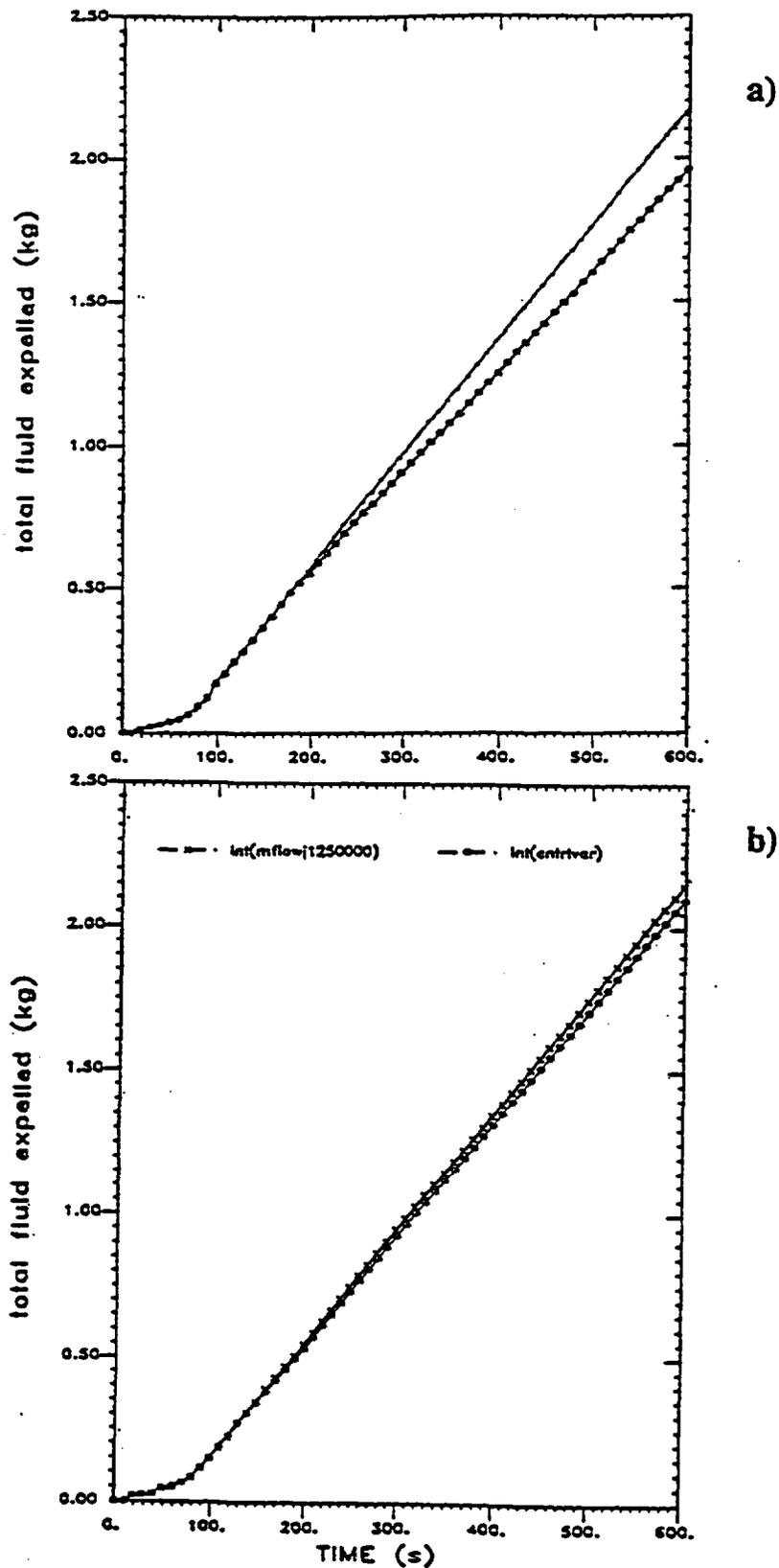


Figure 10: Comparison of the calculated fluid expelled out of the test section by integrating directly the mass flow with that using control variables and the primitive variables (velocities, void fractions and densities): a) RELAP5/MOD2.5 and b) RELAP5/MOD3.

Steady-state calculation:
axial temperature
profiles in the post-CHF
region 'shift' with the
position of the QF.

Conditions:

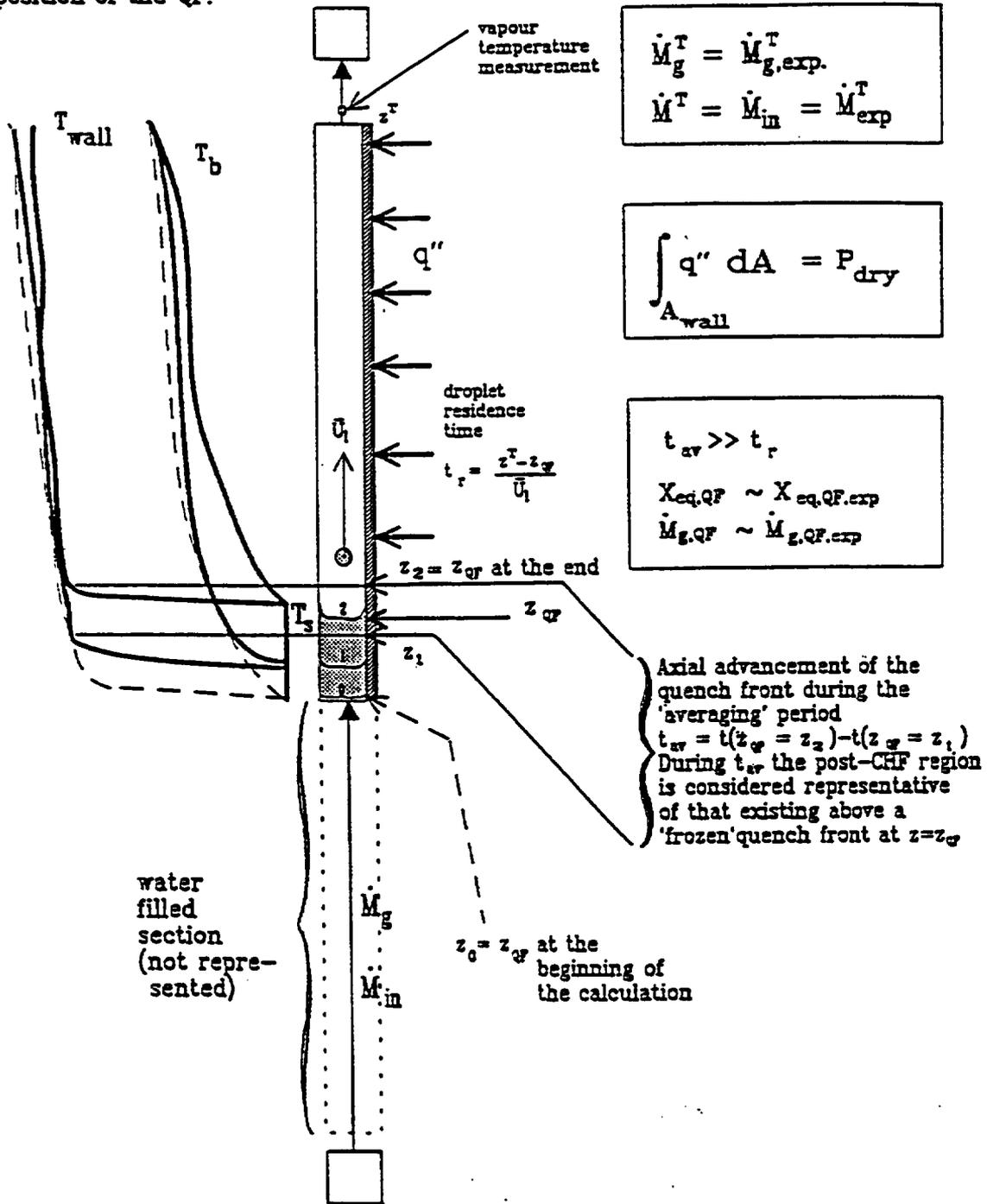


Figure 11: Prerequisites for the 'quasi steady-state' analysis of the post-CHF zone above a slowly advancing quench front (under these conditions the interfacial heat transfer model can be assessed)

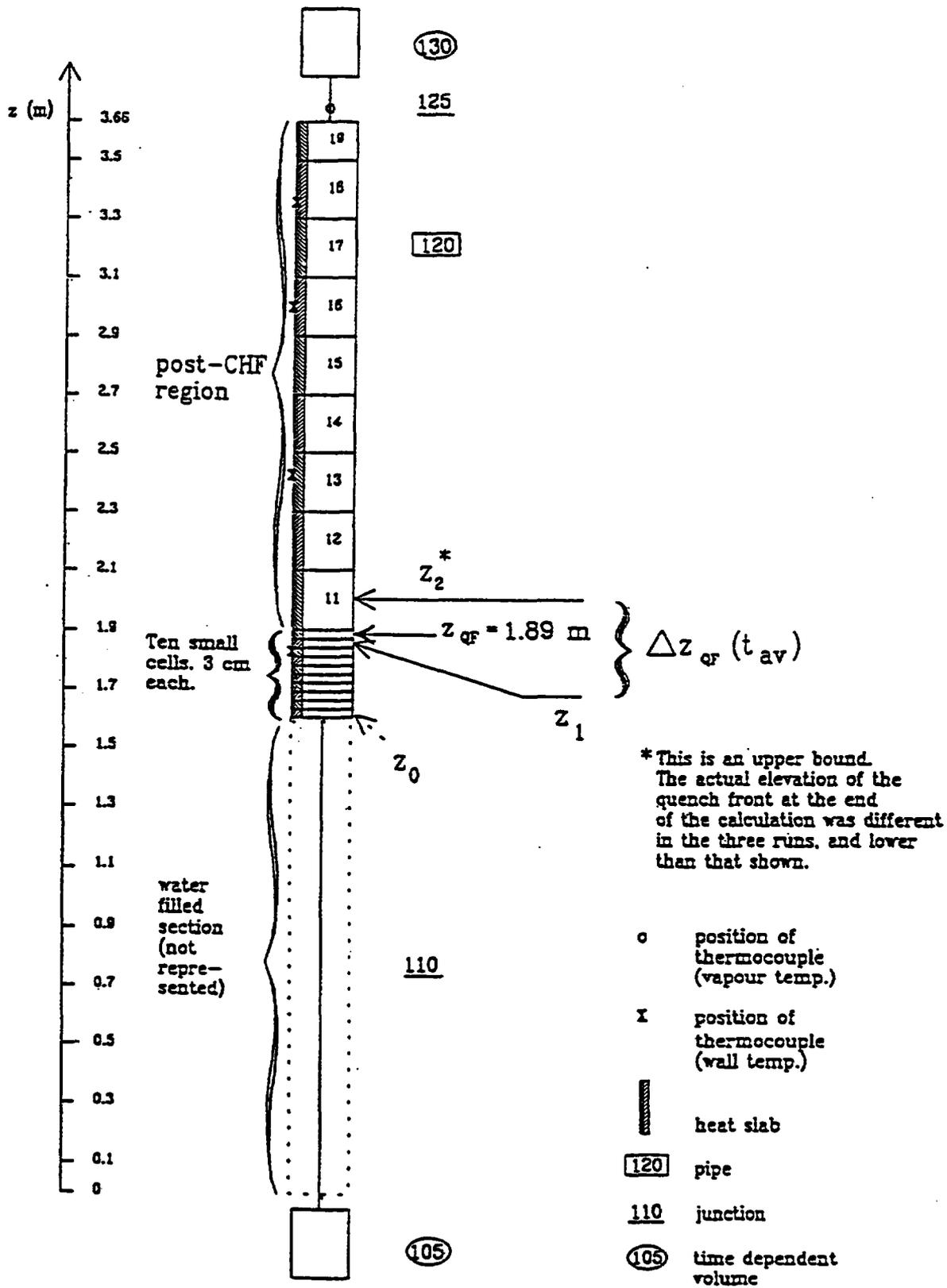


Figure 12: Nodalization used for the steady-state analysis of Run 3051 by RELAP5/MOD2.5.

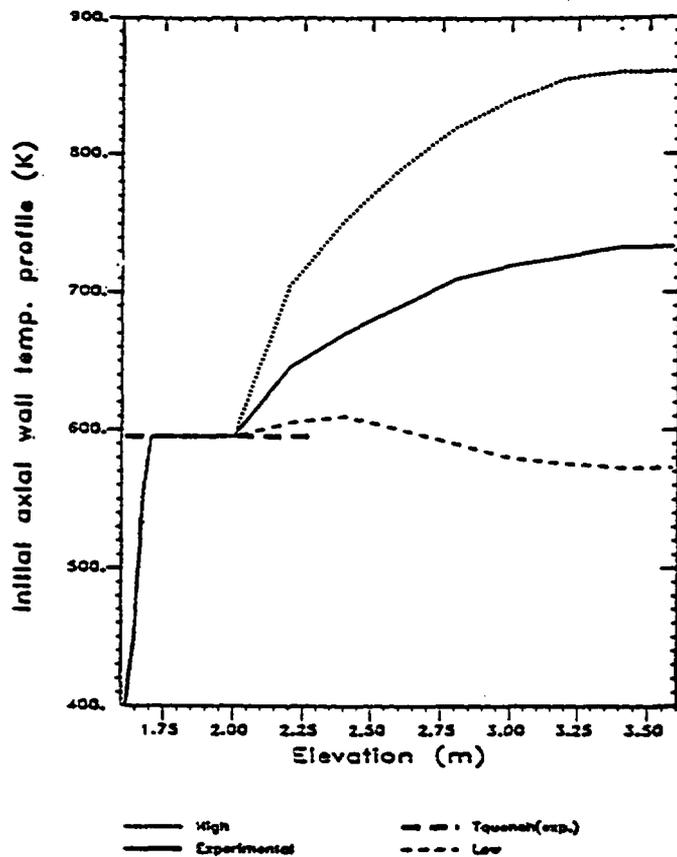


Figure 13: Initial axial wall temperature distributions used for the steady-state analysis of test 3051 at the time of the mid-height quench.

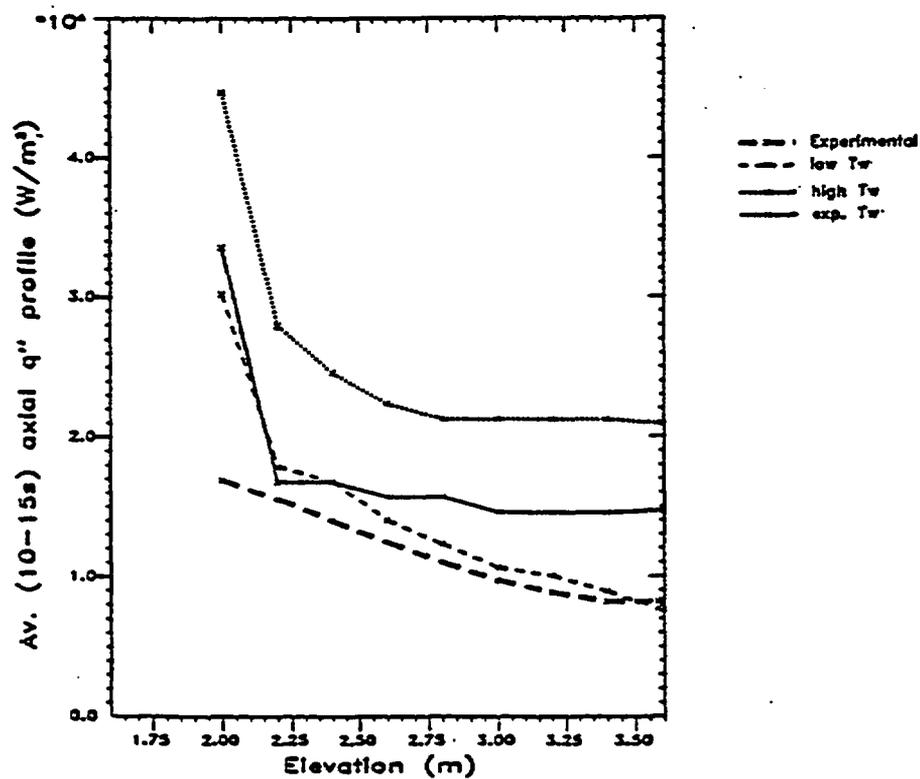


Figure 14: Average axial wall heat flux distributions during the time span 10-15 s for the three initial axial wall temperature distributions shown in Fig. 13.

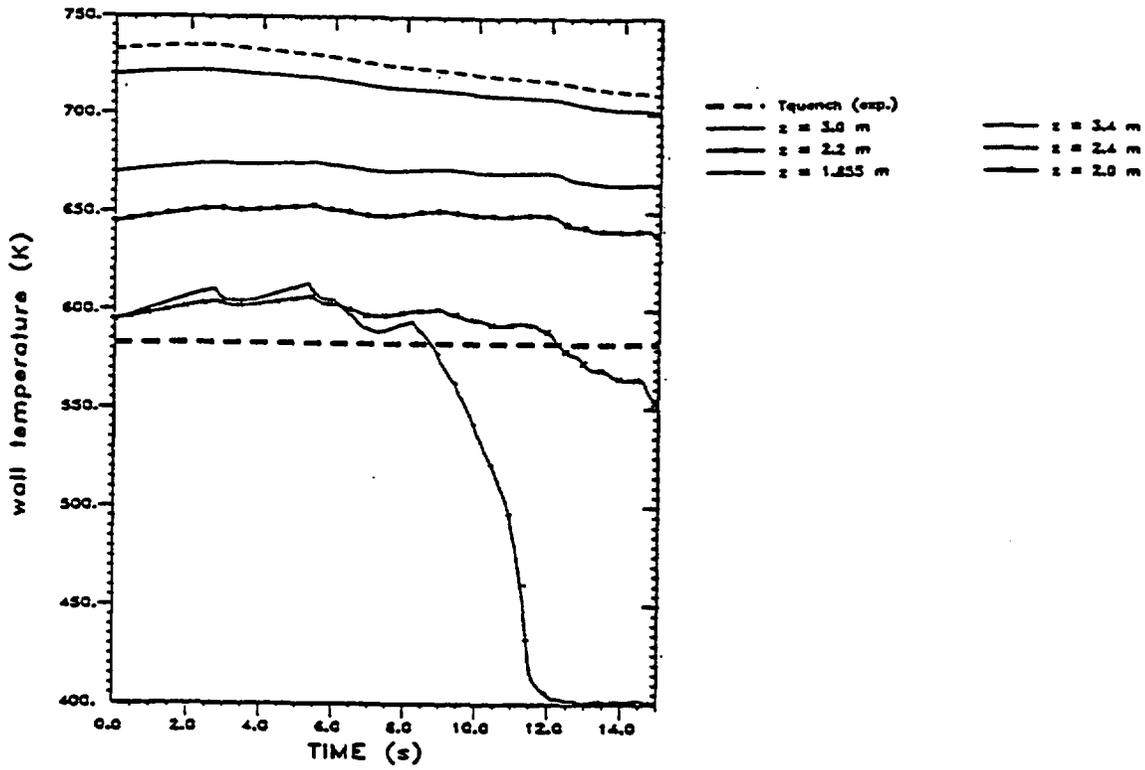


Figure 15: Calculated wall temperatures time-history at various elevations, when the experimental axial wall temperature distribution is imposed.

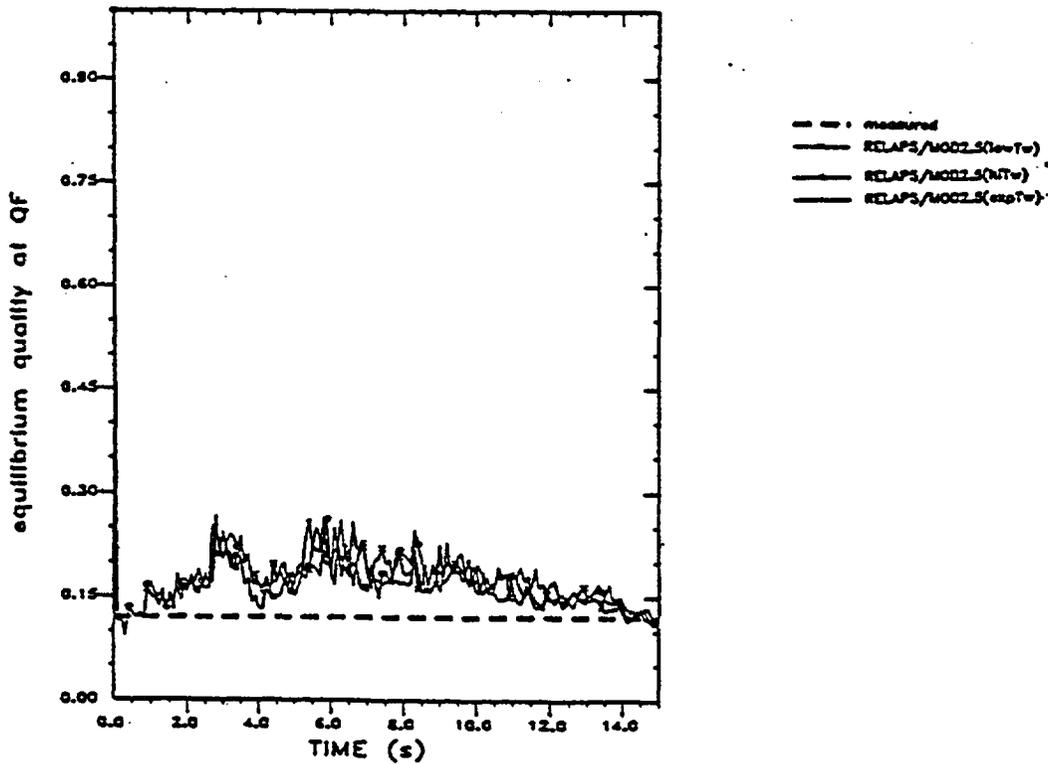


Figure 16: Equilibrium quality at the QF calculated using the three different initial wall temperature distributions (Fig. 13).

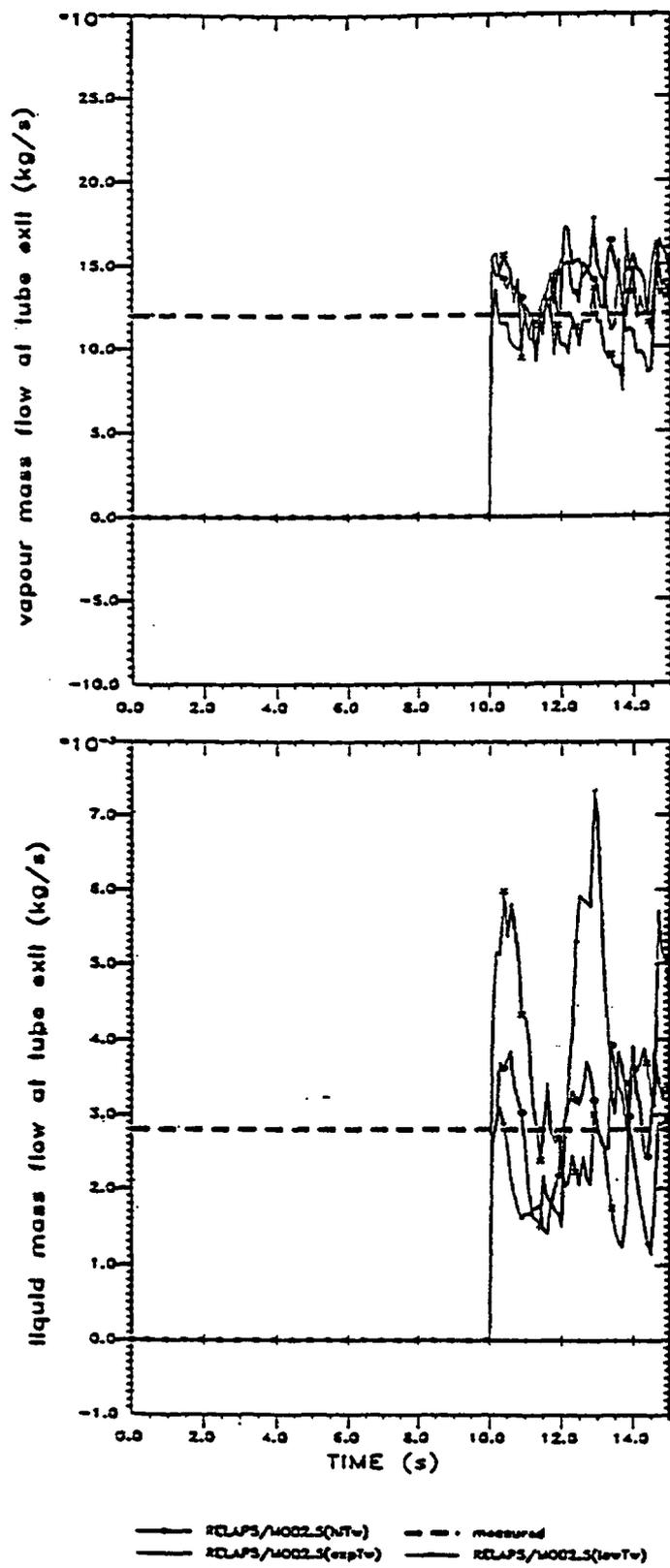


Figure 17: Vapour and liquid mass flow rates at the tube exit.

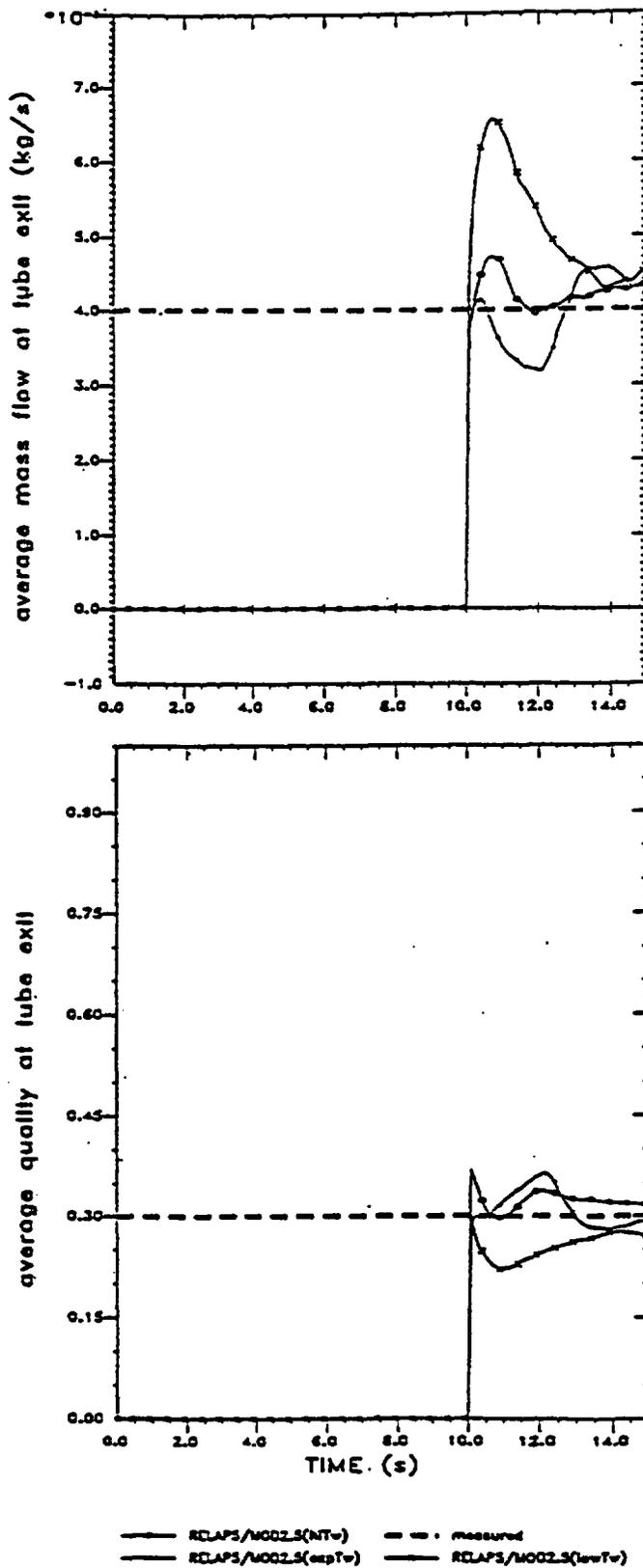


Figure 18: Average values of the total mass flow rates and exit quality over the time span 10-15 sec. The quantities to be averaged are calculated starting from 10 sec.

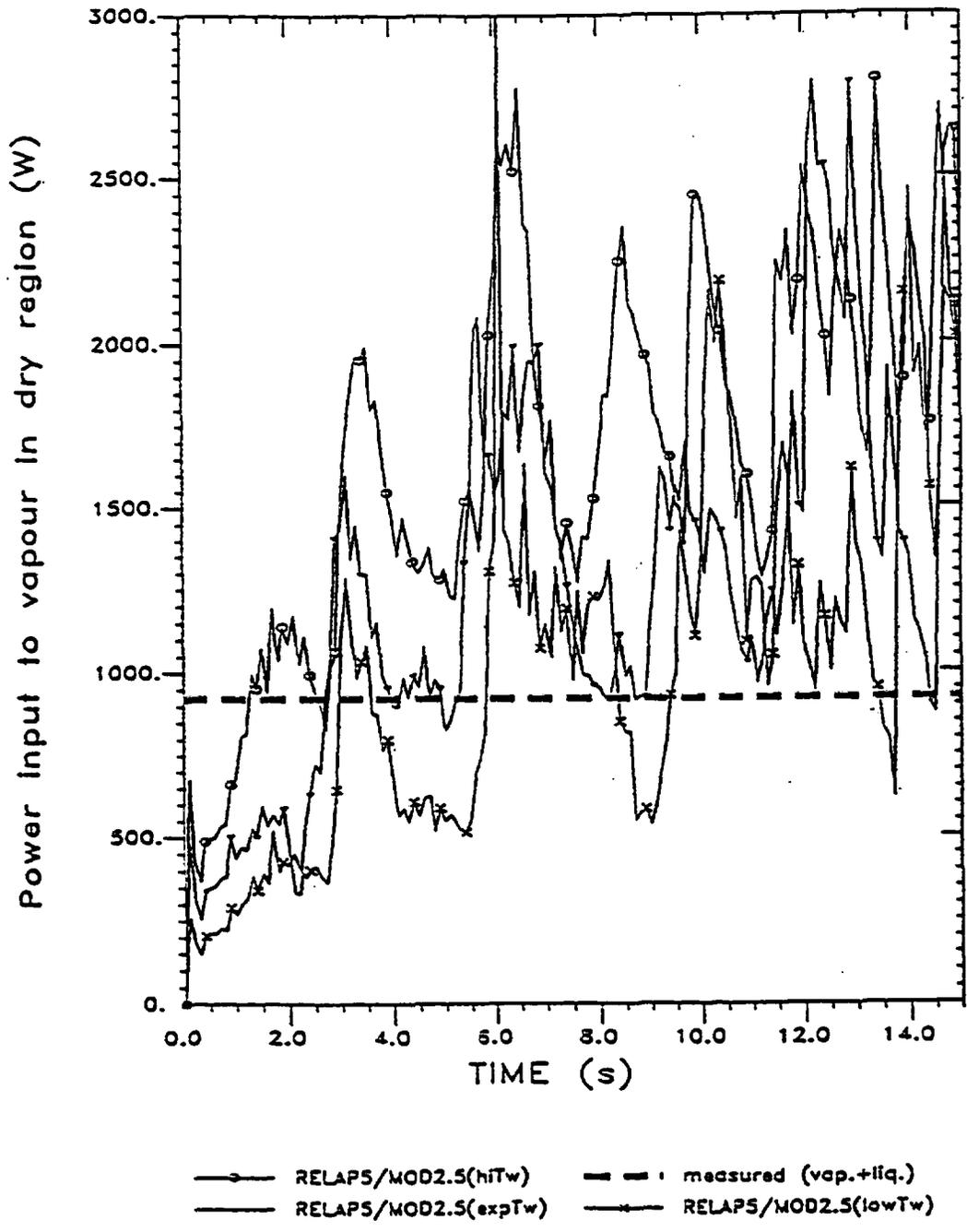


Figure 19: Comparison of the heat transfer rate to vapour with the total power input to the mixture measured in the experiment.

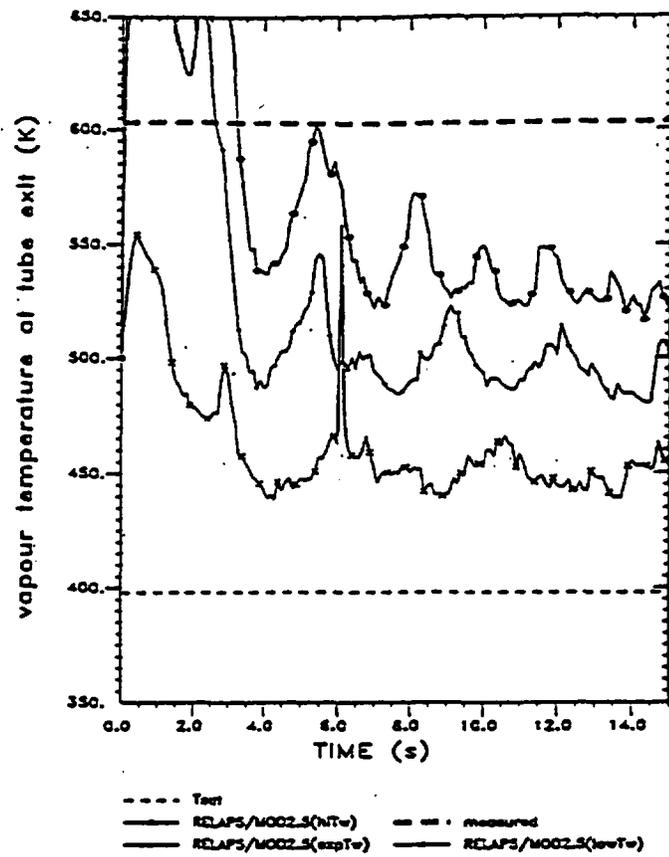


Figure 20: Comparison of the calculated vapour temperature at the tube exit with the experimental value at the time of the mid-height quench.

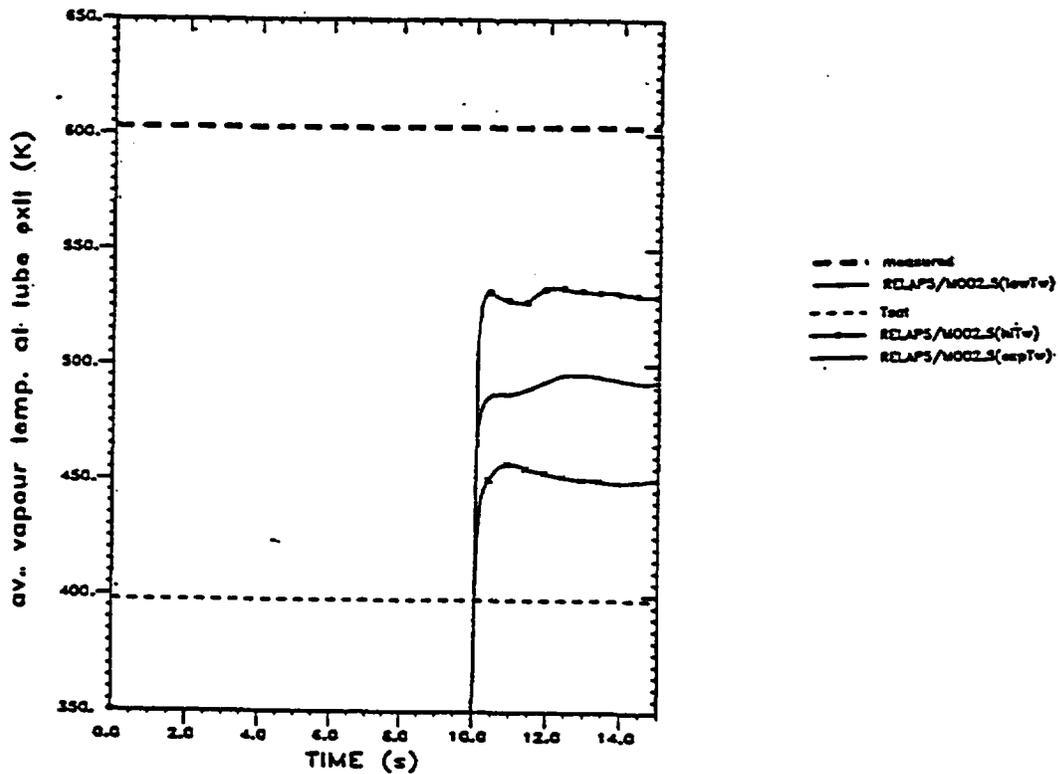


Figure 21: Time-average value of the vapour temperature at the tube exit.

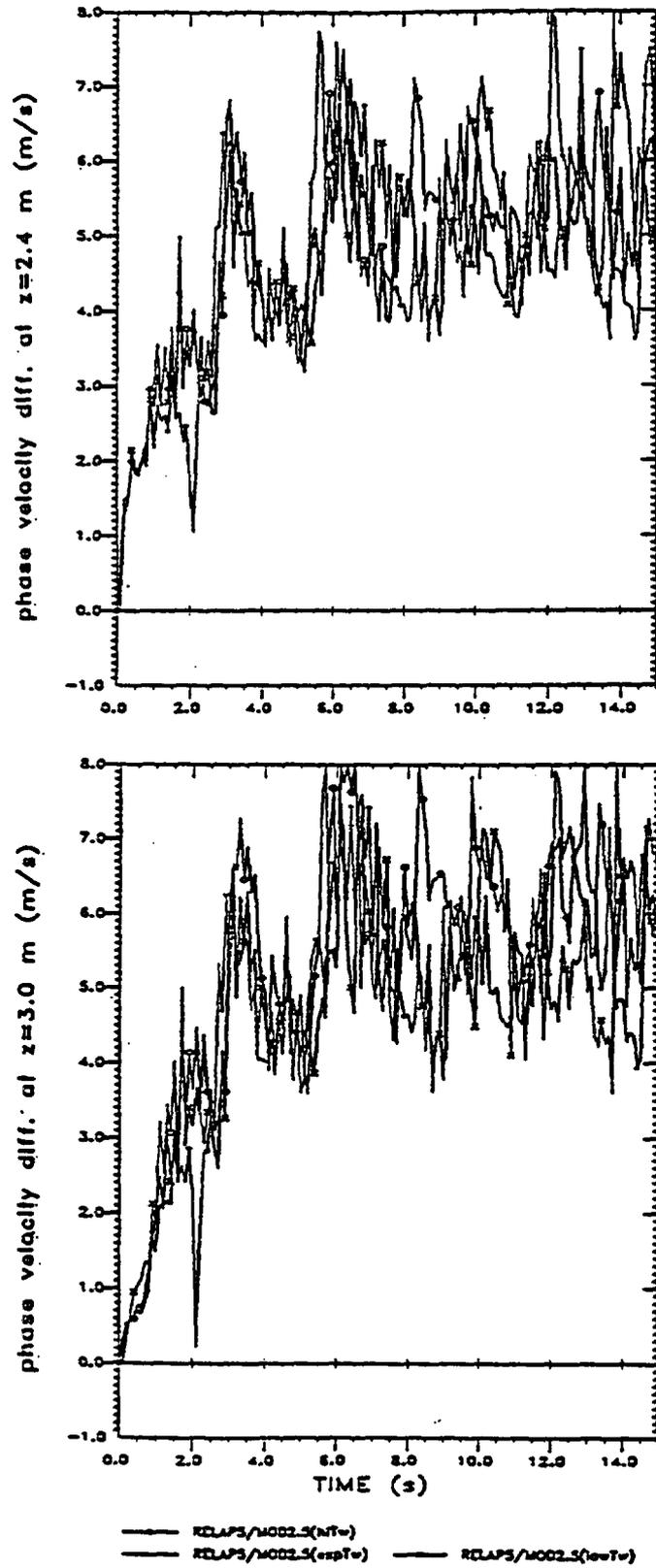


Figure 22: Phase velocity difference at $z=2.4$ m and $z=3.0$ m.

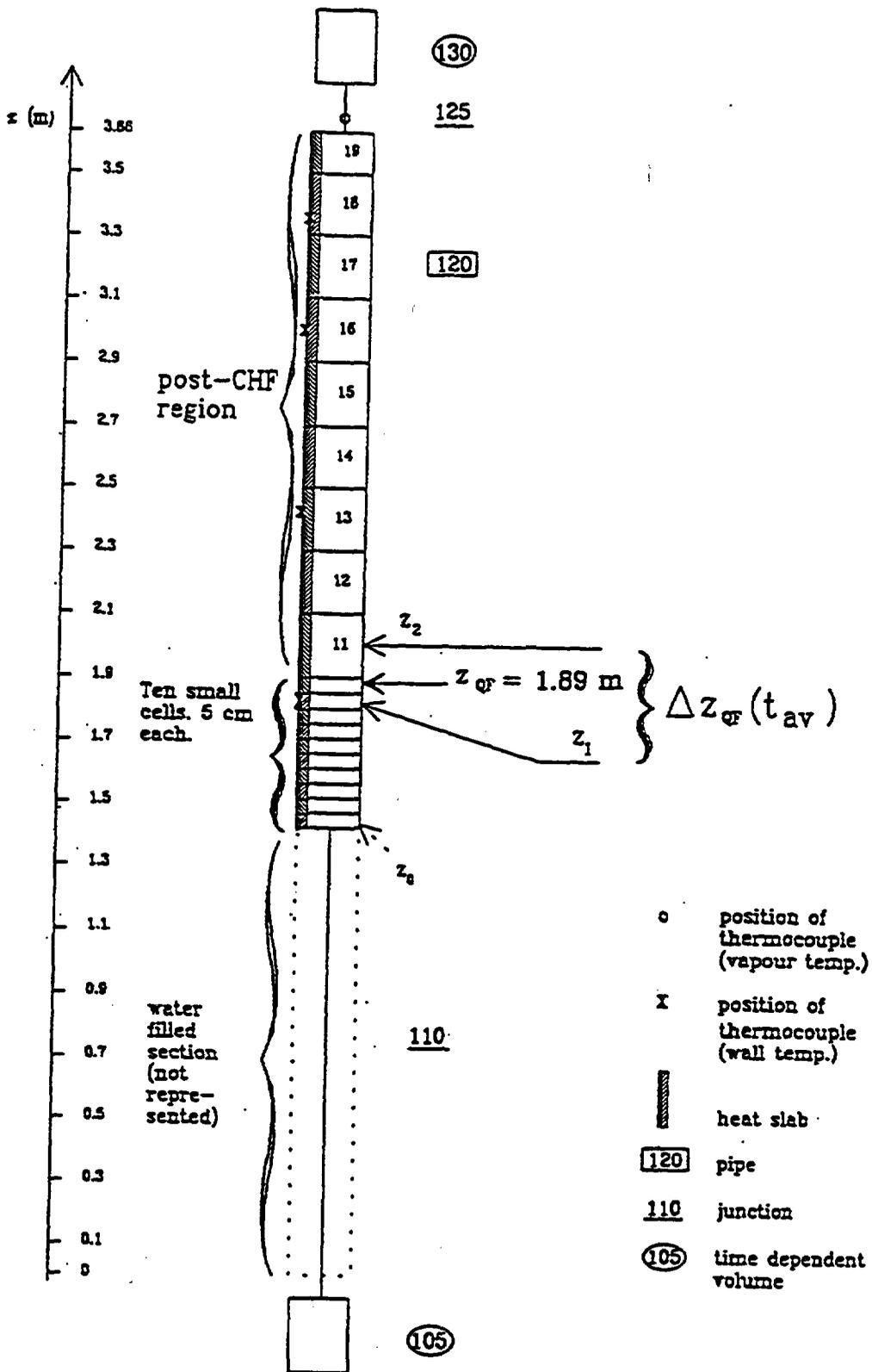


Figure 23: Nodalization used for the steady-state analysis of run 3051 by RELAPS/MOD3.

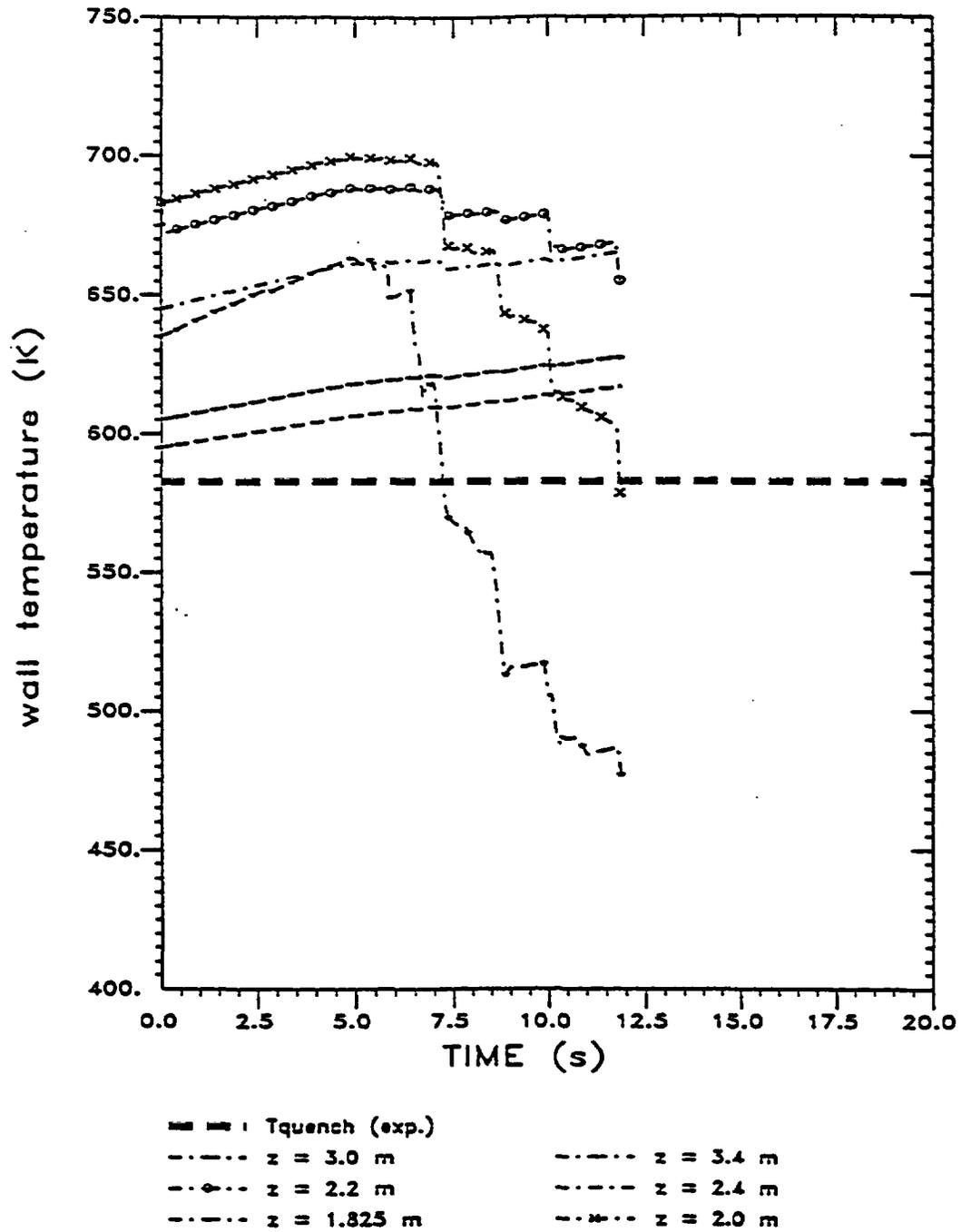


Figure 24: Wall temperatures at various elevations calculated by the frozen version of RELAP5/MOD3.

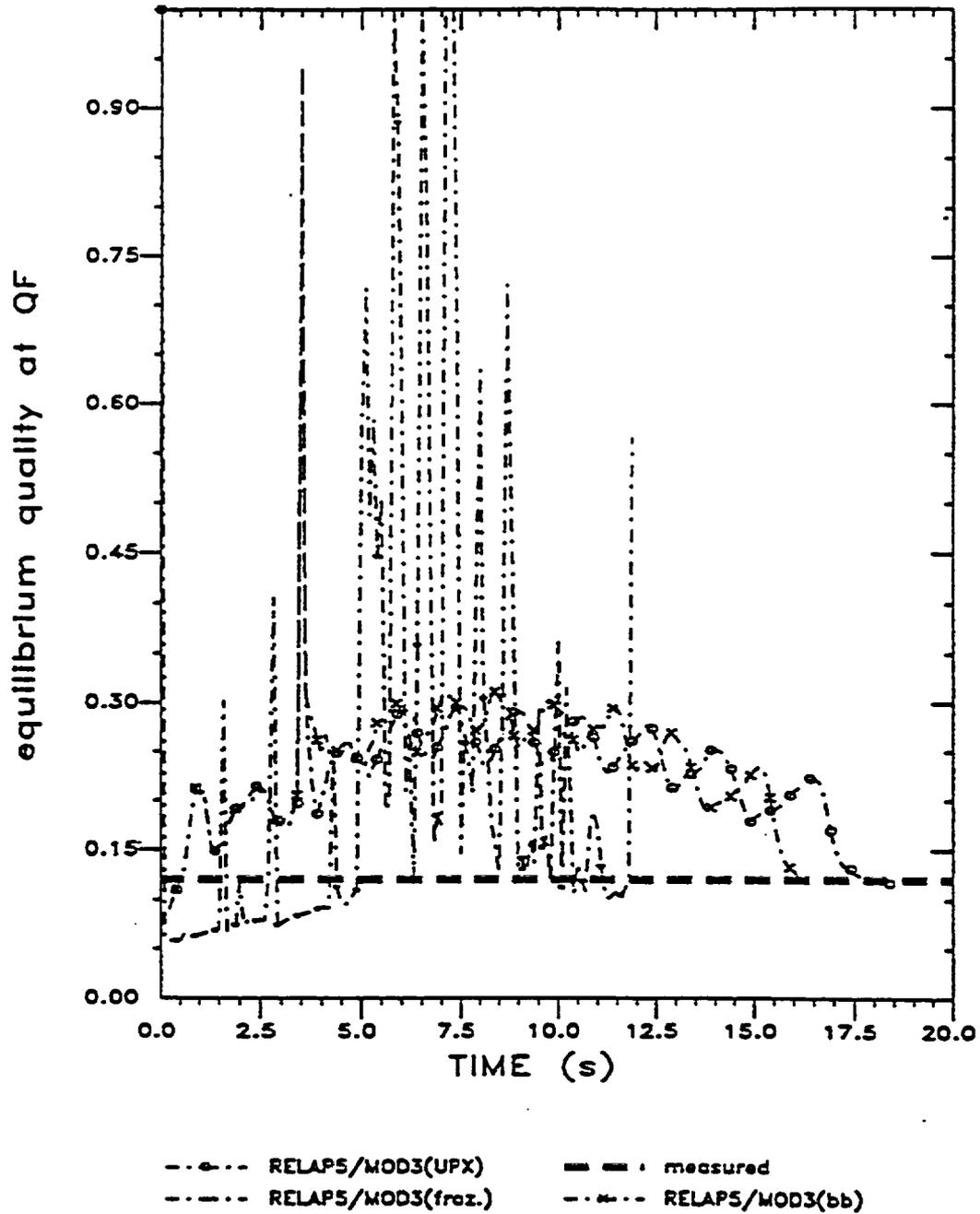


Figure 25: Equilibrium quality at the QF calculated by the frozen RELAP5/MOD3 (· · · · ·), by a version (bb) including the PSI updates (- · × · - : Analytis, 1992) and an 'ad hoc' modification (version UPX) of the RELAP5/MOD3(bb) using a reduced value for the minimum droplet diameter and the correct drag coefficient for droplets (- · o · -).

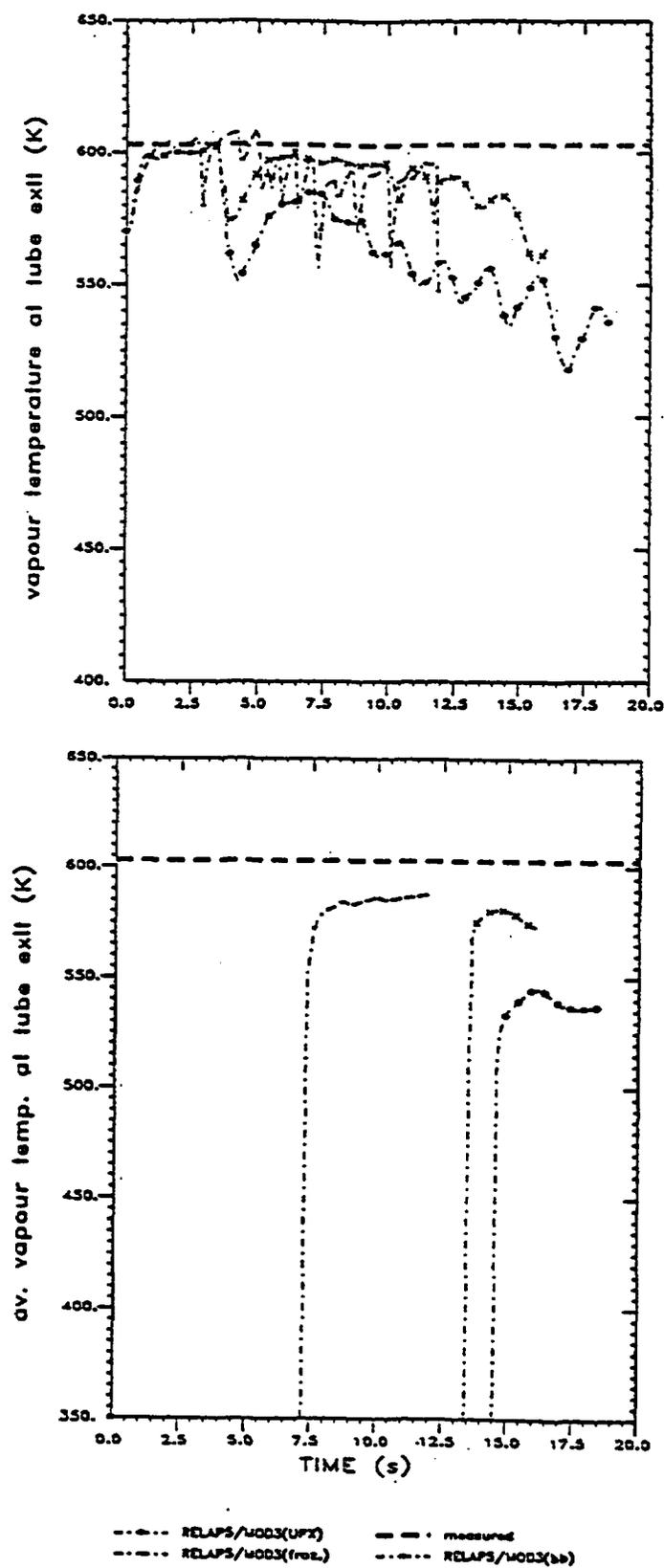


Figure 26: Vapour temperature at the tube exit calculated by three versions of RELAP5/MOD3 (see legend in Fig. 25).

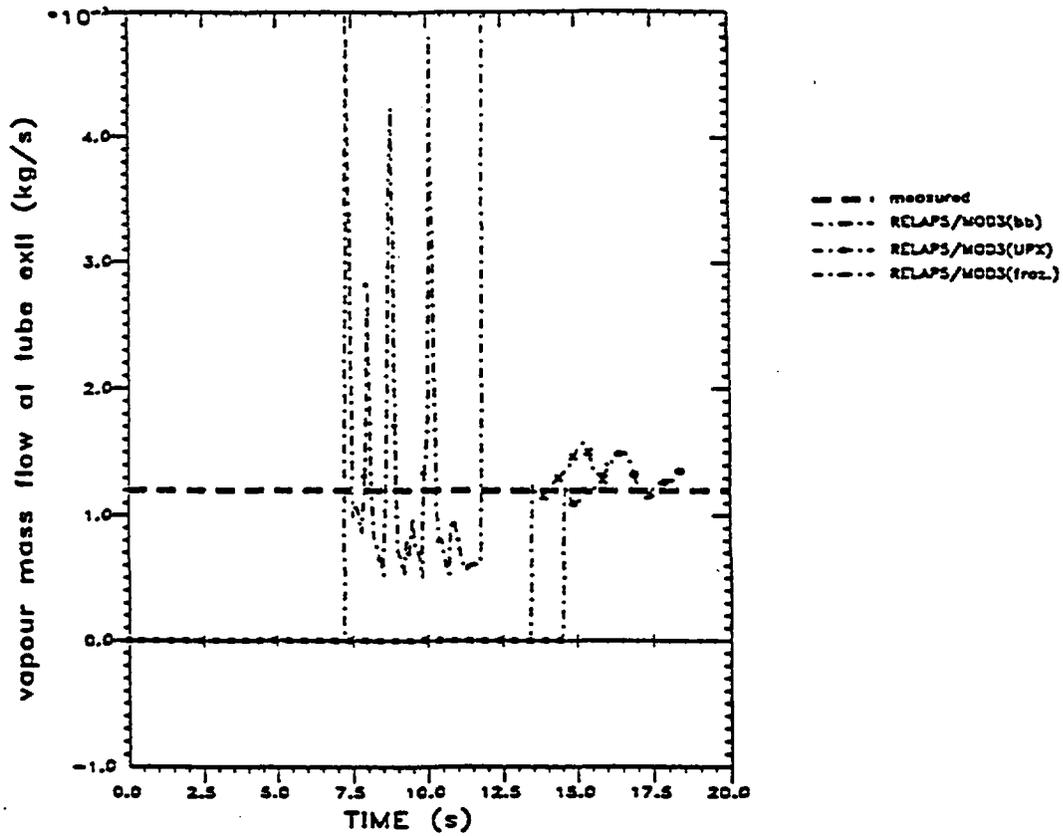


Figure 27: Vapour carry-out calculated by three versions of RELAPS/MOD3 (see legend in Fig. 25).

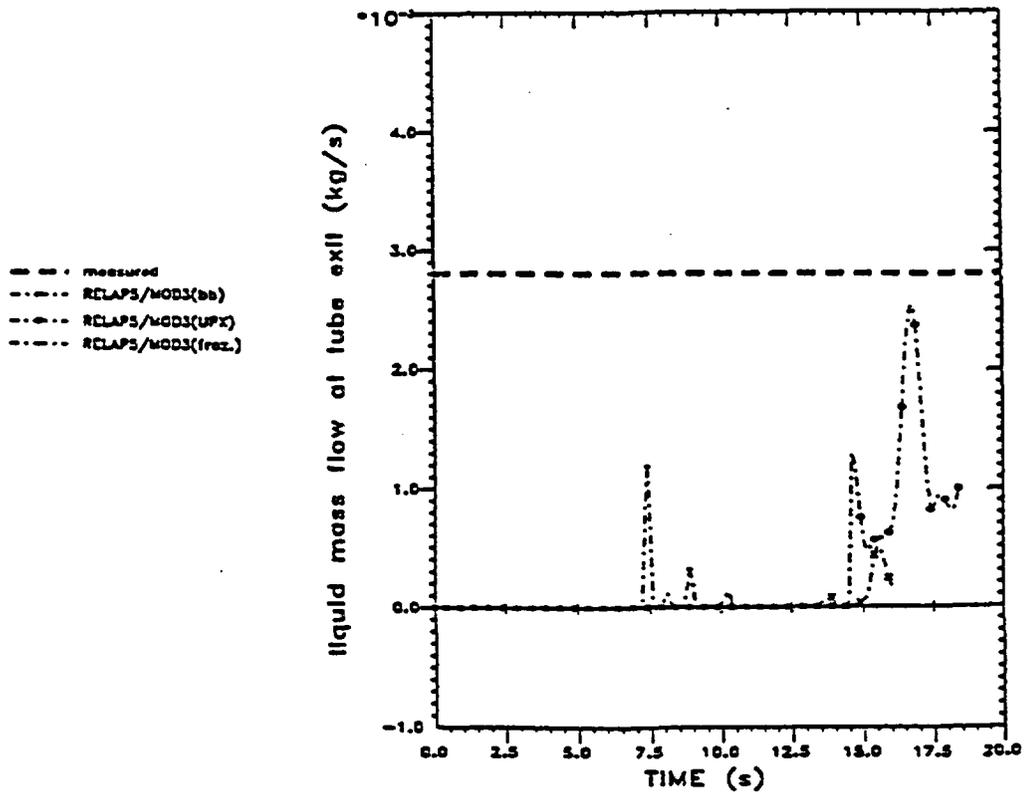
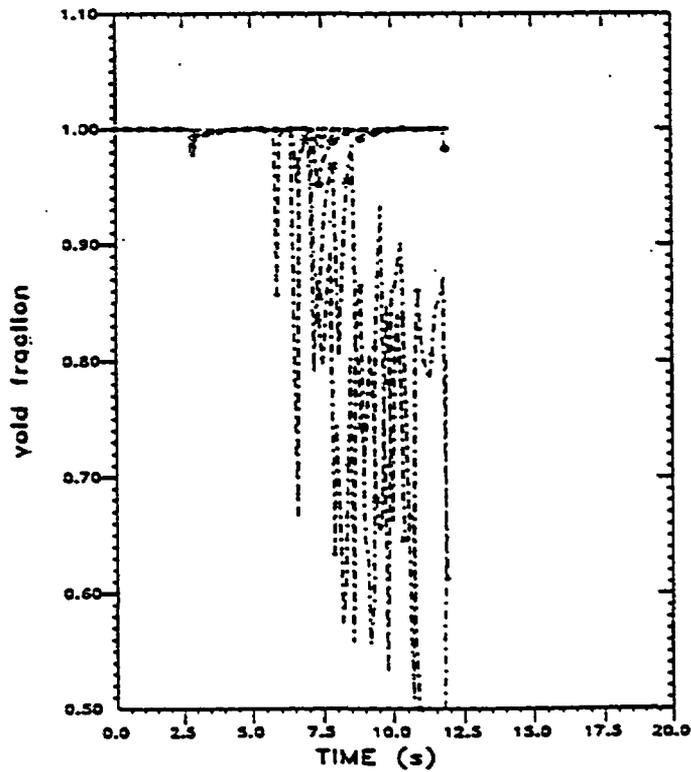
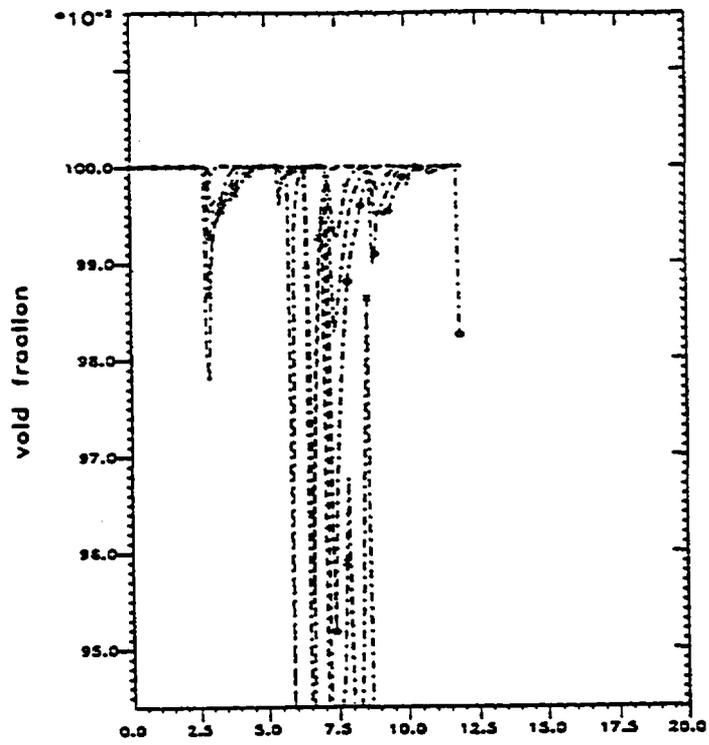


Figure 28: Liquid carry-out calculated by three versions of RELAPS/MOD3 (see legend in Fig. 25).



- - - - - 1.0
 - - - - - 1.1
 - - - - - 1.2
 - - - - - 1.3
 - - - - - 1.4
 - - - - - 1.5
 - - - - - 1.6
 - - - - - 1.7
 - - - - - 1.8
 - - - - - 1.9
 - - - - - 2.0
 - - - - - 2.1
 - - - - - 2.2
 - - - - - 2.3
 - - - - - 2.4
 - - - - - 2.5
 - - - - - 2.6
 - - - - - 2.7
 - - - - - 2.8
 - - - - - 2.9
 - - - - - 3.0

Figure 29: Void fractions (shown on two different scales) at various elevations calculated by the frozen version of RELAP5/MOD3.

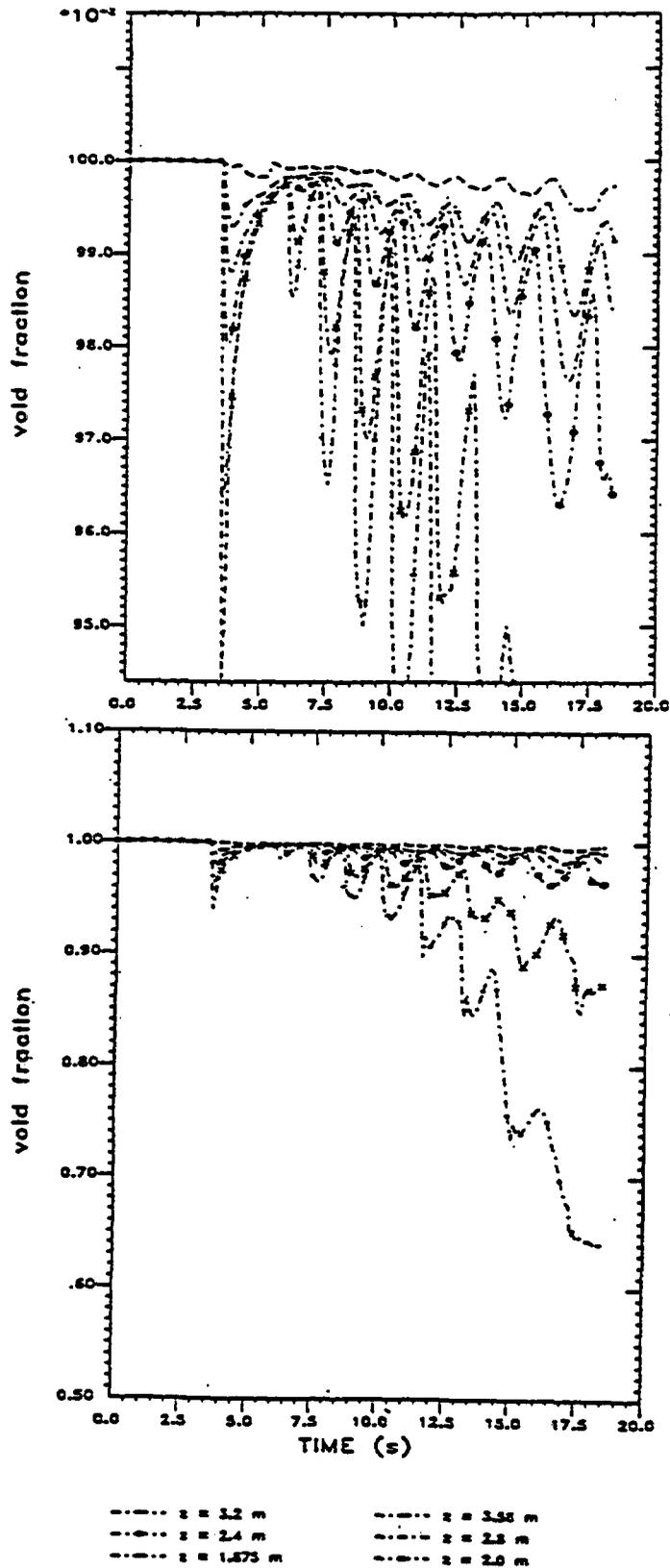


Figure 30: Void fractions (shown on two different scales) at various elevations calculated by RE-LAP5/MOD3+PSI updates (version bb).

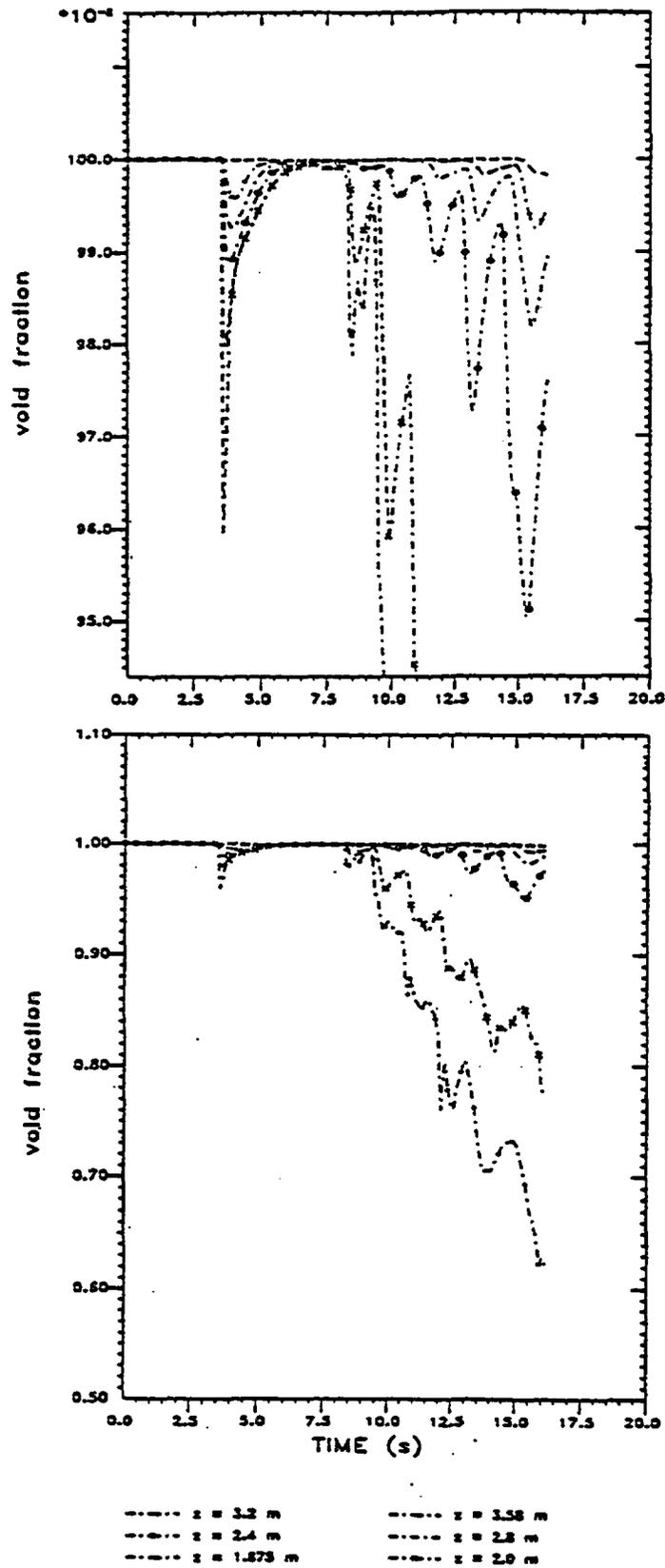


Figure 31: Void fractions (shown on two different scales) at various elevations calculated by RE-LAP5/MOD3+PSI updates+modified C_D and $d_{min} = 1.25$ mm (version UPX).

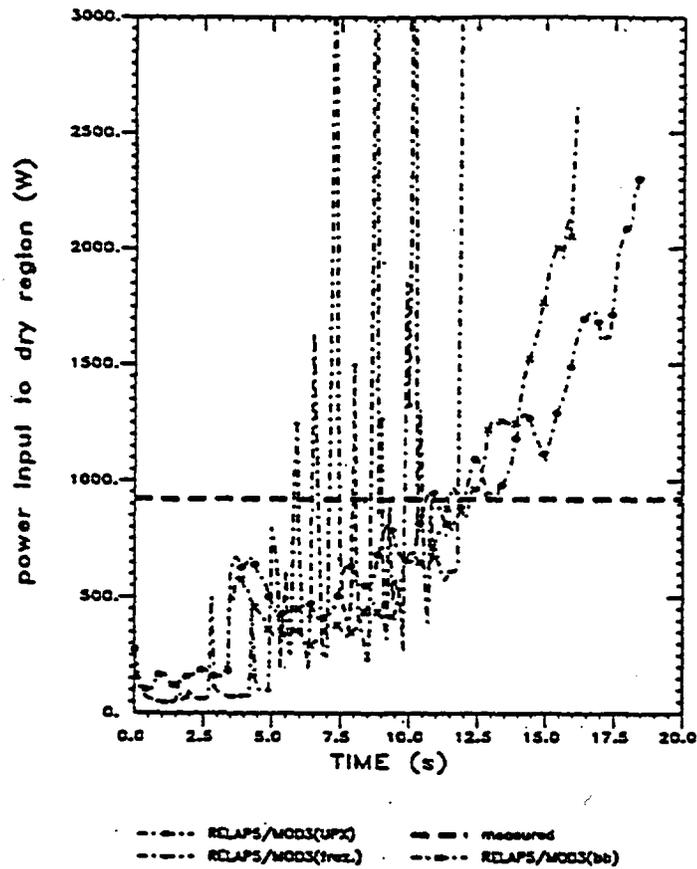


Figure 32: Power input to the mixture in the region above z_{QF} for the three versions of RELAP5/MOD3 (see legend in Fig. 25).

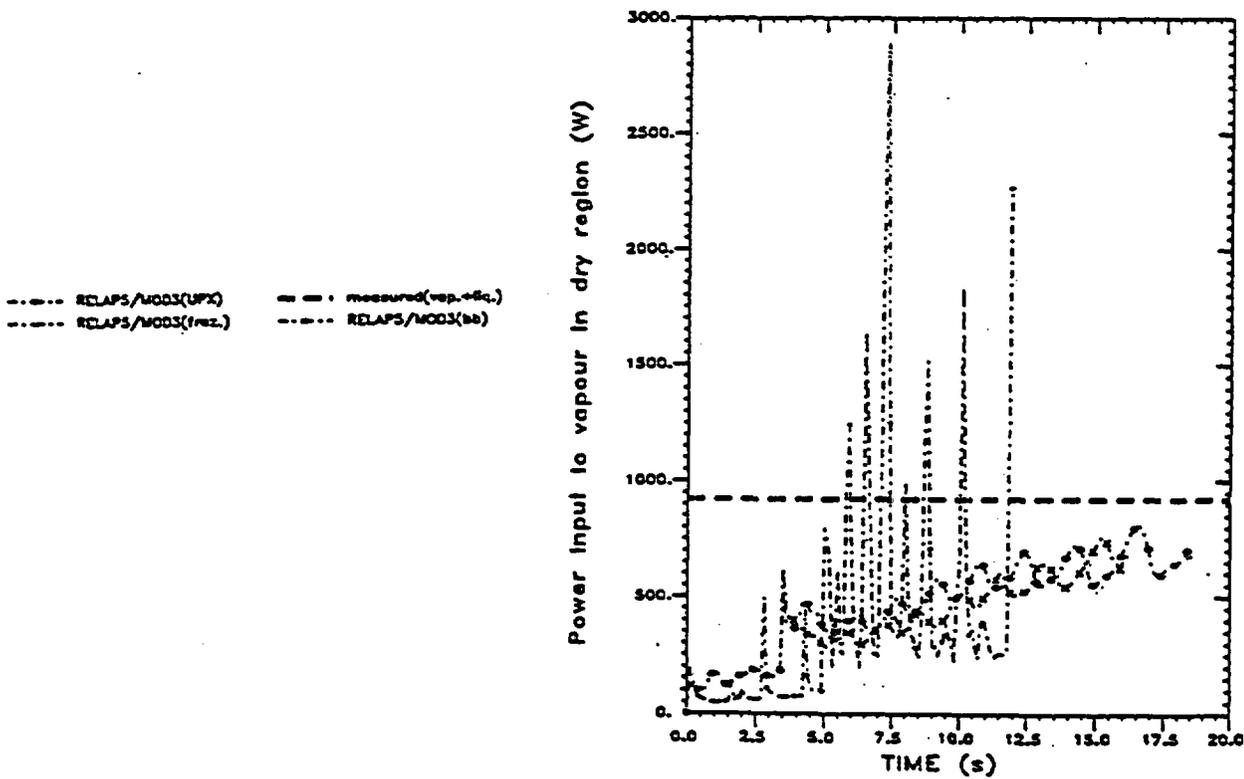


Figure 33: Heat transfer rate to the vapour alone for the three versions of RELAP5/MOD3 (legend in Fig. 25).

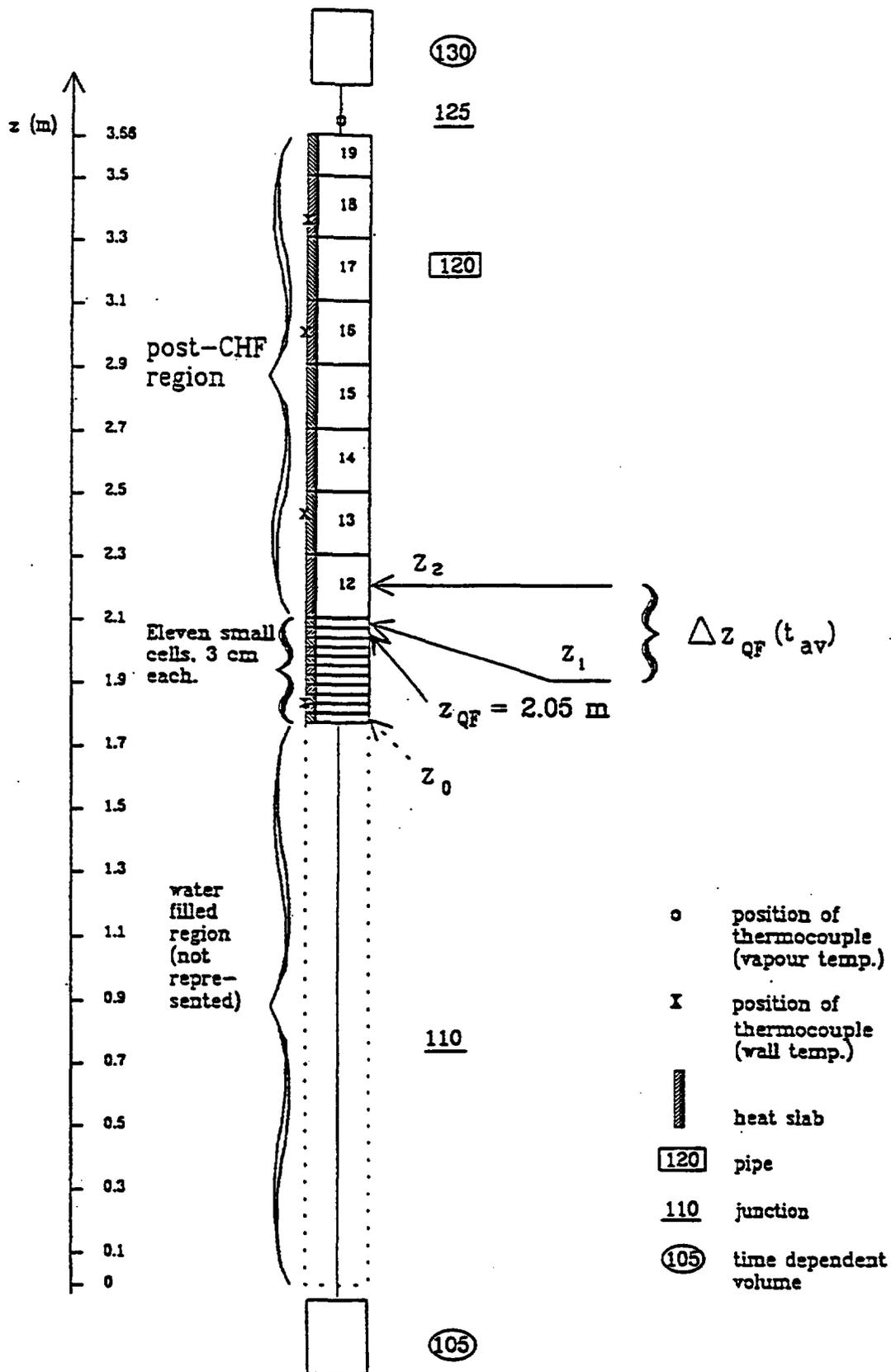


Figure 34: Nodalization used for the steady-state analysis of run 3053 by RELAP5/MOD2.5.

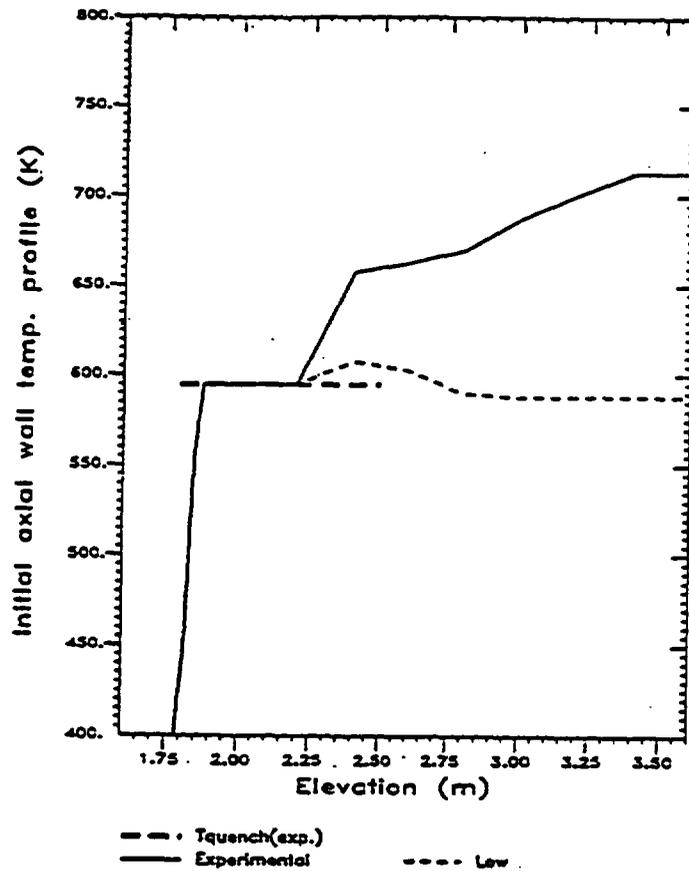


Figure 35: Initial axial wall temperature distributions used for the steady-state analysis of test 3053 at the time of the mid-height quench.

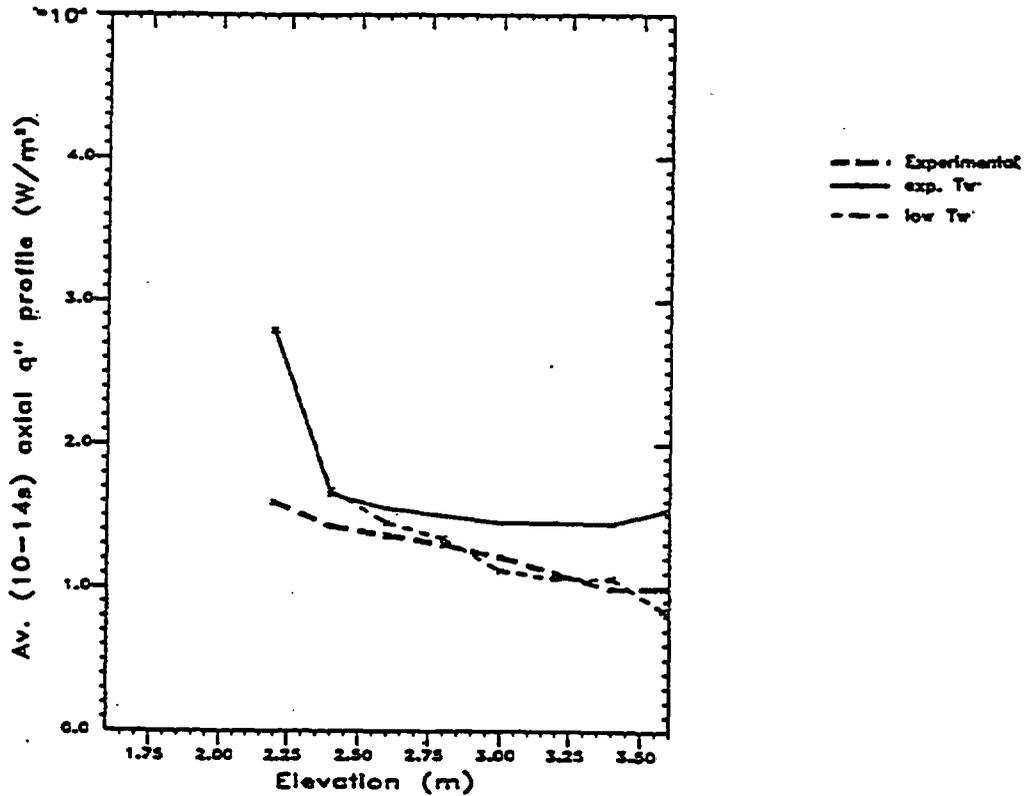


Figure 36: Average axial wall heat flux distributions during the time span 10-14 sec. for the two initial axial wall temperature distributions shown in Fig. 35.

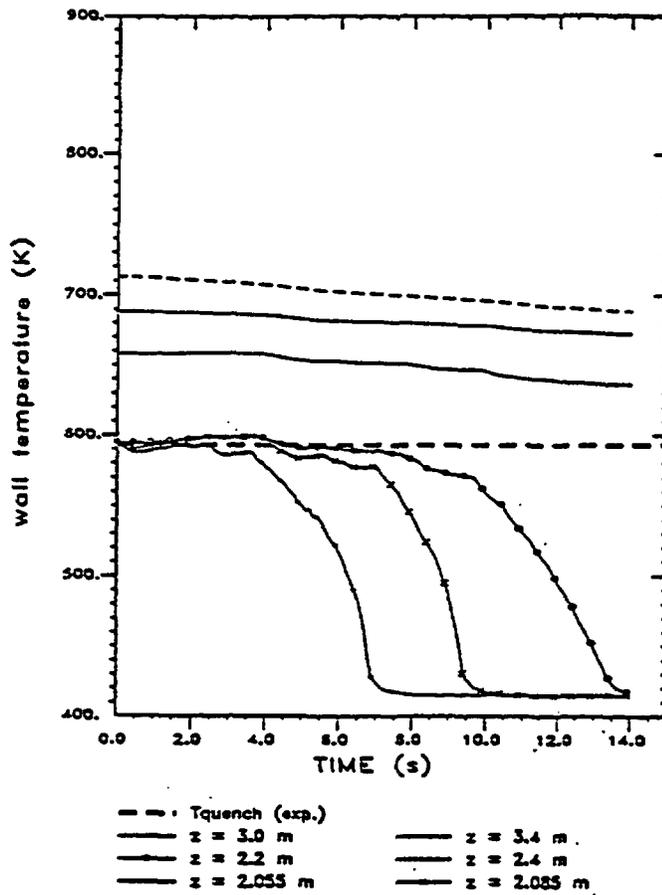


Figure 37: Wall temperatures at various elevations for test 3053 calculated by RELAP5/MOD2.5.

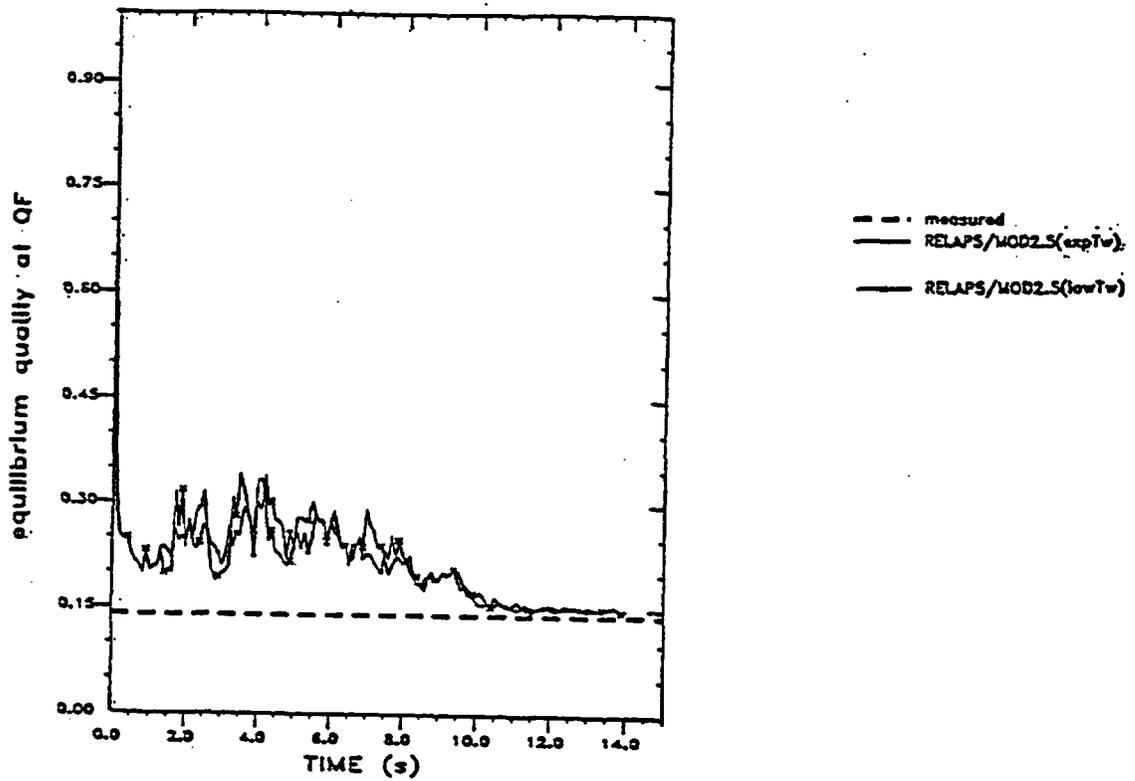


Figure 38: Equilibrium quality at the QF for the two different initial axial wall temperature profiles.

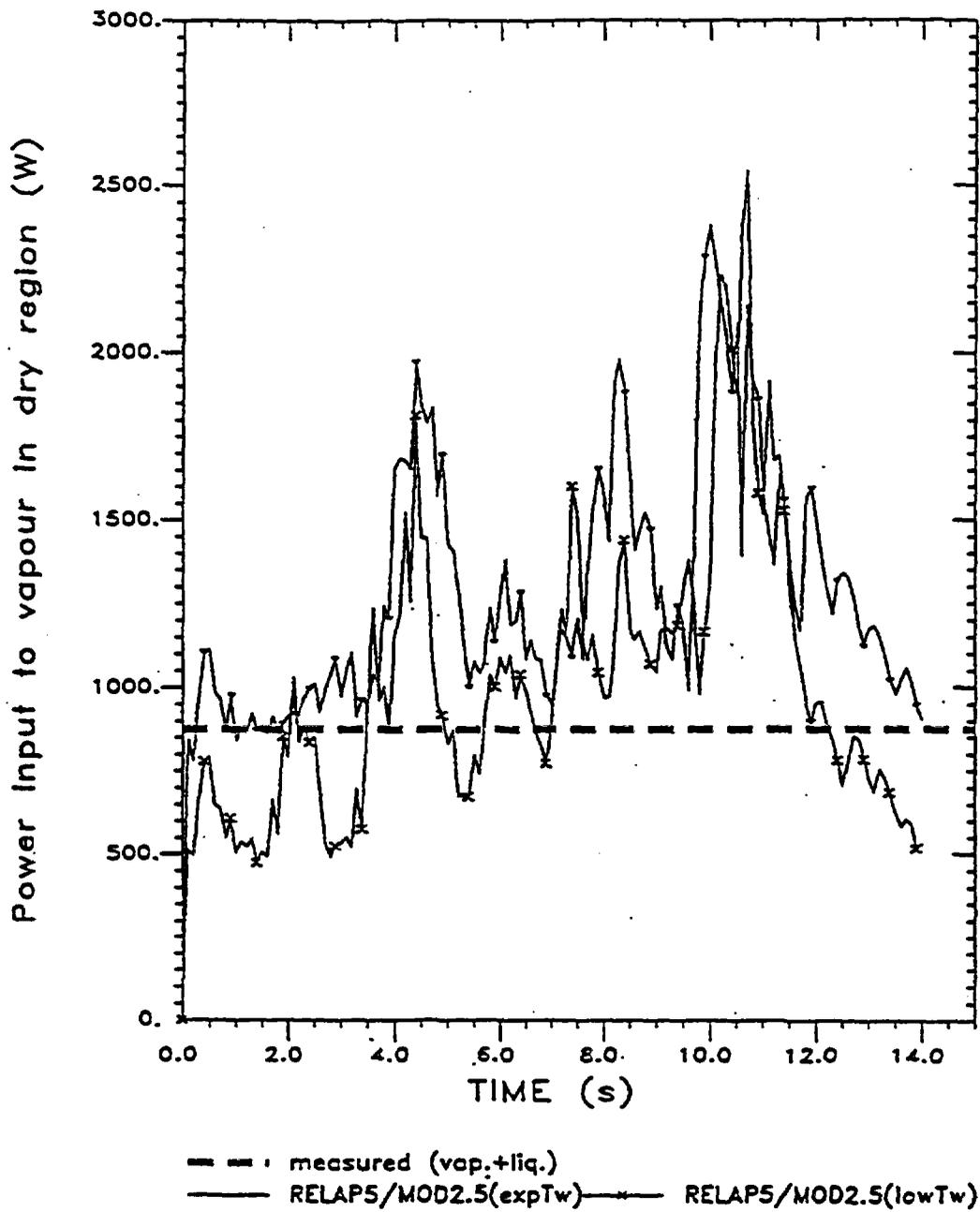


Figure 39: Heat transfer rate from wall to vapour in the post-dryout region for the two initial wall temperature distributions.

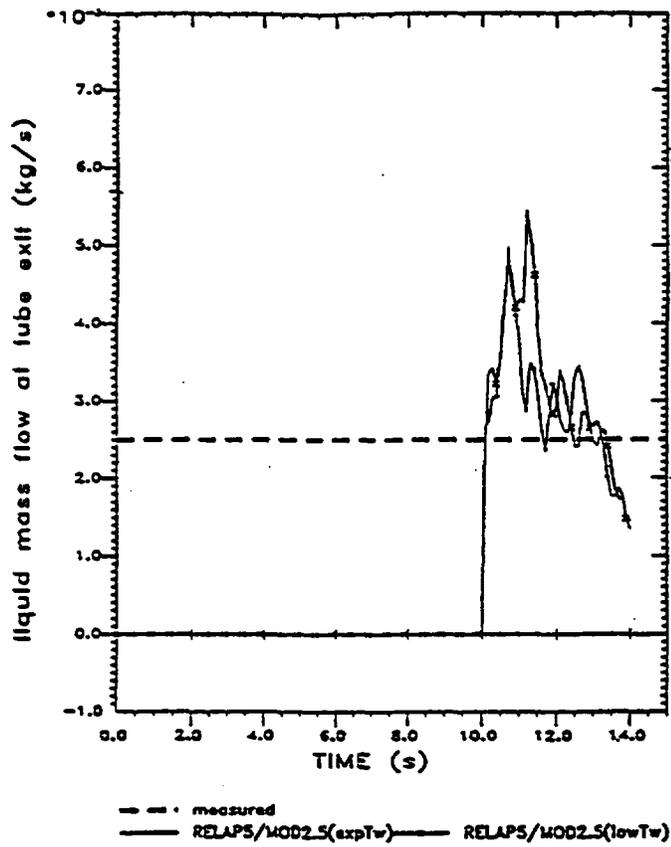


Figure 40: Liquid carry-out for the two initial axial wall temperature distributions.

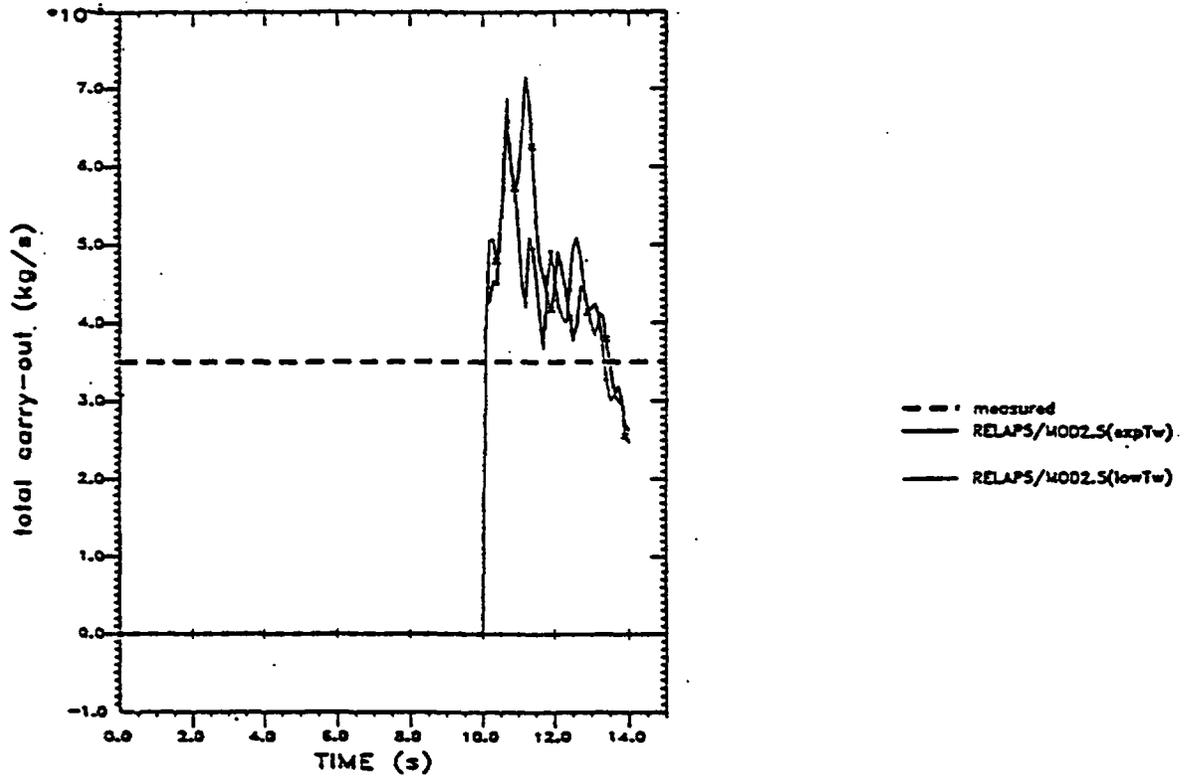


Figure 41: Total mass flow at the tube exit for the two initial axial wall temperature distributions.

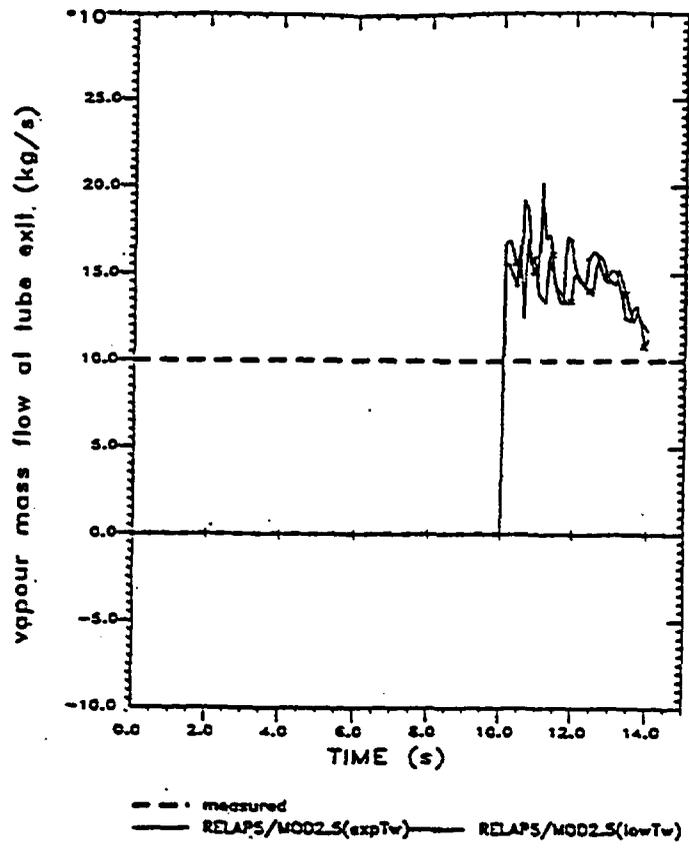


Figure 42: Vapour carry-out for the two initial axial wall temperature distributions.

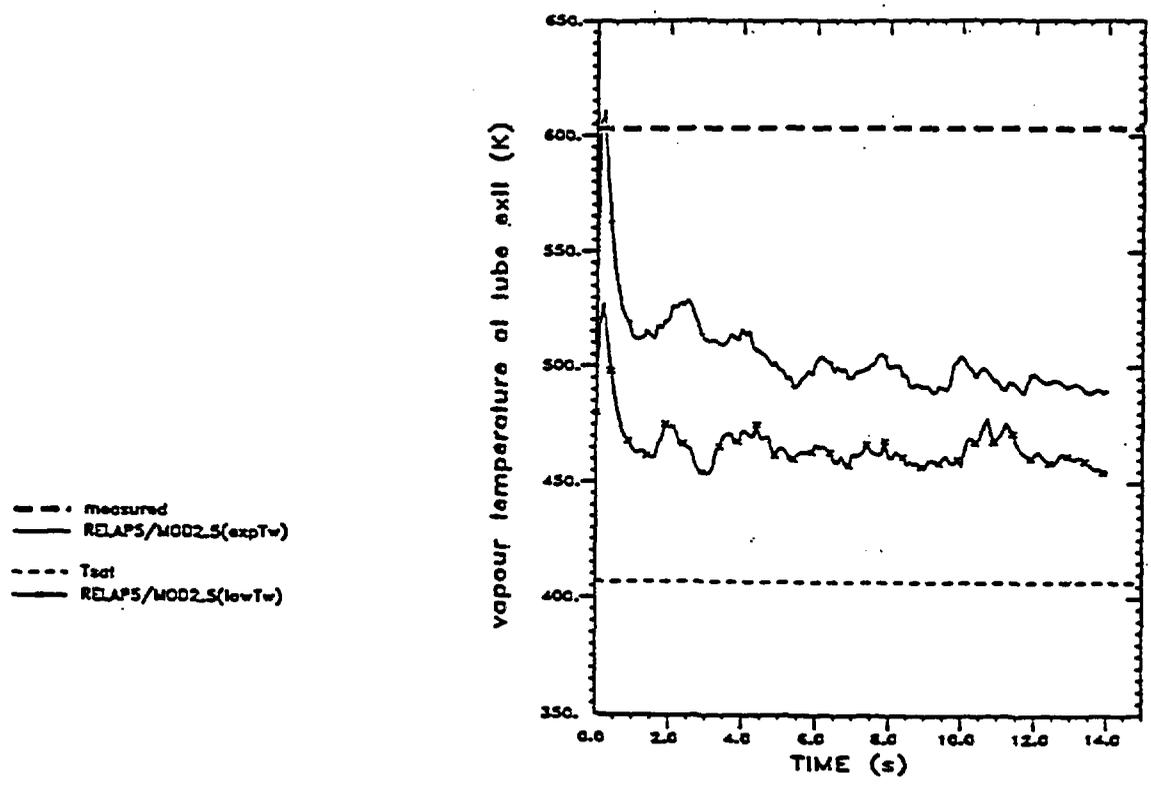


Figure 43: Vapour temperature at the tube exit calculated by RELAPS/MOD2.5 for the two initial wall temperature distributions.

APPENDIX 1

Input decks for transient calculations

[The following content is extremely faint and illegible, appearing to be a large block of text or code, possibly representing input decks for transient calculations. It contains various lines of text that are difficult to discern.]

Appendix IA

RELAP5/Mod 2.5 input deck


```

20510315 0.2 voidf 120140000
20510316 0.2 voidf 120150000
20510317 0.2 voidf 120160000
20510318 0.2 voidf 120170000
20510319 0.2 voidf 120180000
20510320 0.16 voidf 120190000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 112: vapour mass flow from the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511200 entgfl mult 0.00015935 0. 0
20511201 rhogj 125000000
20511202 velgj 125000000
20511203 voidgj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 113: liquid mass flow from the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511300 entffl mult 0.00015935 0. 0
20511301 rhofj 125000000
20511302 velfj 125000000
20511303 voidfj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 114: total mass flow out of the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511400 entflow sum 1. 1. 0
20511401 0.
20511402 1. cntrlvar 112
20511403 1. cntrlvar 113
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 115: quality at the tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511500 exitqual div 1. 1. 0
20511501 cntrlvar 114 cntrlvar 112
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 118: integral of mass flow at tube exit (kg) (R5 var.)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511800 intwgext integral 1. 0. 0
20511801 mflowj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 119: integral of mass flow at tube exit (kg) (cntrlvar)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511900 intwgext integral 1. 0. 0
20511901 cntrlvar 114
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 122: mass of water accumulated (kg)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512200 accwater sum 0.00015935 0. 0
20512201 0.
20512202 0.1 rho 120010000
20512203 0.2 rho 120020000
20512204 0.2 rho 120030000
20512205 0.2 rho 120040000
20512206 0.2 rho 120050000
20512207 0.2 rho 120060000
20512208 0.2 rho 120070000
20512209 0.2 rho 120080000
20512210 0.2 rho 120090000
20512211 0.2 rho 120100000

```


20211006	475.	19.10140
20211007	500.	20.38016
20211008	520.	21.46453
20211009	540.	22.60660
20211010	560.	23.80943
20211011	580.	25.07626
20211012	600.	26.41051
20211013	620.	27.81574
20211014	630.	28.54614
20211015	640.	29.29573
20211016	650.	30.06500
20211017	660.	30.85448
20211018	670.	31.66468
20211019	680.	32.49616
20211020	690.	33.34947
20211021	700.	34.22519
20211022	710.	35.12390
20211023	720.	36.04622
20211024	730.	36.99275
20211025	740.	37.96413
20211026	750.	38.96103
20211027	760.	39.98410
20211028	770.	41.03404
20211029	780.	42.11154
20211030	790.	43.21734
20211031	800.	44.35218
20211032	810.	45.51682
20211033	820.	46.71204
20211034	830.	47.93864
20211035	840.	49.19745
20211036	850.	50.48932
20211037	860.	51.81511
20211038	870.	53.17572
20211039	880.	54.57205
20211040	890.	56.00505
20211041	900.	57.47568
20211042	910.	58.98492
20211043	920.	60.53380
20211044	930.	62.12335
20211045	940.	63.75463
20211046	950.	65.42876
20211047	960.	67.14684
20211048	970.	68.91004
20211049	980.	70.71953
20211050	990.	72.57655
20211051	1000.	74.48233
20211052	1010.	76.43815
20211053	1020.	78.44532
20211054	1030.	80.50520
20211055	1040.	82.61918
20211056	1050.	84.78866
20211057	1060.	87.01512
20211058	1070.	89.30003
20211059	1080.	91.64495
20211060	1090.	94.05144
20211061	1100.	96.52112

*
*---- -1---- -1---- -1---- -1---- -1---- -1----
* tab 120 ambient temperature
*---- -1---- -1---- -1---- -1---- -1---- -1----
*

20212000	temp	
20212001	0.	294.26
20212002	1000.	294.26

Appendix IB

RELAP5/Mod 3 input deck


```

*
1201300 1
1201301 0. 0.0 0. 35 * mass flows
*-----1-----1-----1-----1-----1-----1-----1-----1-
* [125] outlet junction
*-----1-----1-----1-----1-----1-----1-----1-----1-
1250000 "outljun" snljun
1250101 120010000 130000000 0. 0.0 0.0 00000
1250201 1 0. 0.00 0.0
*-----1-----1-----1-----1-----1-----1-----1-----1-
* [130] outlet tank
*-----1-----1-----1-----1-----1-----1-----1-----1-
1300000 "outltan" tmdpv0l
1300101 .01 0. 1000. 0. 0. 0. 0. 0. 0
1300200 003
1300201 0. 2.0+5 812.
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
HEAT STRUCTURES
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
**-----1-----1-----1-----1-----1-----1-----1-----1-*
** [1200] test section(tube)
**-----1-----1-----1-----1-----1-----1-----1-----1-*
11200000 36 10 2 0 0.007125 502 0 16
11200100 0 1
11200101 9 0.00795 * outer tube radius
11200201 1 9
11200301 1. 9
11200400 0
11200401 812. 10 * temperatures of the nodes
11200501 120010000 0 1 1 0.1 1 * inner
11200502 120020000 10000 1 1 0.1 35 * surface
11200503 120190000 0 1 1 0.16 36 *
11200601 -120 0 4110 1 0.1 1 * outer
11200602 -120 0 4110 1 0.1 35 * surface
11200603 -120 0 4110 1 0.16 36 *
11200701 150 0.0273224 0. 0. 01 * power distribution
11200702 150 0.0273224 0. 0. 35 *
11200703 150 0.0437158 0. 0. 36 *
11200801 0 0. 0. 0. 0. 36 * M2
*11200801 0. 11. 11. 0. 0. 0. 0. 1. 18 * M3
*-----1-----1-----1-----1-----1-----1-----1-----1-*
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
CONTROL VARIABLES
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 101: collapsed liquid level in the test section (lower part)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510100 collevl sum 1. 0. 1
20510101 0.
20510102 0.1 voidf 120010000
20510103 0.1 voidf 120020000
20510104 0.1 voidf 120030000
20510105 0.1 voidf 120040000
20510106 0.1 voidf 120050000
20510107 0.1 voidf 120060000
20510108 0.1 voidf 120070000
20510109 0.1 voidf 120080000
20510110 0.1 voidf 120090000
20510111 0.1 voidf 120100000
20510112 0.1 voidf 120110000
20510113 0.1 voidf 120120000
20510114 0.1 voidf 120130000

```

```

20510115 0.1 voidf 120140000
20510116 0.1 voidf 120150000
20510117 0.1 voidf 120160000
20510118 0.1 voidf 120170000
20510119 0.1 voidf 120180000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 102: collapsed liquid level in the test section (upper part)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510200 collevu sum 1. 0. 1
20510201 0.
20510202 0.1 voidf 120190000
20510203 0.1 voidf 120200000
20510204 0.1 voidf 120210000
20510205 0.1 voidf 120220000
20510206 0.1 voidf 120230000
20510207 0.1 voidf 120240000
20510208 0.1 voidf 120250000
20510209 0.1 voidf 120260000
20510210 0.1 voidf 120270000
20510211 0.1 voidf 120280000
20510212 0.1 voidf 120290000
20510213 0.1 voidf 120300000
20510214 0.1 voidf 120310000
20510215 0.1 voidf 120320000
20510216 0.1 voidf 120330000
20510217 0.1 voidf 120340000
20510218 0.1 voidf 120350000
20510219 0.16 voidf 120360000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 103: collapsed liquid level in the test section (m)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510300 collev sum 1. 0. 0
20510301 0.
20510302 1. cntrlvar 101
20510303 1. cntrlvar 102
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 112: vapour mass flow from the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511200 entgfl mult 0.00015935 0. 1
20511201 rhogj 125000000
20511202 velgj 125000000
20511203 voidgj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 113: liquid mass flow from the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511300 entffl mult 0.00015935 0. 1
20511301 rhofj 125000000
20511302 velfj 125000000
20511303 voidfj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 114: total mass flow out of the test section (kg/s)
*-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511400 entflow sum 1. 1. 0
20511401 0.
20511402 1. cntrlvar 112
20511403 1. cntrlvar 113
*

```

```

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 115: quality at the tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511500 exitqual div          1.          1.          0
20511501 cntrlvar 114          cntrlvar 112
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 118: integral of mass flow at tube exit (kg) (R5 var.)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511800 intwgext integral 1.          0.          0
20511801 mflowj 125000000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 119: integral of mass flow at tube exit (kg) (cntrlvar)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511900 intwgext integral 1.          0.          0
20511901 cntrlvar 114
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 120: mass of water accumulated (lower part)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512000 accwatl sum          0.00015935 0.          0
20512001 0.
20512002 0.1 rho          120010000
20512003 0.1 rho          120020000
20512004 0.1 rho          120030000
20512005 0.1 rho          120040000
20512006 0.1 rho          120050000
20512007 0.1 rho          120060000
20512008 0.1 rho          120070000
20512009 0.1 rho          120080000
20512010 0.1 rho          120090000
20512011 0.1 rho          120100000
20512012 0.1 rho          120110000
20512013 0.1 rho          120120000
20512014 0.1 rho          120130000
20512015 0.1 rho          120140000
20512016 0.1 rho          120150000
20512017 0.1 rho          120160000
20512018 0.1 rho          120170000
20512019 0.1 rho          120180000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 121: mass of water accumulated (upper part)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512100 accwatu sum          0.00015935 0.          0
20512101 0.
20512102 0.1 rho          120190000
20512103 0.1 rho          120200000
20512104 0.1 rho          120210000
20512105 0.1 rho          120220000
20512106 0.1 rho          120230000
20512107 0.1 rho          120240000
20512108 0.1 rho          120250000
20512109 0.1 rho          120260000
20512110 0.1 rho          120270000
20512111 0.1 rho          120280000
20512112 0.1 rho          120290000
20512113 0.1 rho          120300000
20512114 0.1 rho          120310000
20512115 0.1 rho          120320000
20512116 0.1 rho          120330000
20512117 0.1 rho          120340000
20512118 0.1 rho          120350000
20512119 0.16 rho          120360000
*

```


20211022	710.	35.12390
20211023	720.	36.04622
20211024	730.	36.99275
20211025	740.	37.96413
20211026	750.	38.96103
20211027	760.	39.98410
20211028	770.	41.03404
20211029	780.	42.11154
20211030	790.	43.21734
20211031	800.	44.35218
20211032	810.	45.51682
20211033	820.	46.71204
20211034	830.	47.93864
20211035	840.	49.19745
20211036	850.	50.48932
20211037	860.	51.81511
20211038	870.	53.17572
20211039	880.	54.57205
20211040	890.	56.00505
20211041	900.	57.47568
20211042	910.	58.98492
20211043	920.	60.53380
20211044	930.	62.12335
20211045	940.	63.75463
20211046	950.	65.42876
20211047	960.	67.14684
20211048	970.	68.91004
20211049	980.	70.71953
20211050	990.	72.57655
20211051	1000.	74.48233
20211052	1010.	76.43815
20211053	1020.	78.44532
20211054	1030.	80.50520
20211055	1040.	82.61918
20211056	1050.	84.78866
20211057	1060.	87.01512
20211058	1070.	89.30003
20211059	1080.	91.64495
20211060	1090.	94.05144
20211061	1100.	96.52112

*
*---- -1----- -1----- -1----- -1----- -1----- -1-----
* tab 120 ambient temperature
*---- -1----- -1----- -1----- -1----- -1----- -1-----
*

20212000	temp	
20212001	0.	294.26
20212002	1000.	294.26

*
*---- -1----- -1----- -1----- -1----- -1----- -1-----
* tab 150 test section power
*---- -1----- -1----- -1----- -1----- -1----- -1-----
*

20215000	power	
20215001	0.	4326.
20215002	1000.	4326.

*


```

= UCB351
*
* UNIV. OF CALIFORNIA BERKELEY (UCB) TEST
* .....
* REFLOOD TEST IN SINGLE TUBE GEOMETRY
* UCB-3051: P=2 BAR, G=25.1 KG/M**2/S, XQ-=0.09
*
* GOAL: CALCULATION OF THE AXIAL VAPOUR
* TEMPERATURE DISTRIBUTION WHEN THE QUENCH FRONT
* IS AT HALF-LENGTH ELEVATION (Z=1.89 M)
*
* REPRESENTATION OF THE SECTION ABOVE QUENCH FRONT ONLY (Z>1.8 M)
* REFLOODING OPTION USED
*
* input: Power and heat losses
*
* RELAP5/MOD2.5 calculation
* for RELAP5/MOD3 only the card nr. 11200801 has to be changed
* .....
*
00000100 new      transnt
00000101 run
*
00000105 5.0      10.0      cpu-lim
00000110 nitrogen      880.
*
* TIME STEP CONTROL CARDS
* .....
* END TIME  MIN DT  MAX DT  OPTN  MNR  MJR  RST
00000201  15      1.0-9  0.01   3    10   400  400
*
* .....
*
* MINOR EDIT REQUESTS
*
* .....
*
00000301 cputime  0
*
00000308 voidg    12010000
00000309 voidg    120110000
00000310 voidg    120120000
00000311 voidg    120130000
00000312 voidg    120140000
00000313 voidg    120150000
00000314 voidg    120170000
00000315 voidg    120190000
*
00000319 tempg    120190000
*
00000320 q          120110000
00000321 qwg     120110000
00000322 q          120120000
00000323 qwg     120120000
00000324 qwg     120130000
00000325 qwg     120140000
00000326 qwg     120150000
00000327 qwg     120160000
00000328 qwg     120170000
00000329 qwg     120180000
00000330 qwg     120190000
*
00000331 mflowj   120110000
00000332 mflowj   120120000
00000333 mflowj   120130000
00000334 mflowj   120140000
00000335 mflowj   120150000
00000336 mflowj   120160000
00000337 mflowj   120170000

```

HYDRODYNAMIC COMPONENTS

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

* [105] feedwater tank - inlet thermodynamic conditions

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

1050000 "feedtan" tmdpvol

1050101 0.00015935 0. 1000. 0. 0. 0. 0. 0

1050200 002 * i.c.: p,X; saturated

1050201 0. 2.15+5 .008

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

* [110] inlet junction - flow boundary conditions

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

1100000 "feedjun" tmdpjun

1100101 105000000 120000000 0.

1100200 1 502 time 0

1100201 -1. 0.0 0.0 0.

1100202 0. 0.0 0.0 0.

1100203 0.0001 0.003707 0.00030 0.

1100204 100.0 0.003707 0.00030 0. *mcell=G*Acell

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

* [120] tube

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

1200000 "testsec" pipe

1200001 19

1200101 0.00015935 19 * Avcell

1200201 0. 18 * Ajunct

1200301 0.03 10 * cell length dz

1200302 0.2 18 *

1200303 0.16 19 *

1200401 0. 19 * cell volume=dz*Avcell

1200501 0. 19 * hor. angle

1200601 90. 19 * vert. angle

1200701 0.03 10 * cell dz

1200702 0.2 18 *

1200703 0.16 19 *

1200801 1.-4 0.01425 19 * friction data

1200901 0. 0. 18 * loss coefficients

*1201002 100 11 * bundle correlation on

1201002 0 19 * tube correlation on

1201101 00000 18

1201201 2 2.00+5 1. 0. 0. 0. 1 * initial

1201202 2 2.00+5 1. 0. 0. 0. 2 * conditions

1201203 2 2.00+5 1. 0. 0. 0. 3

1201204 2 2.00+5 1. 0. 0. 0. 4

1201205 2 2.00+5 1. 0. 0. 0. 5

1201206 2 2.00+5 1. 0. 0. 0. 6

1201207 2 2.00+5 1. 0. 0. 0. 7

1201208 2 2.00+5 1. 0. 0. 0. 8

1201209 2 2.00+5 1. 0. 0. 0. 9

1201210 2 2.00+5 1. 0. 0. 0. 10

1201211 3 2.0+5 420. 0. 0. 0. 11

1201212 3 2.0+5 430. 0. 0. 0. 12

1201213 3 2.0+5 440. 0. 0. 0. 13

1201214 3 2.0+5 450. 0. 0. 0. 14

1201215 3 2.0+5 460. 0. 0. 0. 15

1201216 3 2.0+5 470. 0. 0. 0. 16

1201217 3 2.0+5 480. 0. 0. 0. 17

1201218 3 2.0+5 490. 0. 0. 0. 18

1201219 3 2.0+5 500. 0. 0. 0. 19

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

1201300 1

1201301 0. 0.00030 0. 1 * mass flows

1201302 0. 0.00035 0. 2

1201303 0. 0.00037 0. 3


```

11200413 670. 670. 670. 670. 670. 670. 670. **
+ 670. 670. 670. **
11200414 690. 690. 690. 690. 690. 690. 690. **
+ 690. 690. 690. **
11200415 710. 710. 710. 710. 710. 710. 710. **
+ 710. 710. 710. **
11200416 720. 720. 720. 720. 720. 720. 720. **
+ 720. 720. 720. **
11200417 726. 726. 726. 726. 726. 726. 726. **
+ 726. 726. 726. **
11200418 733. 733. 733. 733. 733. 733. 733. **
+ 733. 733. 733. **
11200419 733. 733. 733. 733. 733. 733. 733. **
+ 733. 733. 733. **
11200501 120010000 10000 1 1 0.03 10 * inner surf
11200502 120110000 10000 1 1 0.2 18 **
11200503 120190000 0 1 1 0.16 19 **
11200601 0 0 0 1 0.03 10 * outer surf
11200602 0 0 2121 1 0.2 11 *
11200603 0 0 2122 1 0.2 12 *
11200604 0 0 2123 1 0.2 13 *
11200605 0 0 2124 1 0.2 14 *
11200606 0 0 2125 1 0.2 15 *
11200607 0 0 2126 1 0.2 16 *
11200608 0 0 2127 1 0.2 17 *
11200609 0 0 2128 1 0.2 18 *
11200610 0 0 2129 1 0.16 19 *
11200701 150 0.00819672 0. 0. 10 * power distribution
11200702 150 0.0546448 0. 0. 18 *
11200703 150 0.0437158 0. 0. 19 *
11200801 0 0. 0. 0. 0. 19 * M2
*11200801 0. 11. 11. 0. 0. 0. 0. 1. 19 * M3
*-----1-----1-----1-----1-----1-----1-----1-----1-----1*
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
* CONTROL VARIABLES
*
*$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$$
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 101: collapsed liquid level in the test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510100 collev sum 1. 0. 1
20510101 0.
20510102 0.03 voidf 120010000
20510103 0.03 voidf 120020000
20510104 0.03 voidf 120030000
20510105 0.03 voidf 120040000
20510106 0.03 voidf 120050000
20510107 0.03 voidf 120060000
20510108 0.03 voidf 120070000
20510109 0.03 voidf 120080000
20510110 0.03 voidf 120090000
20510111 0.03 voidf 120100000
20510112 0.2 voidf 120110000
20510113 0.2 voidf 120120000
20510114 0.2 voidf 120130000
20510115 0.2 voidf 120140000
20510116 0.2 voidf 120150000
20510117 0.2 voidf 120160000
20510118 0.2 voidf 120170000
20510119 0.2 voidf 120180000
20510120 0.16 voidf 120190000
*

```

```

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 102: vapour mass flow from the test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510200 entgfl mult 0.00015935 0. 1
20510201 rhogj 125000000
20510202 velgj 125000000
20510203 voidgj 125000000
20510204 cntrlvar 153
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 103: liquid mass flow from the test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510300 entffl mult 0.00015935 0. 1
20510301 rhofj 125000000
20510302 velfj 125000000
20510303 voidfj 125000000
20510304 cntrlvar 153
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 104: total mass flow out of the test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510400 entflow sum 1. 1.-10 1 1 1.-10
20510401 0.
20510402 1. cntrlvar 102
20510403 1. cntrlvar 103
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 105: quality at the tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510500 exitqual div 1. 0. 1
20510501 cntrlvar 104 cntrlvar 102
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 106: phase velocity difference in volume 13
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510600 diffv13 sum 1. 0. 1
20510601 0.
20510602 1. velg 120130000
20510603 -1. velf 120130000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 107:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510700 abscv106 stdfctn 1. 0. 1
20510701 abs cntrlvar 106
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 108:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510800 cv107s10 mult 1. 0. 1
20510801 cntrlvar 153 cntrlvar 107
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 109:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20510900 cv107int integral 1. 0. 1
20510901 cntrlvar 108
*

```

```

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 110: average phase velocity difference (10-20s) in cell 13
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511000 avdv13 div 1. 0. 1
20511001 cntrlvar 154 cntrlvar 109
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 111: phase velocity difference in volume 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511100 diffv16 sum 1. 0. 1
20511101 0.
20511102 1. velg 120160000
20511103 -1. velf 120160000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 112:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511200 abscv111 stdfctn 1. 0. 1
20511201 abs cntrlvar 111
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 113:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511300 cv112s10 mult 1. 0. 1
20511301 cntrlvar 153 cntrlvar 112
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 114:
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511400 cv113int integral 1. 0. 1
20511401 cntrlvar 113
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 115: average phase velocity difference (10-20s) in cell 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511500 avdv16 div 1. 0. 1
20511501 cntrlvar 154 cntrlvar 114
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 118: vapour mass flow at junction 10
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511800 entgj10 mult 0.00015935 0. 1
20511801 rhogj 120100000
20511802 velgj 120100000
20511803 voidgj 120100000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 119: liquid mass flow at junction 10
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20511900 entfj10 mult 0.00015935 0. 1
20511901 rhofj 120100000
20511902 velfj 120100000
20511903 voidfj 120100000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 120: total mass flow at junction 10
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512000 mflowj10 sum 1. 0. 1
20512001 0.

```

```

20512002 1.          cntrlvar 118
20512003 1.          cntrlvar 119
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 121: quality at junction 10
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512100 qualj10  div      1.          0.          1
20512101 cntrlvar 120      cntrlvar 118
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 122: vapour mass flow at junction 11
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512200 entgj11  mult      0.00015935      0.          1
20512201 rhogj    120110000
20512202 velgj    120110000
20512203 voidgj   120110000
20512204 cntrlvar 153
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 123: liquid mass flow at junction 11
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512300 entfj11  mult      0.00015935      0.          1
20512301 rhofj    120110000
20512302 velfj    120110000
20512303 voidfj   120110000
20512304 cntrlvar 153
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 124: total mass flow at junction 11
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512400 mflowj11 sum      1.          1.-10      1          1          1.-10
20512401 0.
20512402 1.          cntrlvar 122
20512403 1.          cntrlvar 123
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 125: quality at junction 11
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512500 qualj11  div      1.          0.          1
20512501 cntrlvar 124      cntrlvar 122
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 126: vapour mass flow at junction 16
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512600 entgj16  mult      0.00015935      0.          1
20512601 rhogj    120160000
20512602 velgj    120160000
20512603 voidgj   120160000
20512604 cntrlvar 153
*
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
* cv 127: liquid mass flow at junction 16
*---- -1----- -1----- -1----- -1----- -1----- -1----- -1-----
*
20512700 entfj16  mult      0.00015935      0.          1
20512701 rhofj    120160000
20512702 velfj    120160000
20512703 voidfj   120160000
20512704 cntrlvar 153
*

```

```

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 128: total mass flow at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512800 mflowj16 sum      1.      1.-10    1      1      1.-10
20512801 0.
20512802 1.      cntrlvar  126
20512803 1.      cntrlvar  127
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 129: quality at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20512900 qualj16  div      1.      0.      1
20512901 cntrlvar  128      cntrlvar  126
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 130: integral of vapour mass flow at junction 11
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513000 intwgj11  integral 1.      0.      0
20513001 cntrlvar  122
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 131: integral of liquid mass flow at junction 11
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513100 intwfj11  integral 1.      0.      0
20513101 cntrlvar  123
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 132: integral of total mass flow at junction 11
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513200 intwtj11  sum      1.      1.-10    1      1      1.-10
20513201 0.
20513202 1.      cntrlvar  130
20513203 1.      cntrlvar  131
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 133: average total mass flow at junction 11
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513300 avwtj11  div      1.      0.      0
20513301 cntrlvar  154      cntrlvar  132
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 134: average quality at junction 11
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513400 avxaj11  div      1.      0.      0
20513401 cntrlvar  132      cntrlvar  130
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 135: integral of vapour mass flow at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513500 intwgj16  integral 1.      0.      0
20513501 cntrlvar  126
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 136: integral of liquid mass flow at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513600 intwfj16  integral 1.      0.      0
20513601 cntrlvar  127
*

```

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*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 137: integral of total mass flow at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513700 intwtj16 sum 1. 1.-10 1 1 1.-10
20513701 0.
20513702 1. cntrlvar 135
20513703 1. cntrlvar 136
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 138: average total mass flow at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513800 avwtj16 div 1. 0. 0
20513801 cntrlvar 154 cntrlvar 137
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 139: average quality at junction 16
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20513900 avxaj16 div 1. 0. 0
20513901 cntrlvar 137 cntrlvar 135
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 140: integral of vapour mass flow at tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514000 intwgext integral 1. 0. 0
20514001 cntrlvar 102
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 141: integral of liquid mass flow at tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514100 intwfext integral 1. 0. 0
20514101 cntrlvar 103
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 142: integral of total mass flow at tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514200 intwtext sum 1. 1.-10 1 1 1.-10
20514201 0.
20514202 1. cntrlvar 140
20514203 1. cntrlvar 141
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 143: average total mass flow at tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514300 avwtext div 1. 0. 0
20514301 cntrlvar 154 cntrlvar 142
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 144: average quality at tube exit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514400 avxaext div 1. 0. 0
20514401 cntrlvar 142 cntrlvar 140
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 149: equilibrium quality at the quench front
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20514900 xeqqf sum 1. 1. 0. 0
20514901 0.075
20514902 1.13355258-4 cntrlvar 150
*

```

```

*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 150: power input to wet portion of test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20515000 qinwet      sum      1.      0.      1
20515001 0.
20515002 1.          qq      120010000
20515003 1.          qq      120020000
20515004 1.          qq      120030000
20515005 1.          qq      120040000
20515006 1.          qq      120050000
20515007 1.          qb      120060000
20515008 1.          qq      120070000
20515009 1.          qq      120080000
20515010 1.          qq      120090000
20515011 1.          q      120100000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 151: power input to dry portion of test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20515100 qindry      sum      1.      0.      1
20515101 0.
20515102 1.          q      120110000
20515103 1.          qq      120120000
20515104 1.          qq      120130000
20515105 1.          qq      120140000
20515106 1.          qq      120150000
20515107 1.          qq      120160000
20515108 1.          qq      120170000
20515109 1.          qq      120180000
20515110 1.          q      120190000
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 152: total heat input to fluid in test section
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20515200 qintot      sum      1.      0.      1
20515201 0.
20515202 1.          cntrlvar 150
20515203 1.          cntrlvar 151
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 153: trip unit
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20515300 trip      tripunit 1.      0.      0
20515301 503
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 154: delay
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
20515400 delay      sum      1.      0.0000001 1
20515401 -10.
20515402 1.          time      0
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 160: heat added to wet part (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516000 htwet10      mult      1.      0.      0
20516001 cntrlvar 150      cntrlvar 153
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 161: heat added to dry part (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-

```

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*
20516100 htdry10 mult 1. 0. 0
20516101 cntrlvar 151 cntrlvar 153
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 162: heat added to fluid (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516200 httot10 mult 1. 0. 0
20516201 cntrlvar 152 cntrlvar 153
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 163: total heat added to fluid in wet region (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516300 htinwet integral 1. 0. 1
20516301 cntrlvar 160
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 164: total heat added to fluid in dry region (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516400 htindry integral 1. 0. 1
20516401 cntrlvar 161
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 165: total heat added to fluid (after 10 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516500 htintot integral 1. 0. 1
20516501 cntrlvar 162
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 166: average heat transfer rate in wet region (10-20 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516600 avhtwet div 1. 0. 0
20516601 cntrlvar 154 cntrlvar 163
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 167: average heat transfer rate in dry region (10-20 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516700 avhtdry div 1. 0. 0
20516701 cntrlvar 154 cntrlvar 164
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 168: average heat transfer rate to fluid (10-20 s)
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516800 avhttot div 1. 0. 0
20516801 cntrlvar 154 cntrlvar 165
*
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
* cv 169: heat transfer to vapour only in the dry region
*-----1-----1-----1-----1-----1-----1-----1-----1-----1-
*
20516900 qdryvap sum 1. 0. 1
20516901 0.
20516902 1. qwg 120110000
20516903 1. qwg 120120000
20516904 1. qwg 120130000
20516905 1. qwg 120140000
20516906 1. qwg 120150000
20516907 1. qwg 120160000
20516908 1. qwg 120170000
20516909 1. qwg 120180000
20516910 1. qwg 120190000
*

```



```

*
20211500 htrnrate
20211501 0.      10980.
20211502 1000.  10980.
*
20211600 htrnrate
20211601 0.      9700.
20211602 1000.  9700.
*
20211700 htrnrate
20211701 0.      8843.
20211702 1000.  8843.
*
20211800 htrnrate
20211801 0.      8195.
20211802 1000.  8195.
*
20211900 htrnrate
20211901 0.      8195.
20211902 1000.  8195.
*
20212100 htrnrate
20212101 0.      8494.
20212102 1000.  8494.
*
20212200 htrnrate
20212201 0.      9802.
20212202 1000.  9802.
*
20212300 htrnrate
20212301 0.      11174.
20212302 1000.  11174.
*
20212400 htrnrate
20212401 0.      12553.
20212402 1000.  12553.
*
20212500 htrnrate
20212501 0.      13823.
20212502 1000.  13823.
*
20212600 htrnrate
20212601 0.      14973.
20212602 1000.  14973.
*
20212700 htrnrate
20212701 0.      15737.
20212702 1000.  15737.
*
20212800 htrnrate
20212801 0.      16317.
20212802 1000.  16317.
*
20212900 htrnrate
20212901 0.      16317.
20212902 1000.  16317.
*
*-----1-----1-----1-----1-----1-----1-----
* tab 150 test section power
*-----1-----1-----1-----1-----1-----1-----
*
20215000 power
20215001 0.      4326.
20215002 1000.  4326.
*

```


NRC FORM 335 (2-89) NRCM 1102, 3201, 3202	U.S. NUCLEAR REGULATORY COMMISSION	1. REPORT NUMBER (Assigned by NRC, Add Vol., Supp., Rev., and Addendum Numbers, if any.)												
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11. ABSTRACT (200 words or less) For conditions of low quality at the quench front and low mass flux, the interfacial heat transfer calculated assuming a uniform distribution of droplets over the cross-sectional area of the channel is necessarily overpredicted, and the vapour superheat is strongly underpredicted. The purpose of the present paper is to show that limitations of the 1-D approach, obtained from the steady-state analyses of slow reflooding experiments, has some impact on the performance of the 1-D transient computer codes like RELAP5/MOD2.5 and RELAP5/MOD3. As an example, the transient analysis of a low flooding rate experiment in a tube was performed. An early completion of the quench process and a fast desuperheating of the vapour at the tube exit was obtained by both codes. The too high quench front velocity (four times higher than in the experiment) could not, however, be put univocally in relation to the underprediction of the vapour temperature, and the consequent increase of the precursory cooling, as many coupled thermal and hydraulic transient effects prevailed. Quasi steady-state analyses of two runs, where the boundary conditions for the post-dryout region could be better controlled for a predetermined position of the quench front, were thus performed. These analyses show that the vapour superheat at the tube exit is strongly underpredicted, confirming the limitations of the 1-D model for interfacial heat transfer in the dispersed flow region.														
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